

La physique du mouillage



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Outline

➤ Surface tension

➤ Statics of wetting

- ✓ Contact angle
- ✓ Hysteresis
- ✓ Structured surfaces
- ✓ Nanoscale wetting
- ✓ Wetting of fibers

➤ Wetting dynamics: the moving contact line issue

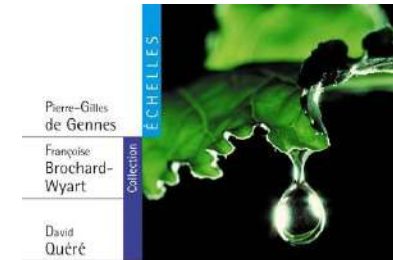
- ✓ Dynamical contact angle
- ✓ Spreading dynamics
- ✓ Simulations
- ✓ Inertial wetting
- ✓ Impact

➤ Wetting in complex situations

- ✓ Volatile liquids
- ✓ Reactive liquids
- ✓ Surfactants

➤ Wetting and adsorption

- ✓ Contamination by airborne molecules or particles
- ✓ Cleaning



Gouttes, bulles,
perles et ondes

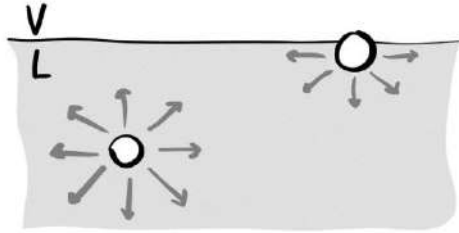
NOUVELLE ÉDITION AVEC CD-ROM

Belin

De Gennes, *Rev. Mod. Phys.* 1986
Bonn *et al.*, *Rev. Mod. Phys.* 2009

Surface tension

➤ Surface energy



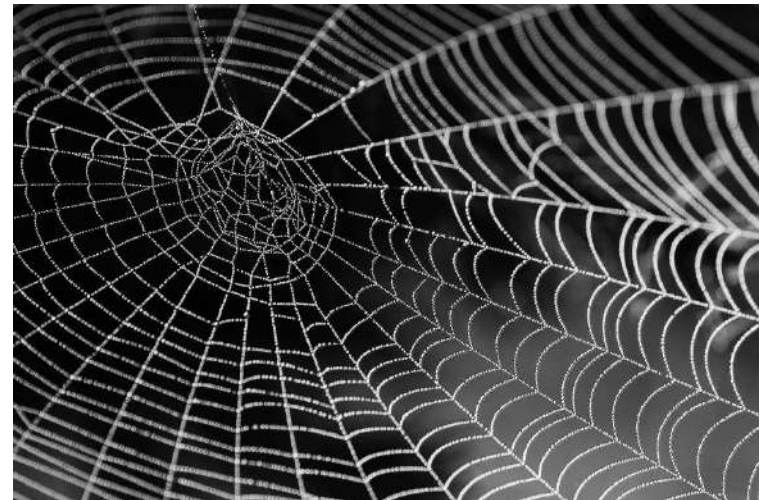
$$\gamma = \frac{\text{energy}}{\text{surface}} \sim \frac{kT}{a^2}$$

$$\gamma \sim \frac{4 \cdot 10^{-21}}{(2 \cdot 10^{-10})^2} \sim 100 \text{ mJ/m}^2$$



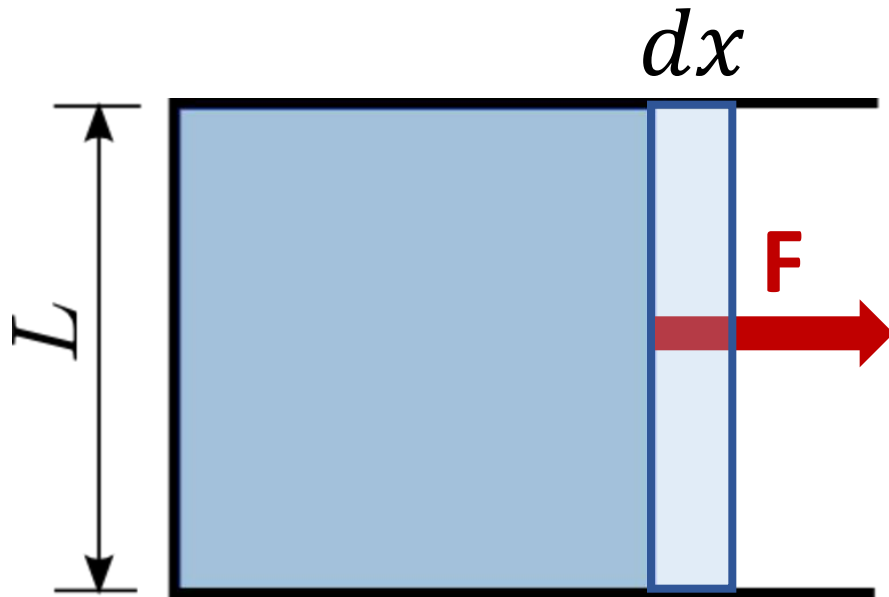
Surface area minimization

✓ Ex: Rayleigh-Plateau instability



Interfacial tension

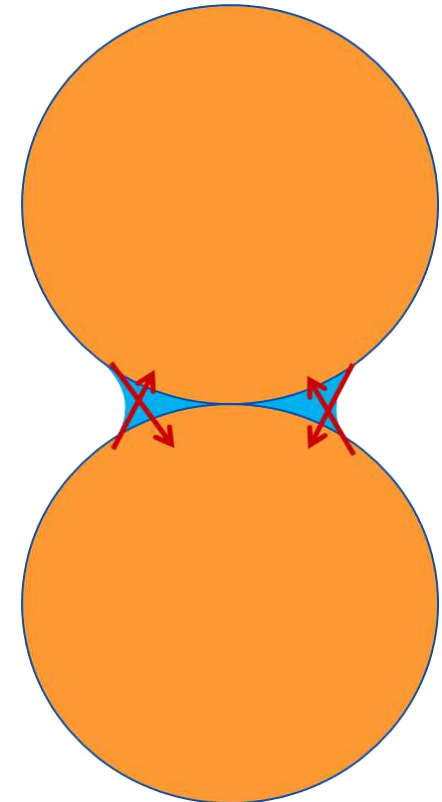
➤ Force per unit length



$$\delta W = \gamma L dx = F \cdot dx$$

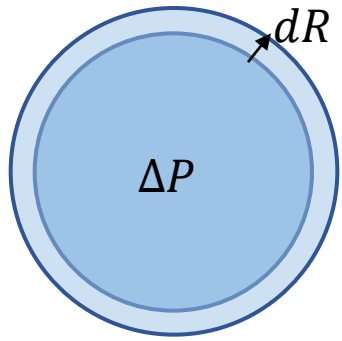
$$\gamma = \frac{F}{L}$$

✓ Ex: sand castle



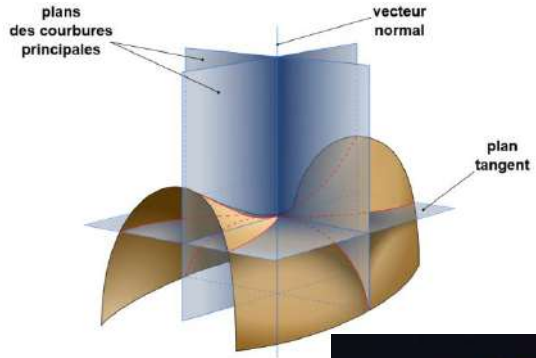
Interfacial tension

➤ Laplace pressure



$$\delta W = \Delta P \cdot 4\pi R^2 dR = \gamma d(4\pi R^2) = \gamma 8\pi R dR$$

$$\Delta P = \frac{2\gamma}{R}$$



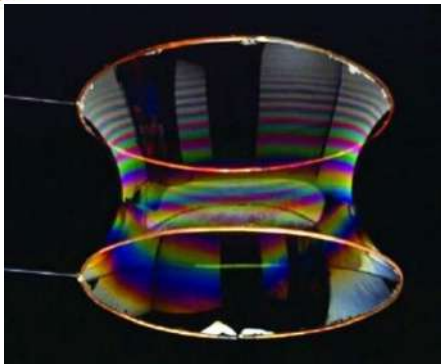
$$\Delta P = \gamma \left(\frac{1}{R_1} + \frac{1}{R_2} \right)$$

R > 0 or < 0

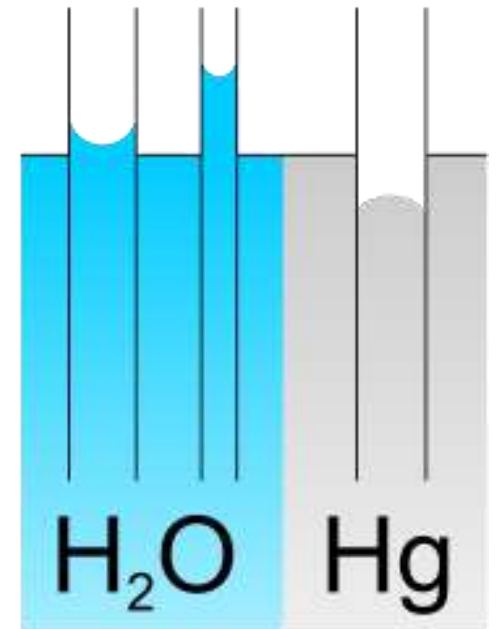
✓ Ex : loi de Jurin

$$\Delta P_{Laplace} + \rho g h = 0$$

$$h = \pm \frac{2\gamma}{\rho g R}$$

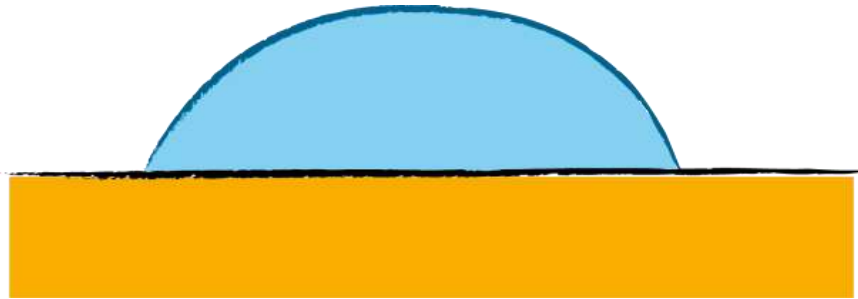


$$\begin{array}{cc} \Delta P_{Laplace} & \\ -\frac{2\gamma}{R} & \frac{2\gamma}{R} \\ h > 0 & h < 0 \end{array}$$



Capillary length

➤ Laplace pressure vs. hydrostatic pressure



$$\frac{2\gamma}{h} \sim \rho g h$$

$$K^{-1} = \sqrt{\frac{\gamma}{\rho g}} \sim 2 \text{ mm}$$

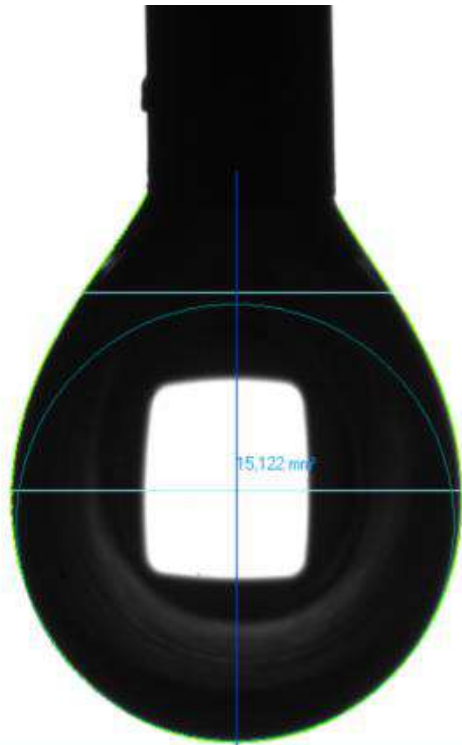


Surface tension measurement

➤ Measure tension

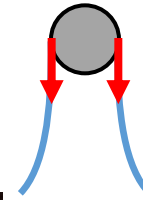
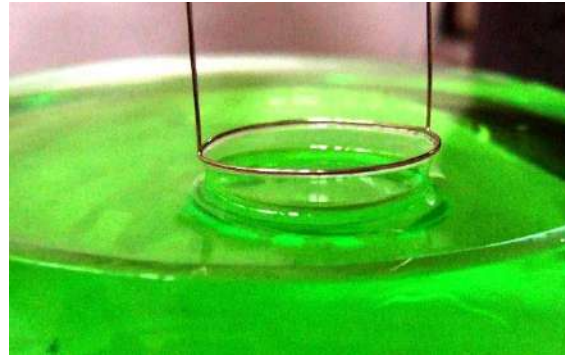
✓ Pendant drop method

Balance surface tension force - weight



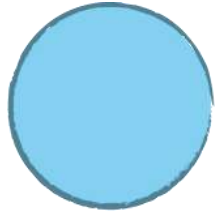
✓ Du Noüy tensiometer

$$F = 2 * 2\pi r\gamma$$



Liquid	γ (mN/m)
Water	72,8
Glycerol	63
Ethylene glycol	48
Ethanol	22,4
Hexane	17,9
PDMS	20,4
Mercure	486

Wetting: thermodynamics



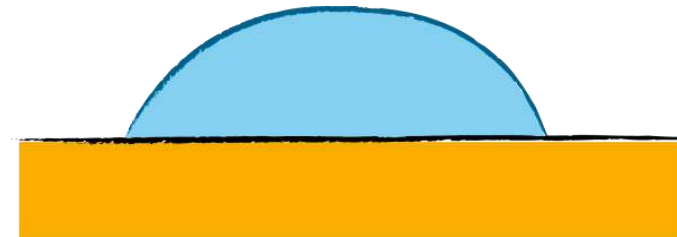
➤ Spreading parameter

$$S = \gamma_{SG} - (\gamma_{SL} + \gamma)$$

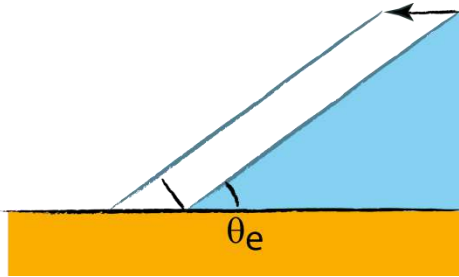
- $S > 0$ Total wetting



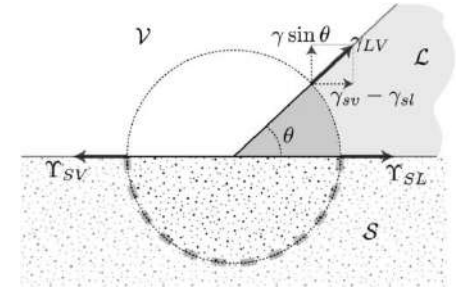
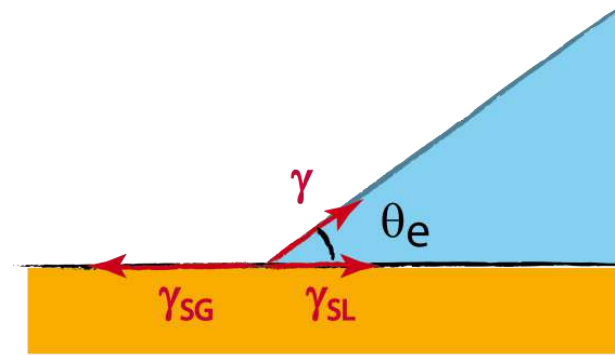
- $S < 0$ Partial wetting



Partial wetting



$$\cos \theta_e = \frac{\gamma_{SG} - \gamma_{SL}}{\gamma}$$



Marchand et al., Am. J Phys. 2011

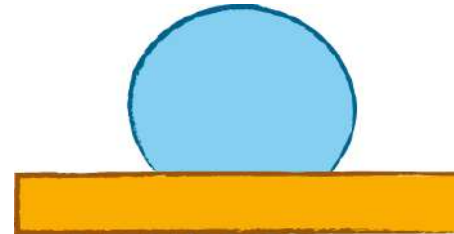
$$dE = (\gamma_{SL} - \gamma_{SG})dx + \gamma \cdot \cos \theta \cdot dx$$

Hydrophilic



$$\theta_e < 90^\circ$$

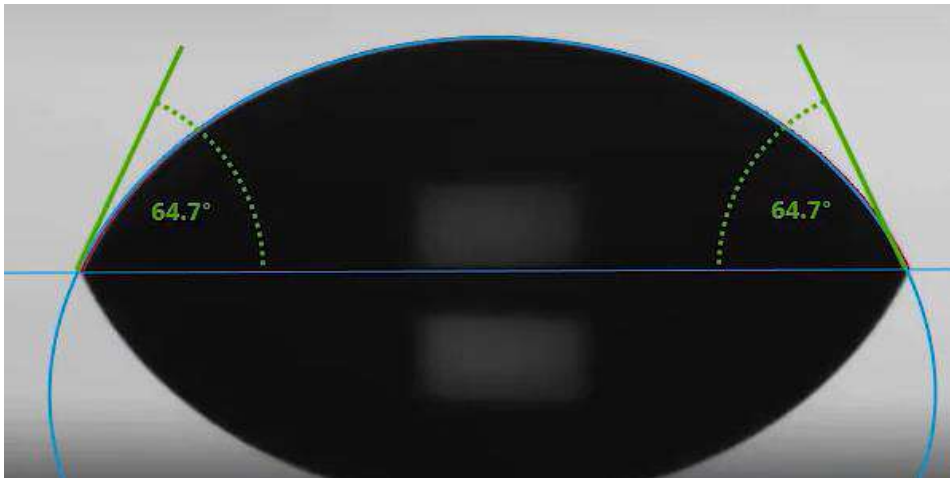
Hydrophobic



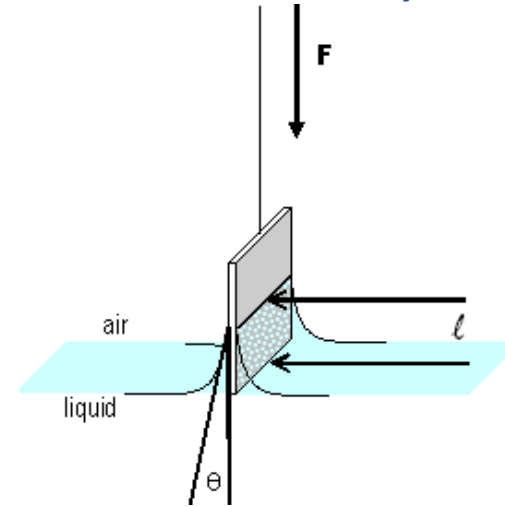
$$\theta_e > 90^\circ$$

Contact angle measurements

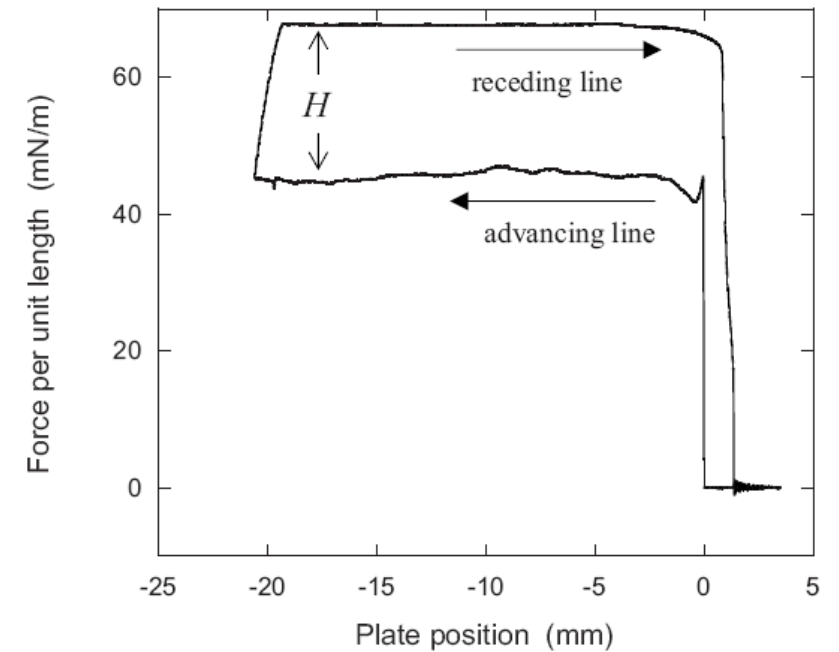
➤ Goniometer



➤ Wilhelmy balance

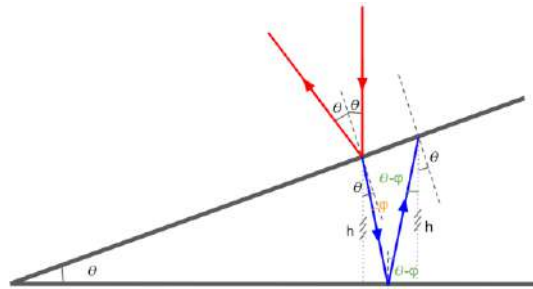
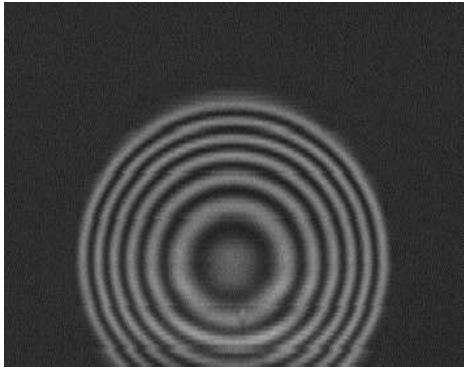


$$F = 2L\gamma \cos \theta$$



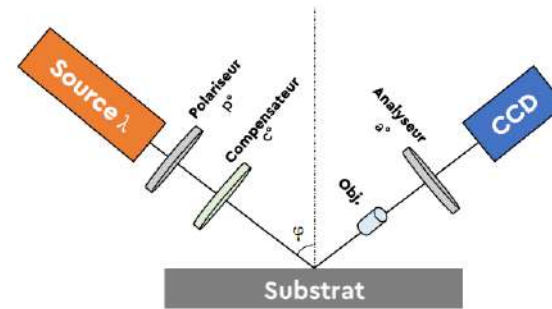
Contact angle measurements

➤ Interferences

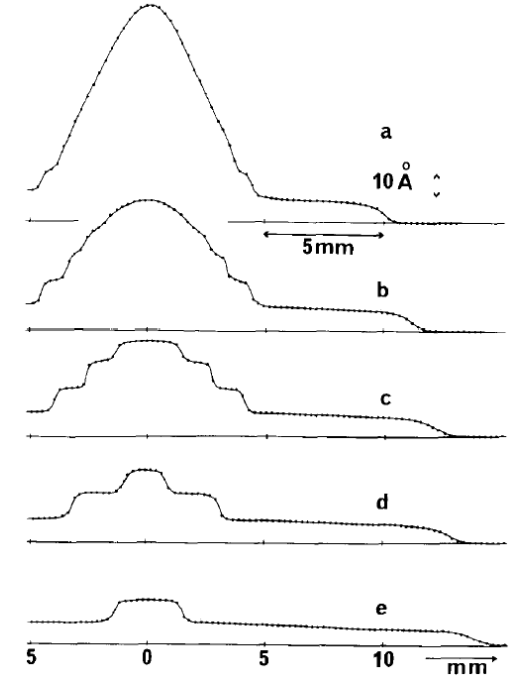


$$\Delta h = \frac{\lambda}{2n}$$

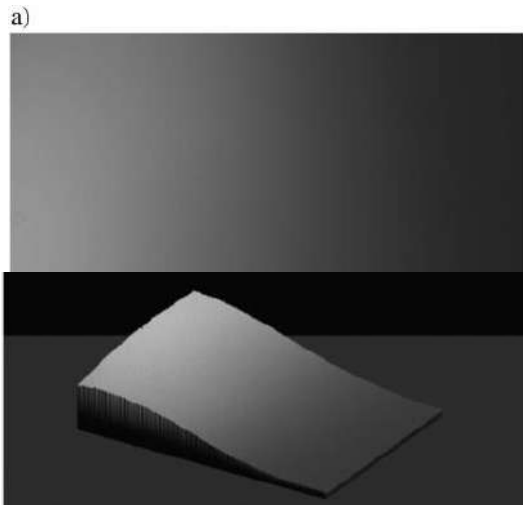
➤ Ellipsométrie



Heslot *et al.*, Nature 1989

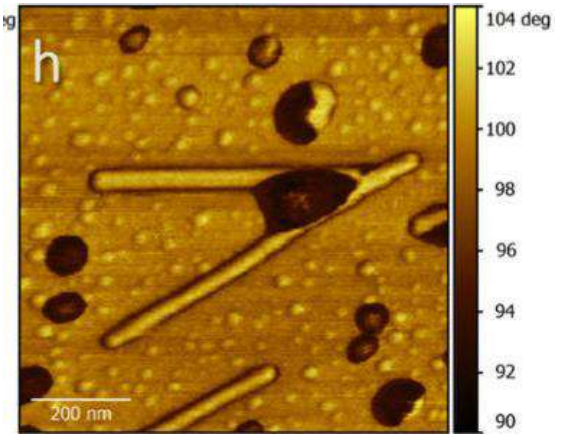
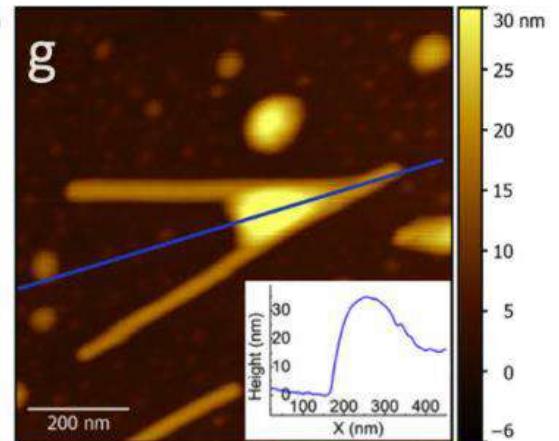


➤ Fluorescence



Huang *et al.*, Phys. Rev. Lett. 2011

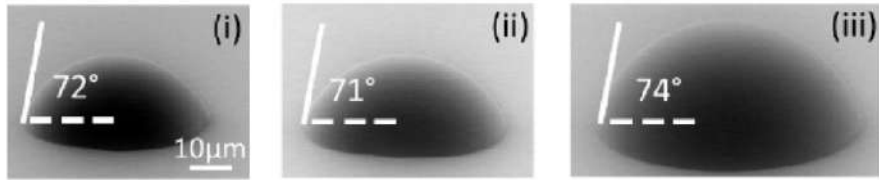
➤ AFM



Calo *et al.*, Molecules. 2021

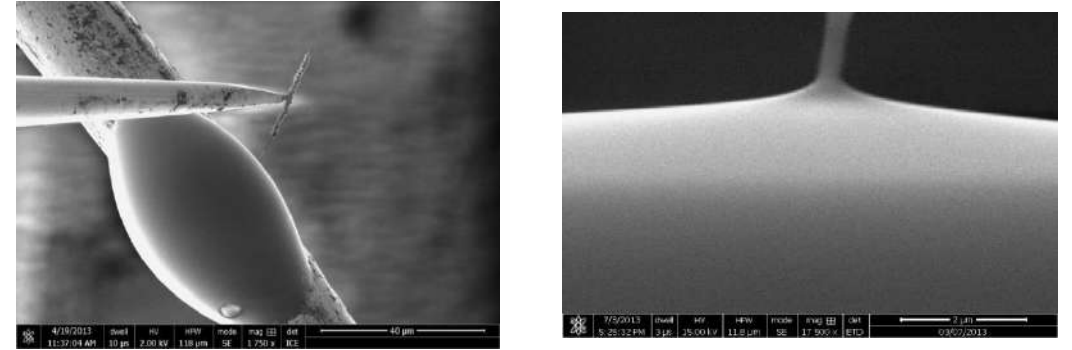
Contact angle measurements

➤ ESEM



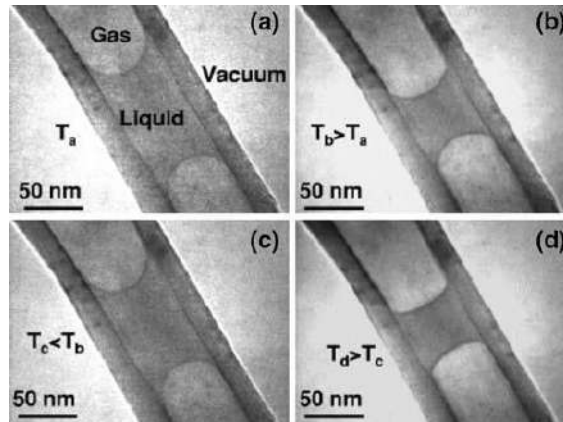
Chatre *et al.*, *Langmuir* 2023

➤ SEM with ionic liquids



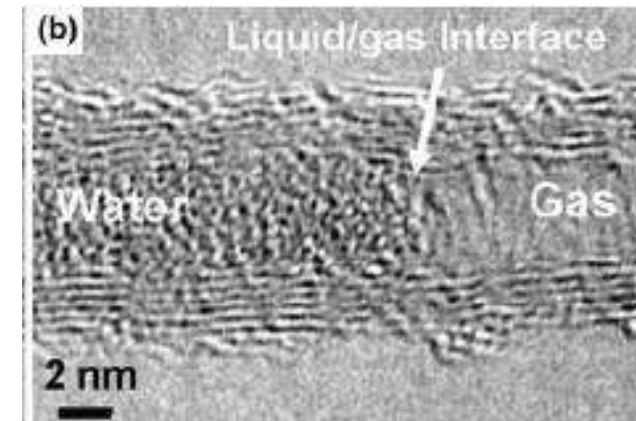
Dupré de Baubigny *et al.*, *Langmuir*. 2015

➤ TEM



Gogotsi *et al.*, *Appl. Phys. Lett.* 2001

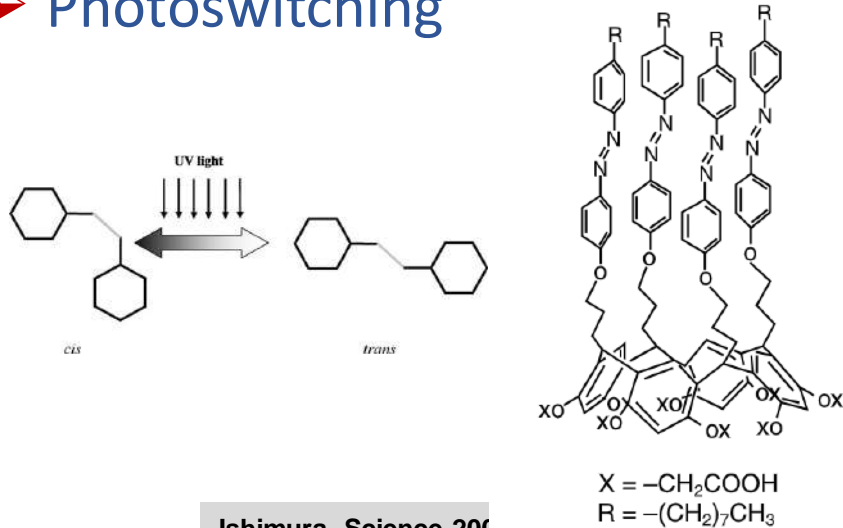
➤ HRTEM



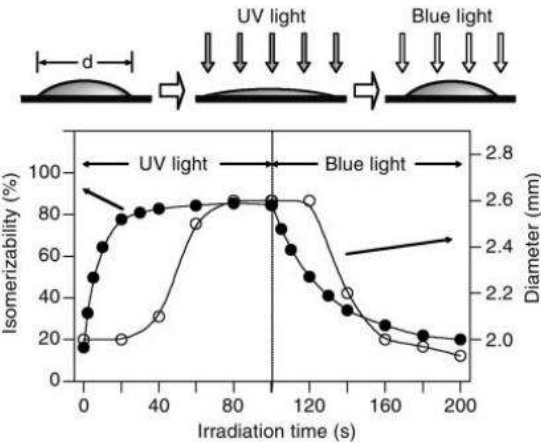
Naguib *et al.*, *Nano. Lett.* 2004

Switchable surfaces

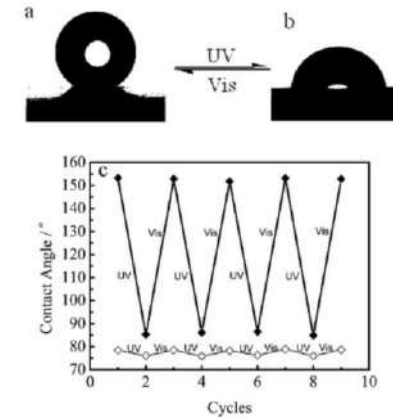
➤ Photoswitching



Ishimura, Science 200

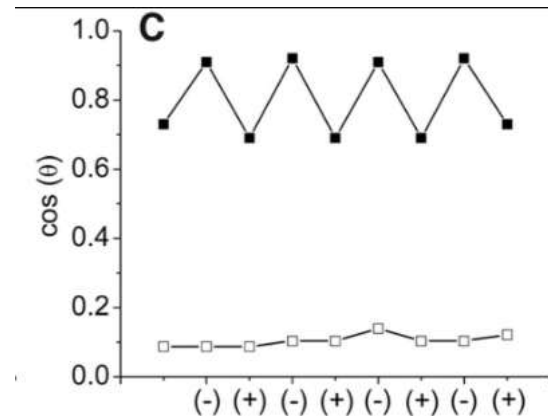
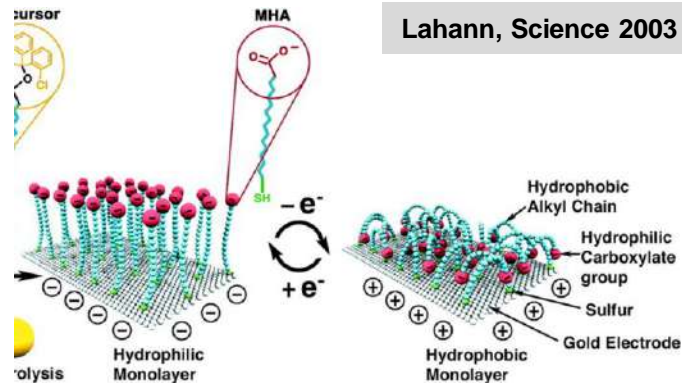


On rough surfaces



Jiang et al., Chem. Comm. 2005

➤ Photoswitching

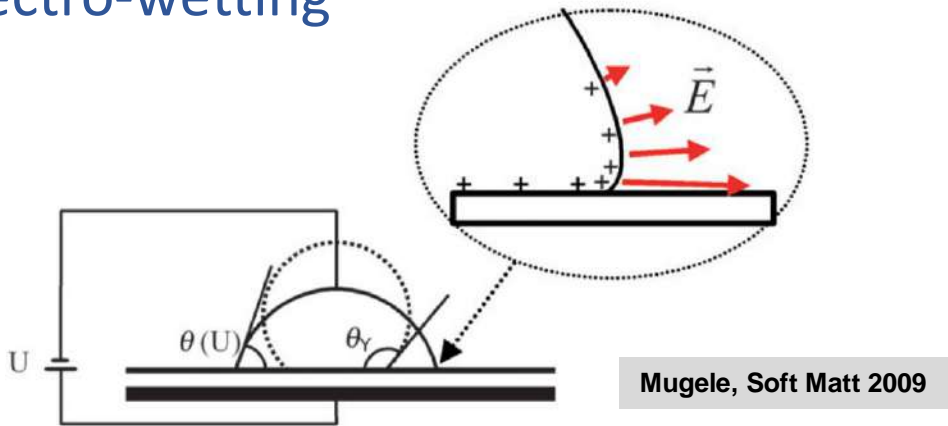


➤ Others

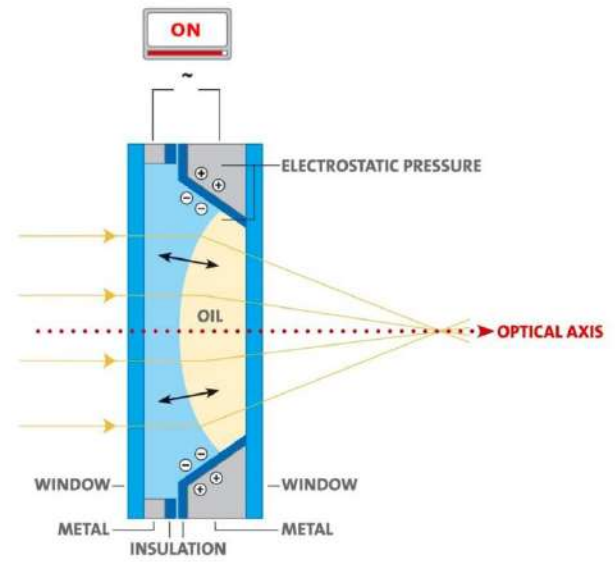
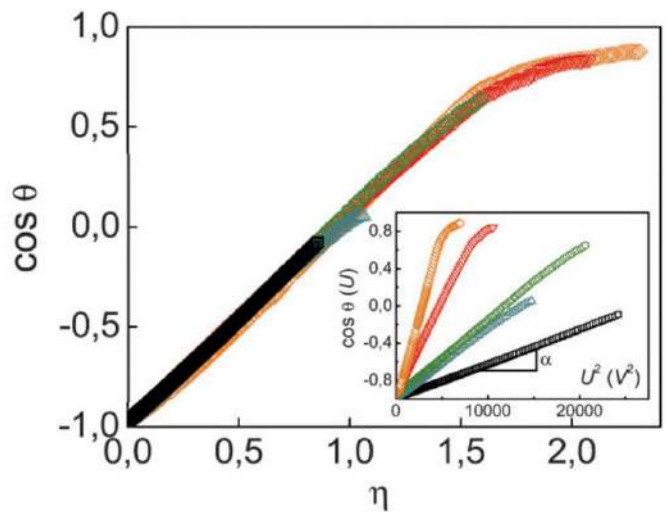
Gras et al., ChemPhysChem 2007

Controlling contact angle

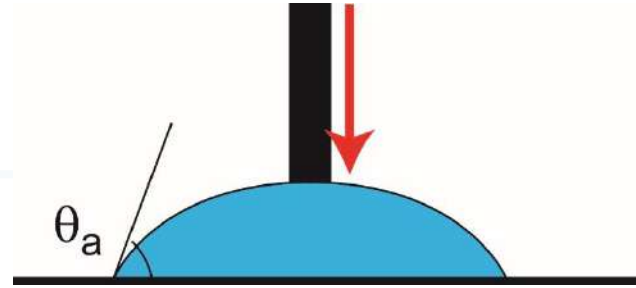
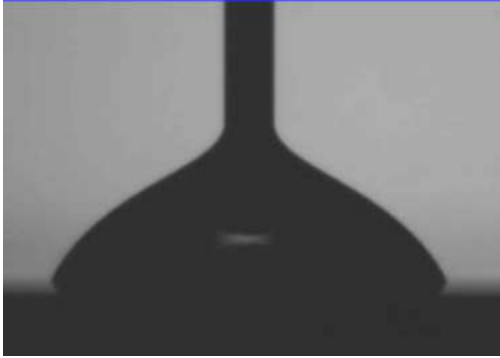
➤ Electro-wetting



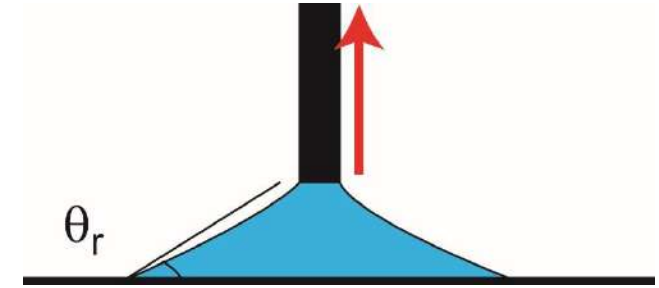
$$\cos \theta = \cos \theta_e + \frac{\epsilon \epsilon_0 U^2}{2\gamma d}$$



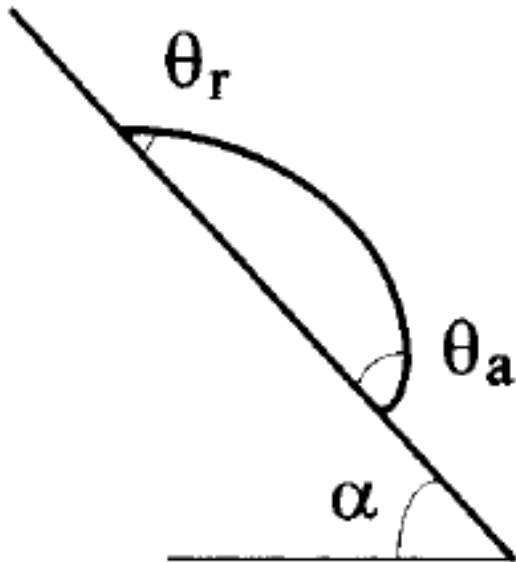
Contact angle hysteresis



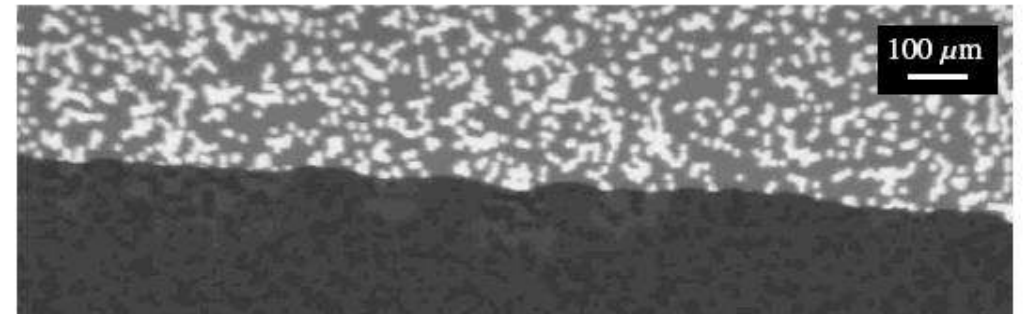
Advancing



Receding



Origin: pinning of the contact line on surface defects

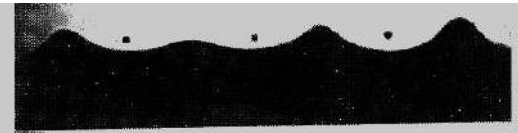


$$\pi r \gamma (\cos \theta_r - \cos \theta_a) \geq \rho \Omega g \sin \alpha$$

Contact angle hysteresis: individual defect

➤ Contact line elastic energy

$$E_{el} = \frac{1}{4} \gamma \sin^2 \theta \int |q| |\tilde{\eta}(q)|^2 \frac{dq}{2\pi}.$$

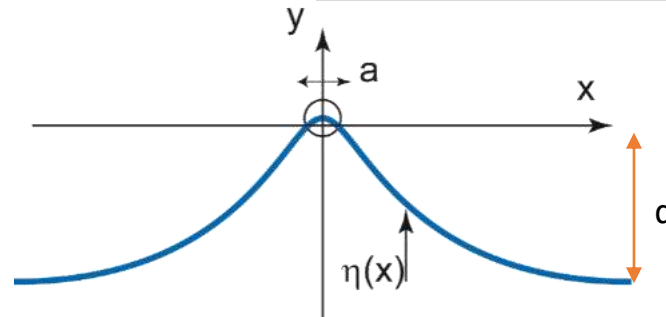


Ondarçuhu, Veyssié *Nature* 1991

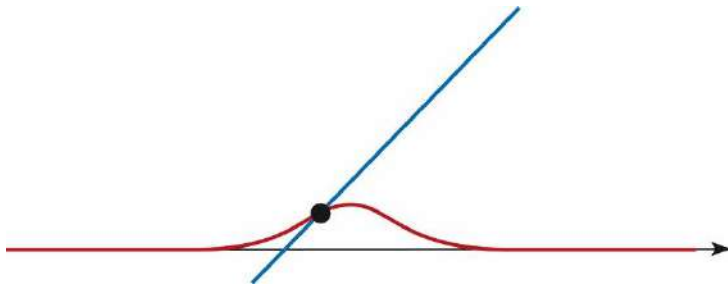
For localized defect: $F = k \cdot y$ with

$$k = \frac{\pi \gamma \sin^2 \theta_e}{\ln(L/a)}$$

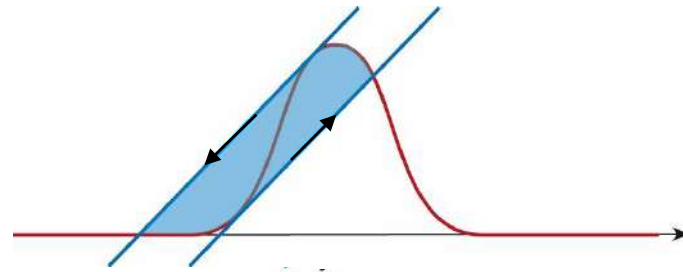
Joanny and de Gennes *J. Chem. Phys.* 1984



➤ Balance of forces



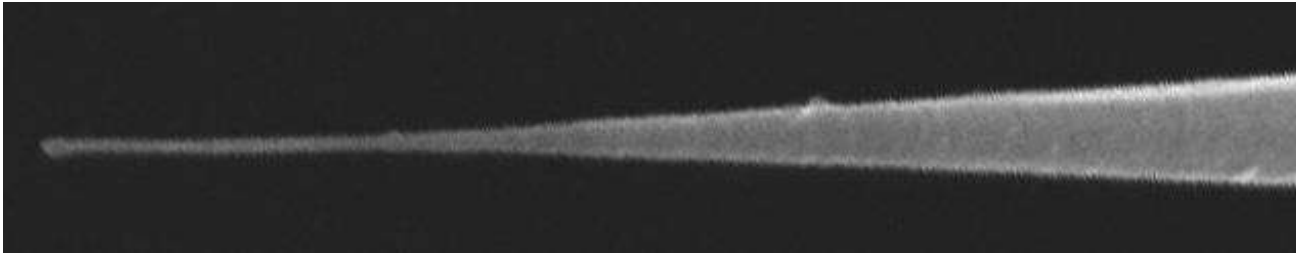
Weak defect



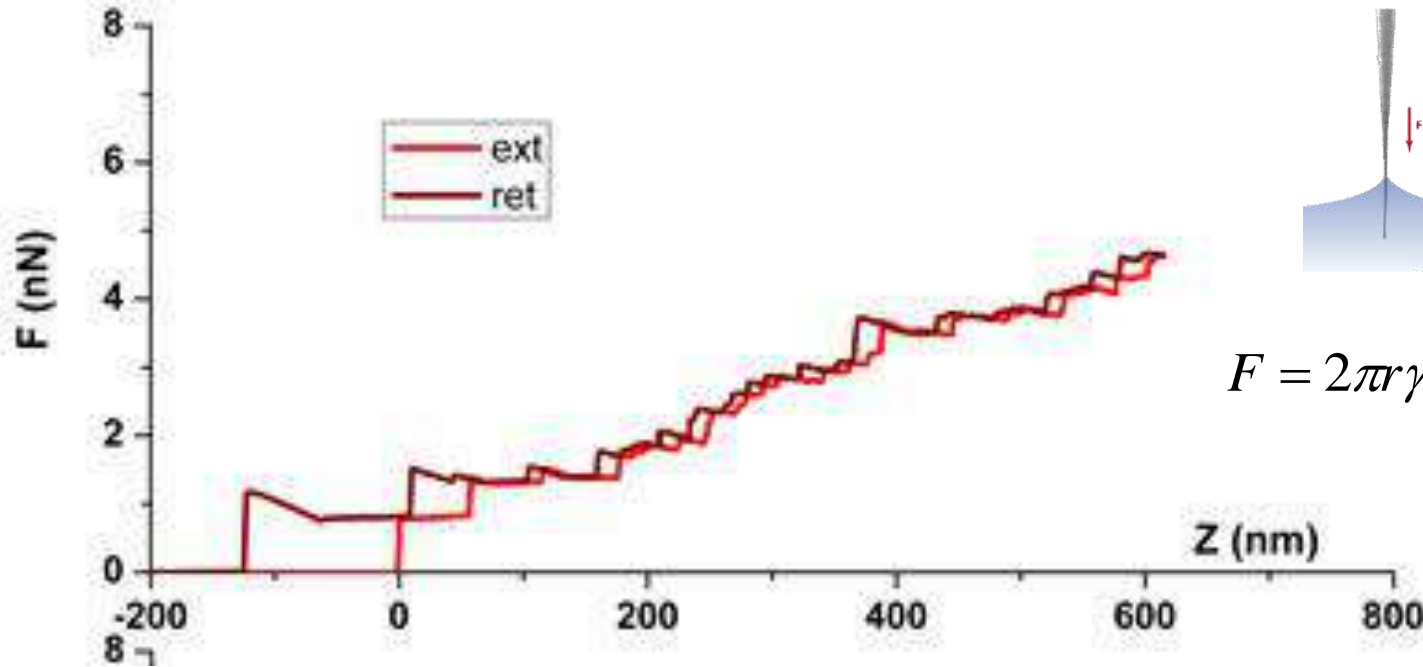
Strong defect

Contact angle hysteresis: individual defect

$\phi = 11 \text{ nm}$

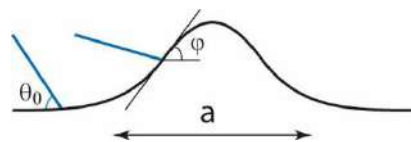
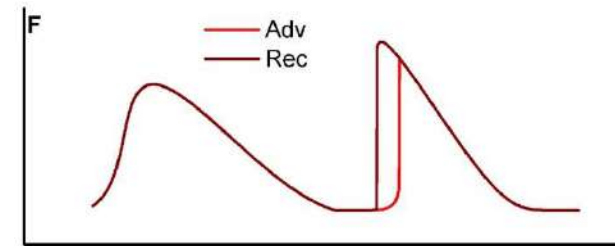
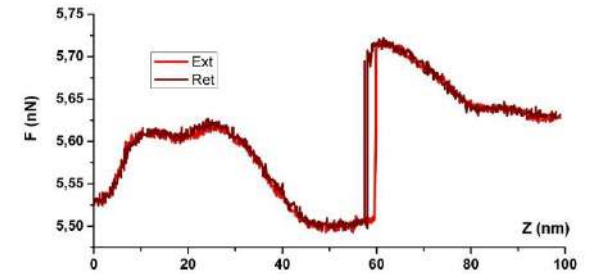


Delmas *et al.* PRL 2011



$$F = 2\pi r \gamma \cos \theta$$

Weak defect Strong defect

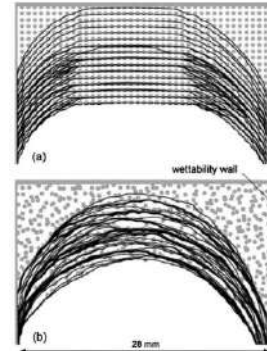
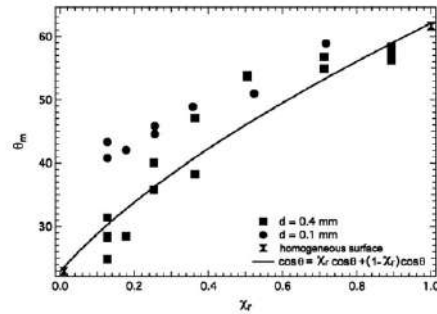


$$W = \frac{F_0^2}{2k} > kT \implies$$

$$a_c = \left(\frac{4\pi}{\text{Ln}(L/a) \sin^2 \varphi} \right)^{1/2} b \approx 1,5 b$$

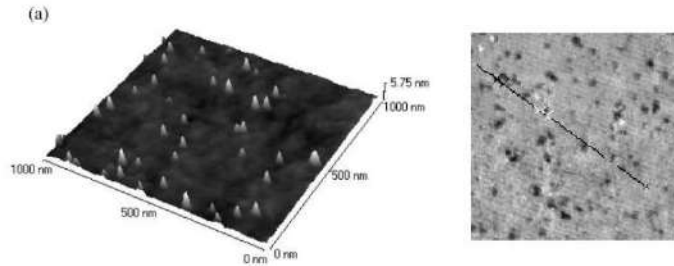
Contact angle hysteresis: collective effects

➤ Patterned substrates ($> 10 \mu\text{m}$)



Cubaud, Fermigier *JCIS* 2004

➤ Nanometric defects

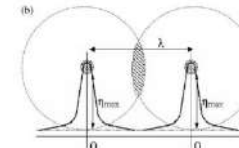
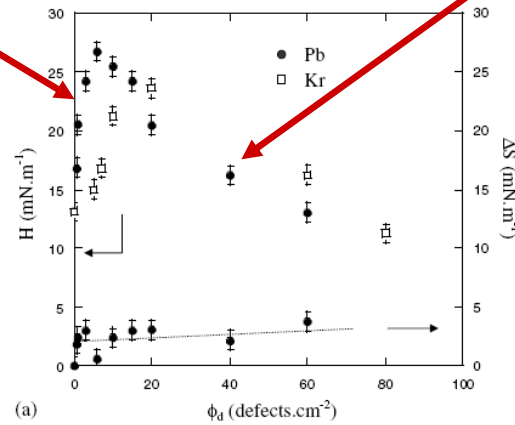


$\varnothing \sim 20\text{-}200 \text{ nm}$

Ramos et al *Surf. Sci.* 2003
Eur. Phys. J E 2006

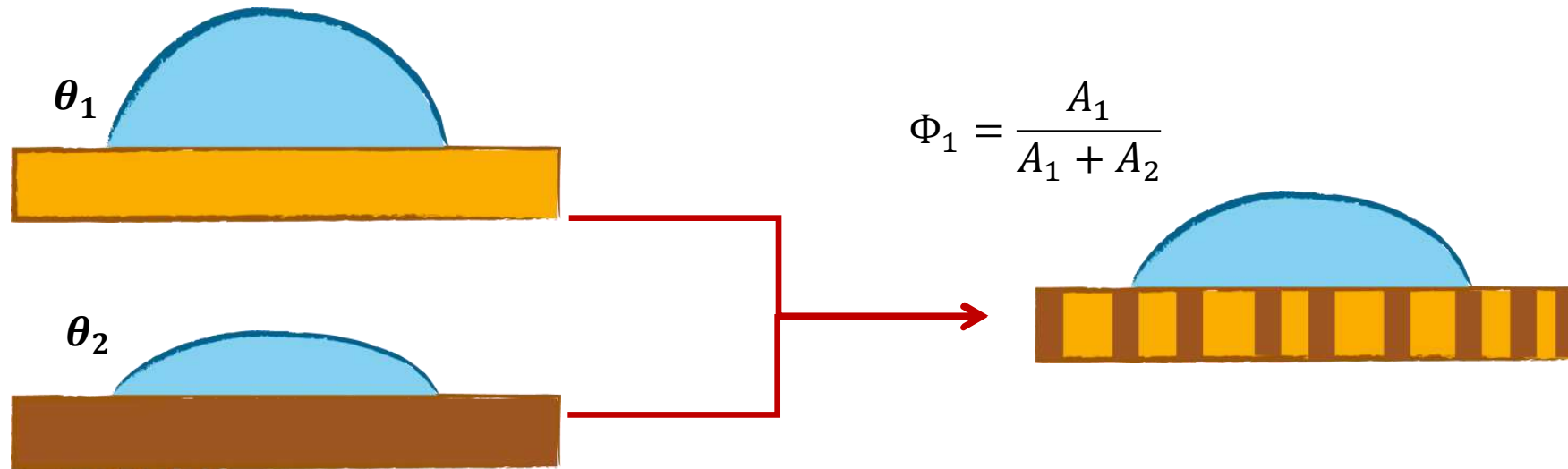
Individual pinning

Collective pinning



➤ Pb : reference surface without hysteresis

Chemical heterogeneity: Cassie-Baxter model

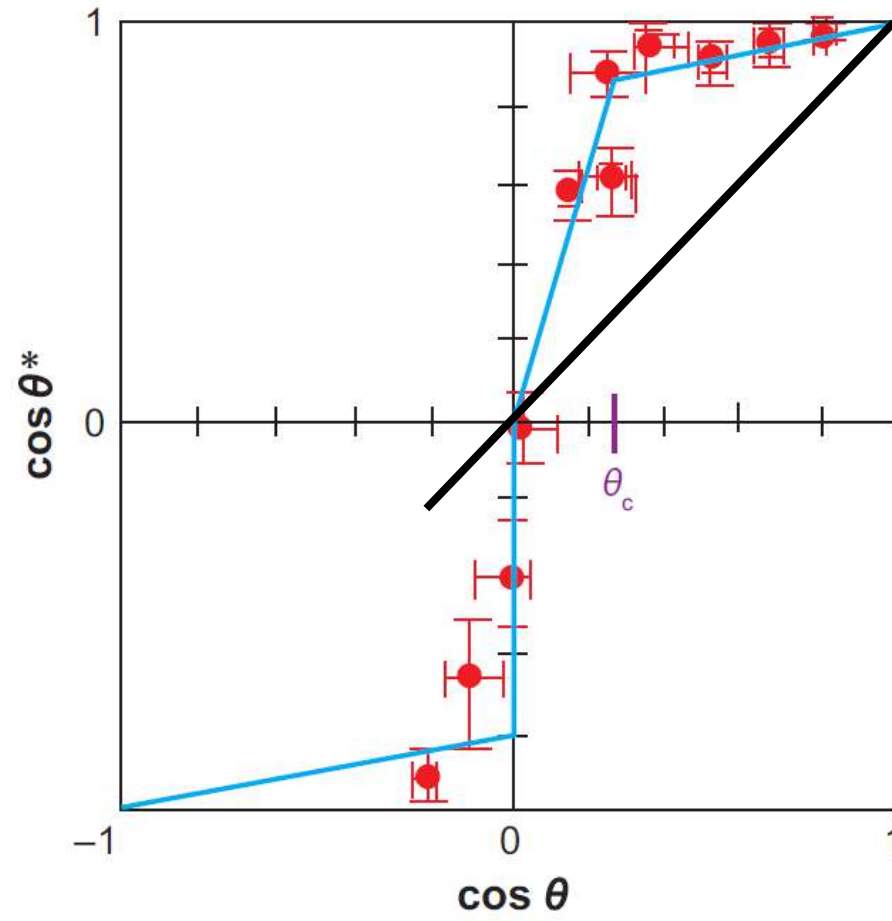


$$dE = [\Phi_1(\gamma_{1L} - \gamma_{1G}) + \Phi_2(\gamma_{2L} - \gamma_{2G})]. dx + \gamma. \cos \theta. dx$$

$$\cos \theta^* = \Phi_1 \cos \theta_1 + (1 - \Phi_1) \cos \theta_2$$

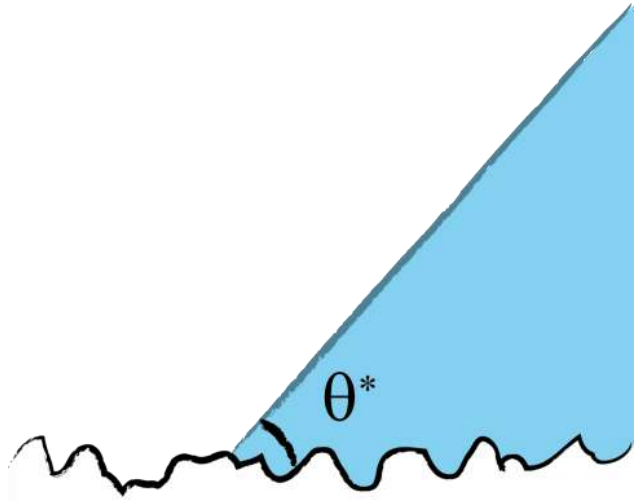
Rough substrates

Quéré, *Ann. Rev. Mat. Res.* 2009



Onda *et al.*, *Langmuir* 1996

Rough substrates: Wenzel state

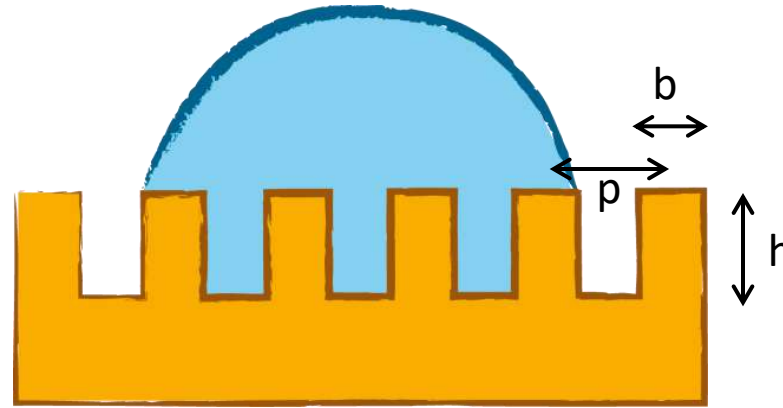


$$dE = r(\gamma_{SL} - \gamma_{SV})dx + \gamma \cdot \cos \theta \cdot dx$$

$$\cos \theta^* = r \cdot \cos \theta_e$$

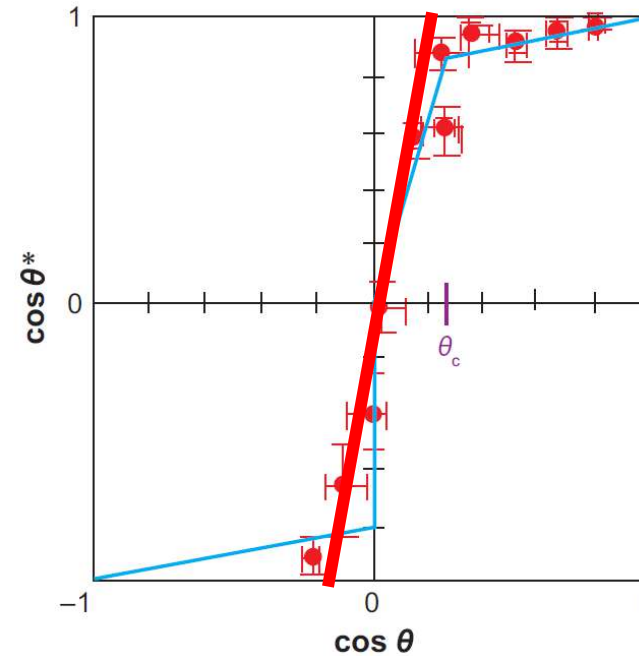
Wenzel, *Ind. Eng. Chem.* 1936

$$r = \frac{\text{actual surface area}}{\text{projected surface area}} = \text{roughness factor}$$

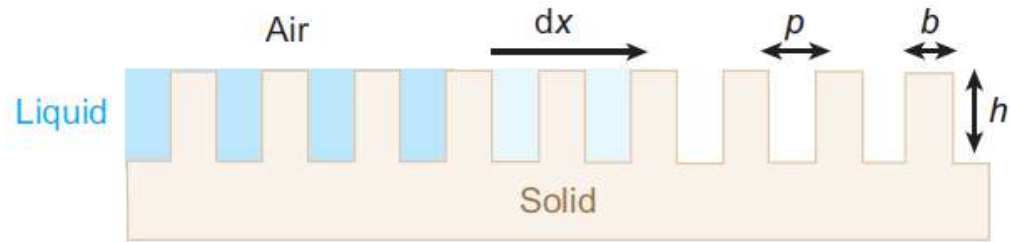


$$r = \frac{p + 2h}{p} = 1 + 2\frac{h}{p}$$

$$\Phi_s = \frac{b}{p}$$



Rough hydrophilic: superhydrophilicity

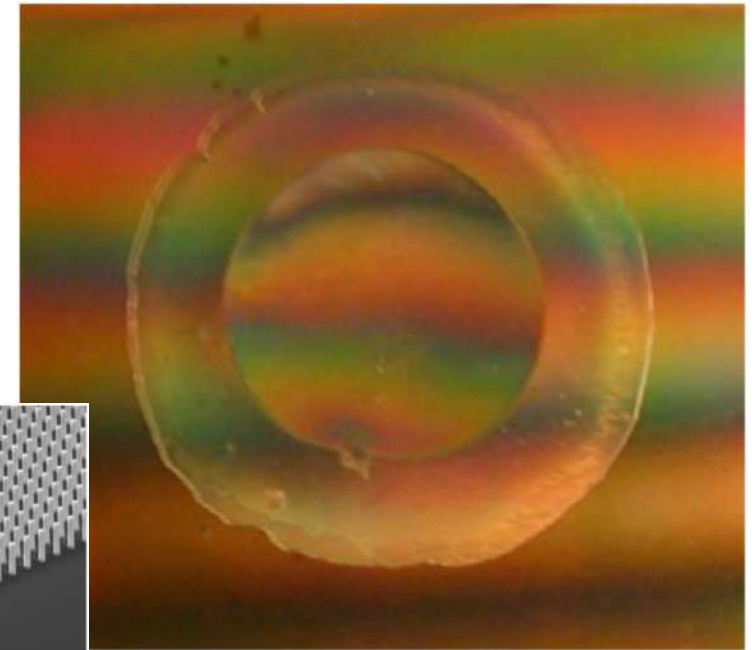


$$r = 1 + 2 \frac{h}{p}$$

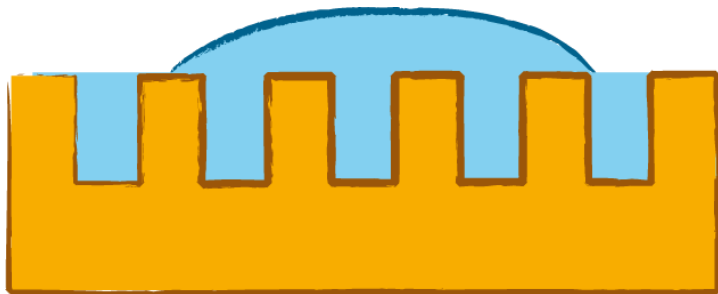
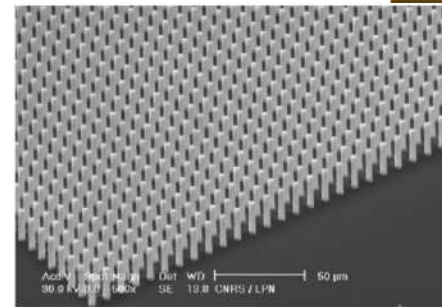
$$\Phi_s = \frac{b}{p}$$

$$dE = (\gamma_{SL} - \gamma_{SG})(r - \Phi_s) + \gamma(1 - \Phi_s)$$

$$\cos \theta_c = \frac{1 - \phi_s}{r - \phi_s}$$



Ichino et al., EPL 2007



$$\theta_1 = \theta_e$$

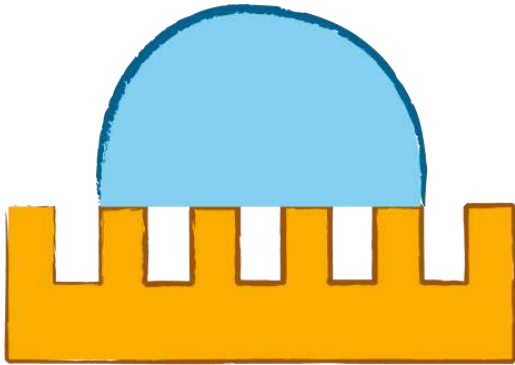
$$\theta_2 = 0^\circ$$

$$\cos \theta^* = \phi_s \cos \theta_e + (1 - \phi_s)$$

$$\cos \theta^* = 1 - \phi_s(1 - \cos \theta_e)$$

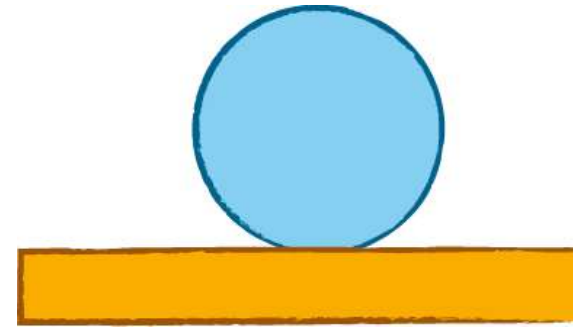
Rough hydrophobic: superhydrophobicity

$$\cos \theta_c = -\frac{1 - \Phi_S}{r - \Phi_S}$$

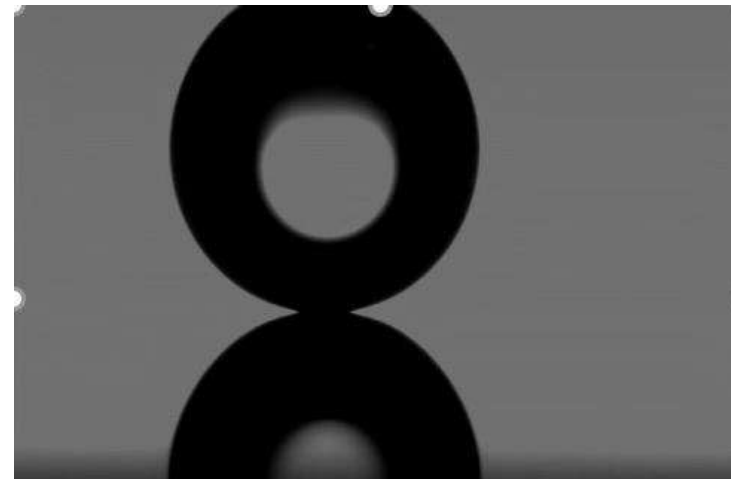
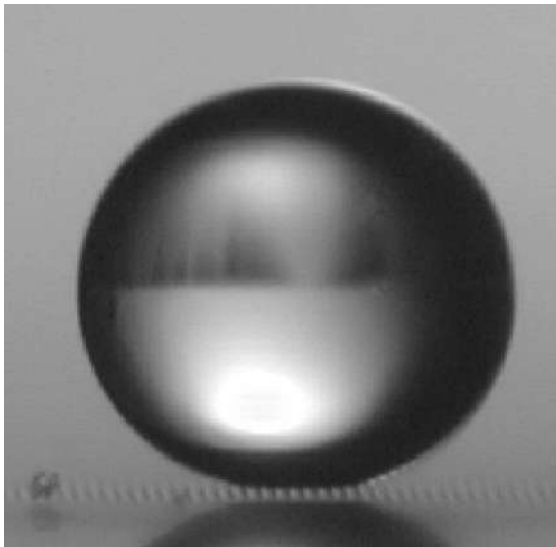


$$\begin{aligned}\theta_1 &= \theta_e \\ \theta_2 &= 180^\circ\end{aligned}$$

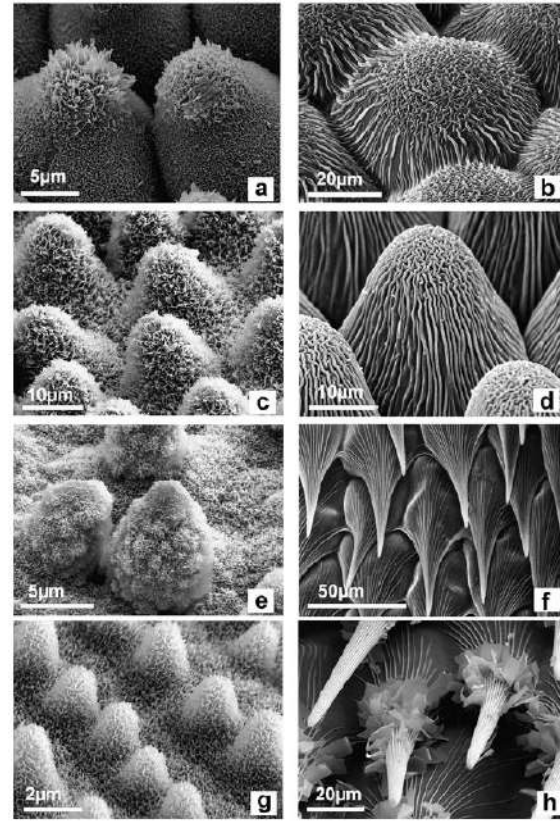
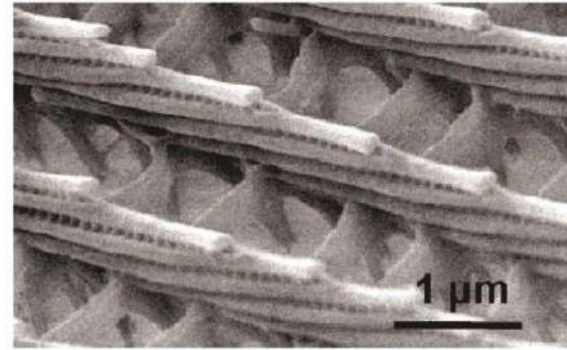
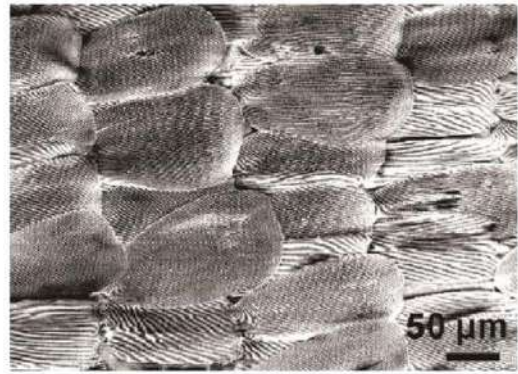
$$\cos \theta^* = -1 + \phi_S(1 - \cos \theta_e)$$



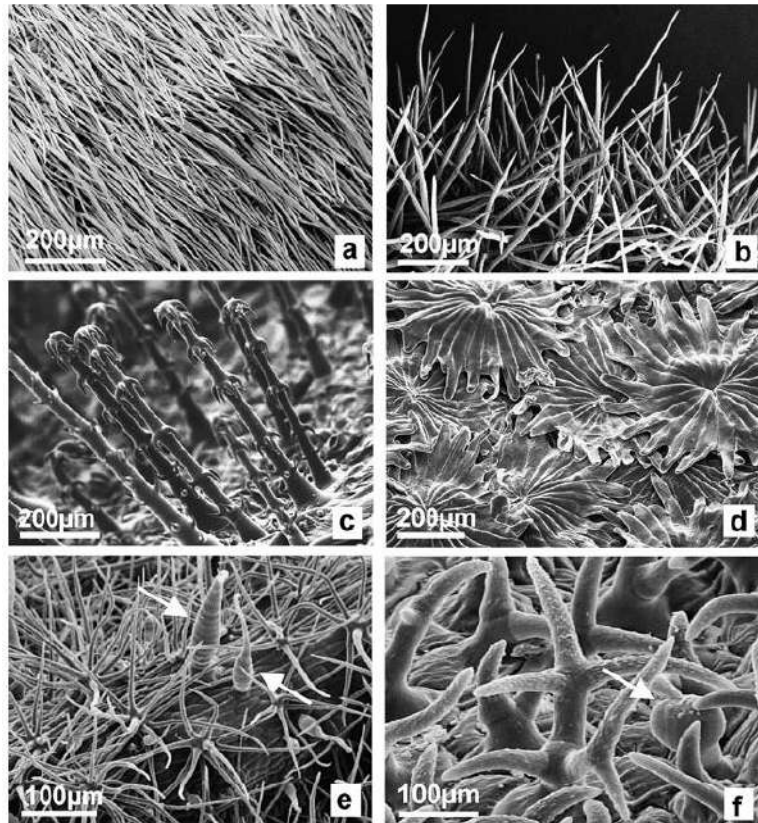
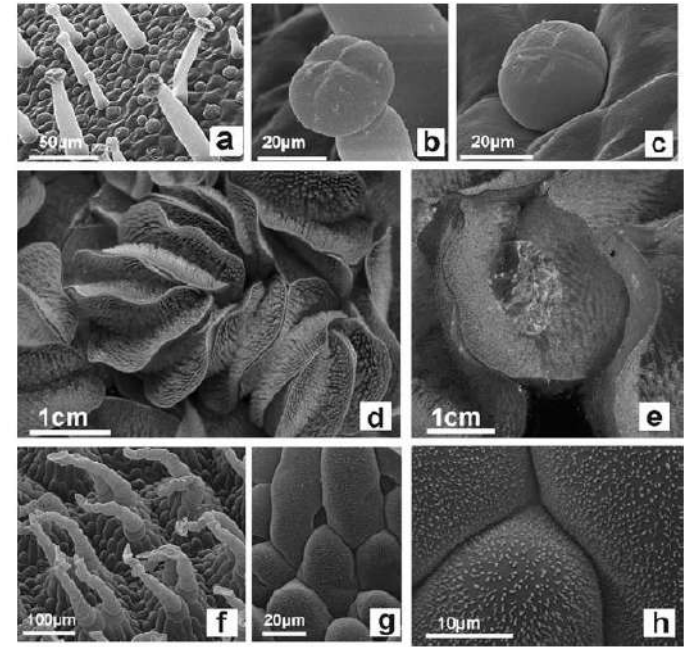
$$\theta > 160^\circ$$



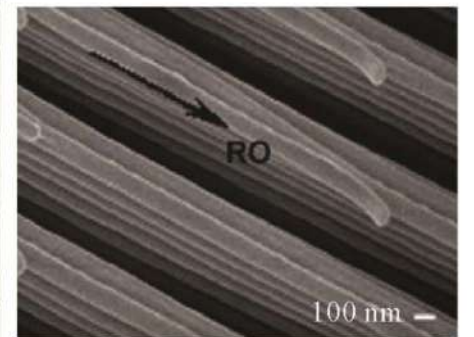
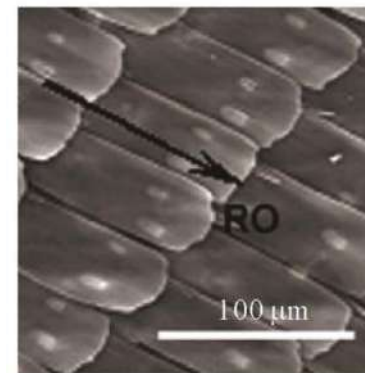
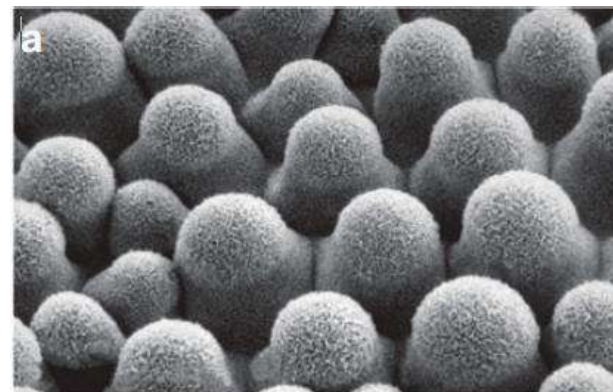
Superhydrophobicity in nature



Koch et al., *Soft Matt.* 2008



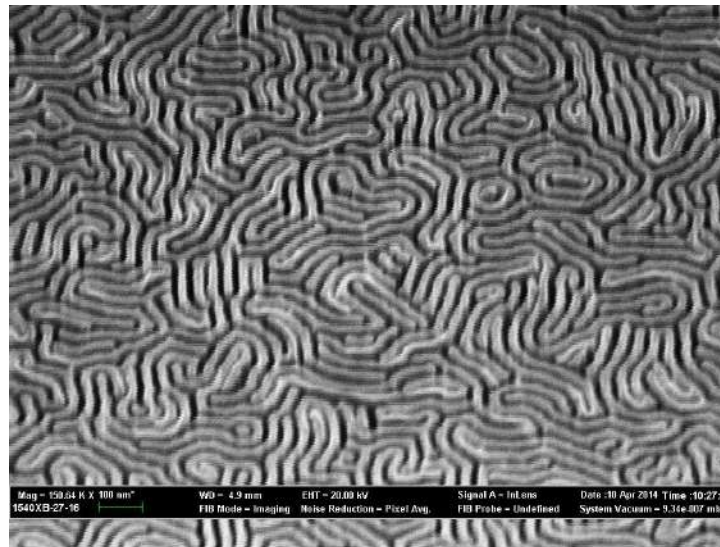
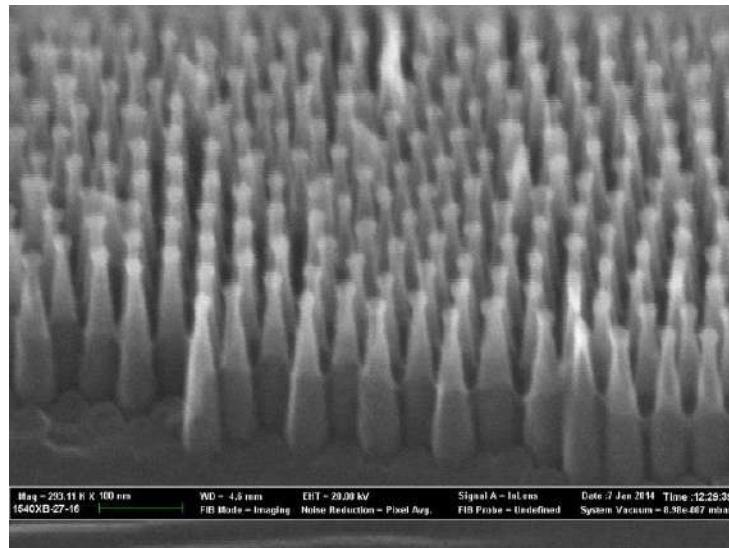
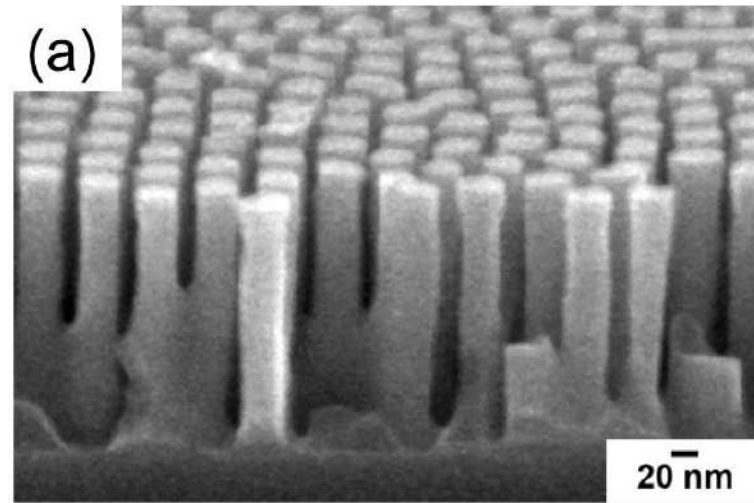
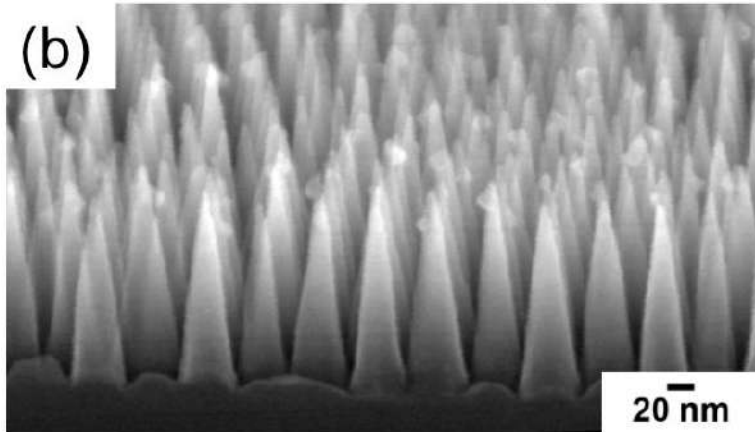
Micro + nano



50 μm

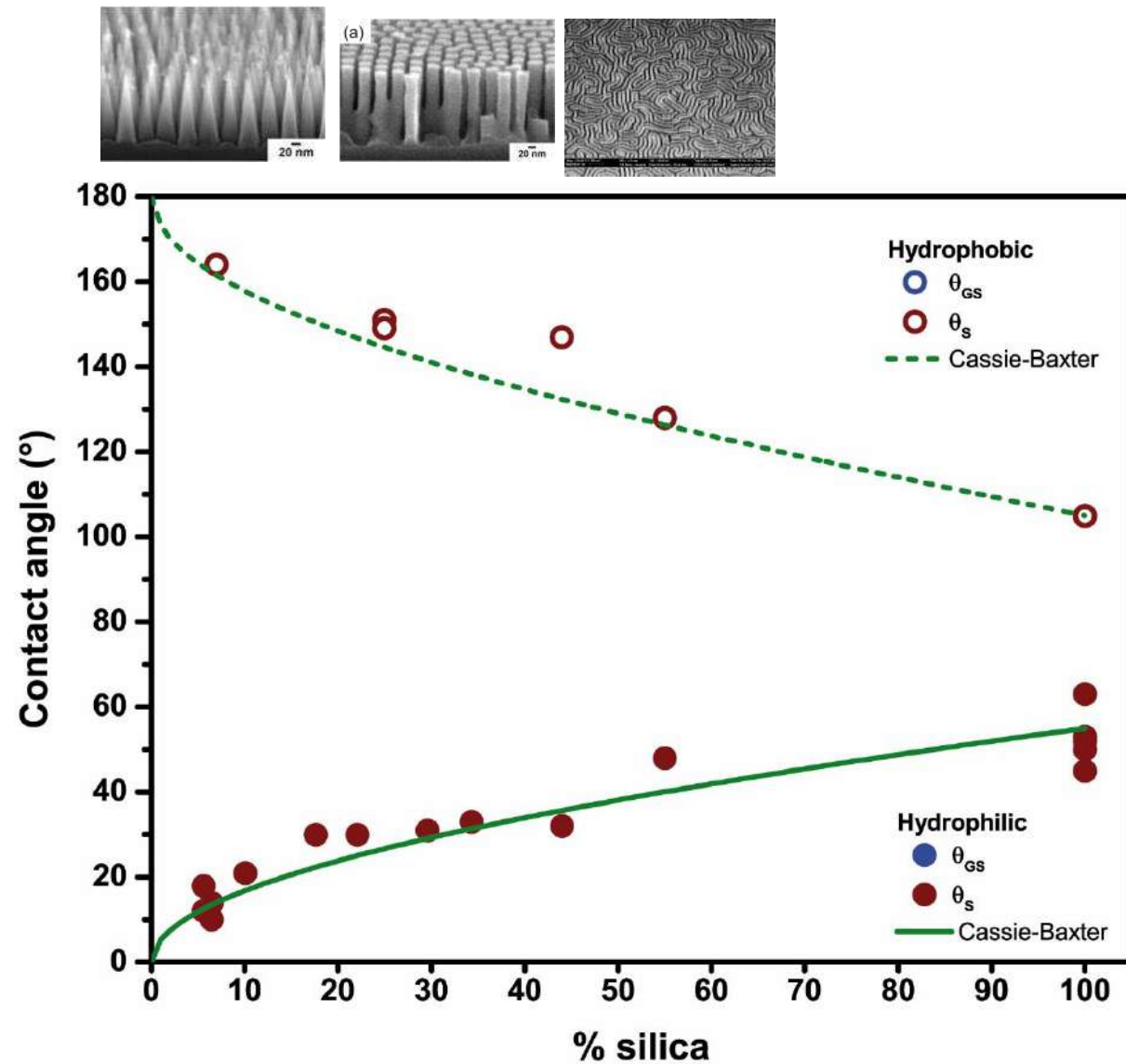
Artificial superhydro-phobic (philic) surfaces

Checco *et al.*, Adv. Mat. 2014

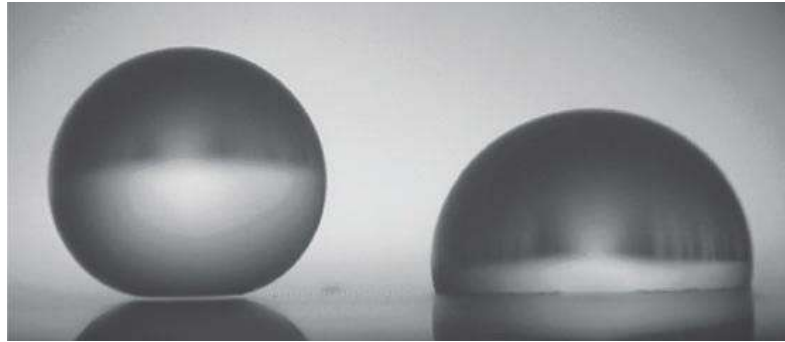


Superhydrophilicity - superhydrophobicity

Ondarçuhu *et al.*, Sci. Rep. 2016



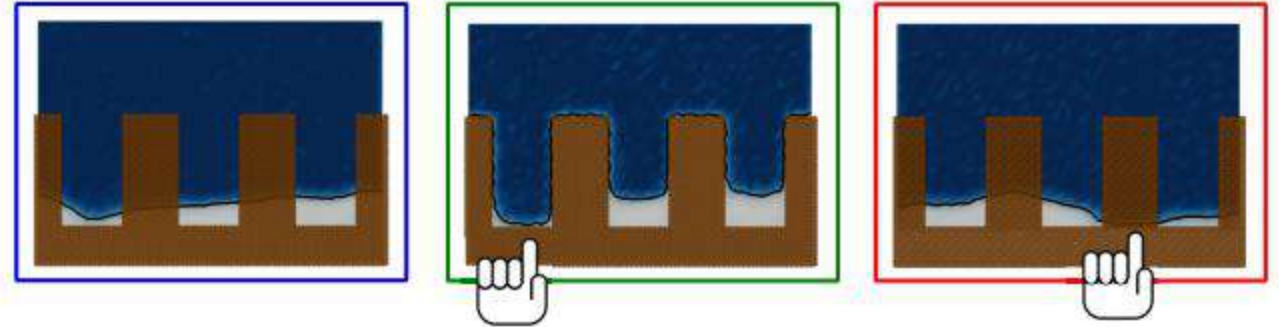
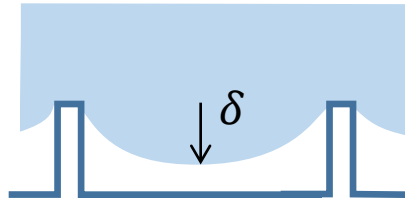
Metastability of Cassie state



Cassie

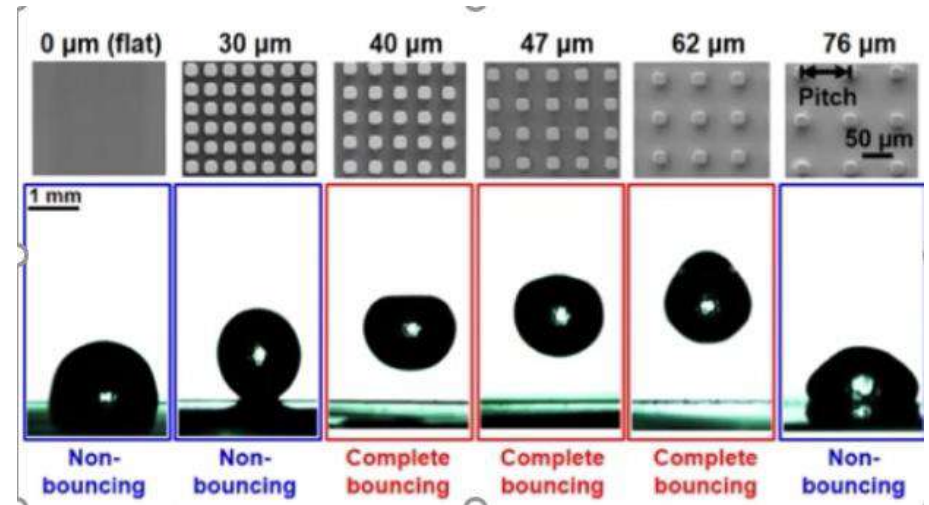
Wenzel

$$\Delta E \approx 2\pi b h / p^2 \gamma \cos \theta$$

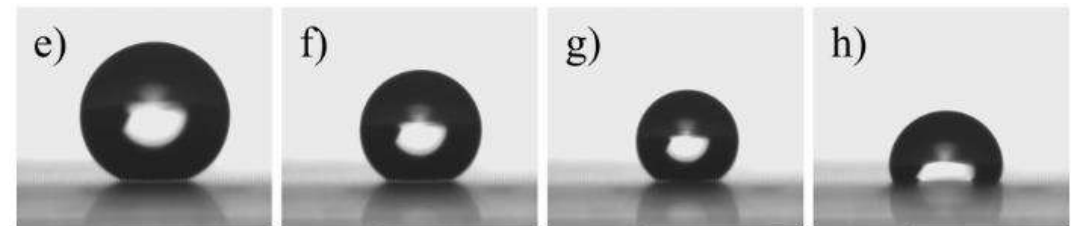


Amabili *et al.*, PRF 2017

Impact



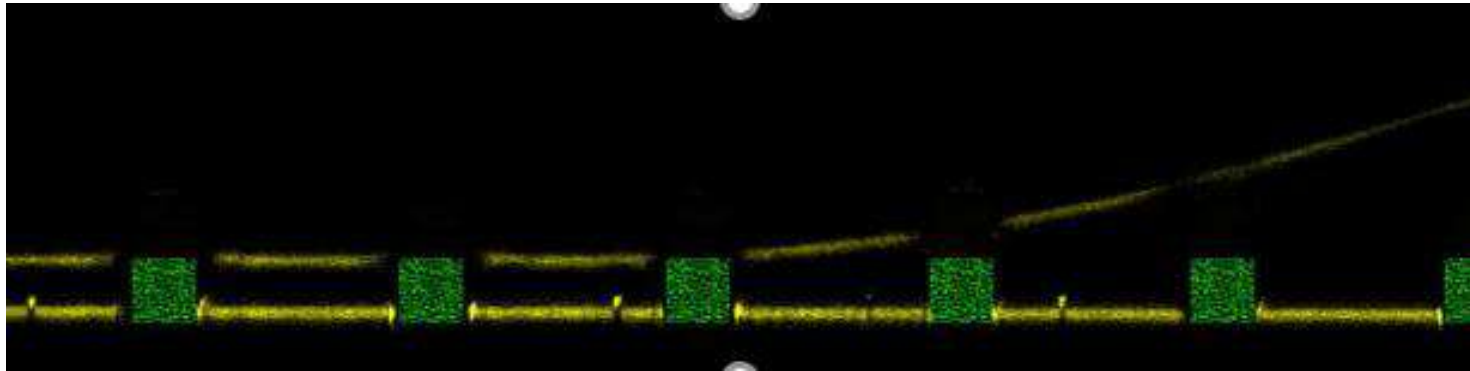
Evaporation



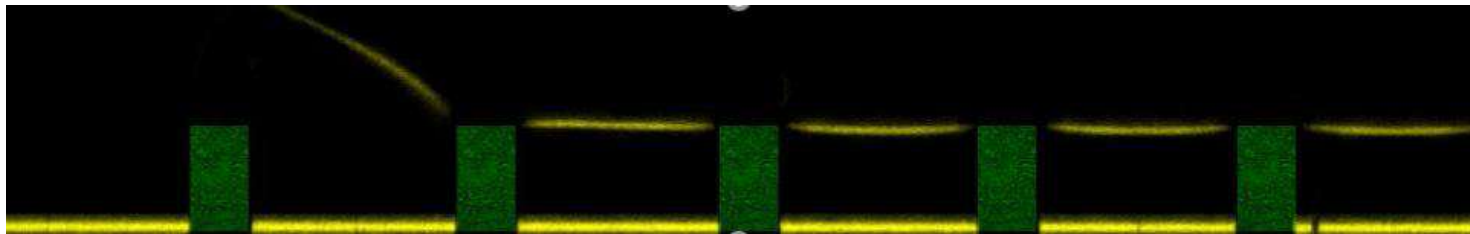
Mc Hale *et al.*, Langmuir 2004

At microscopic scale

Schellenberger *et al.*, Sci. Rep. 2016



$$\theta_{av} \sim 180^\circ$$



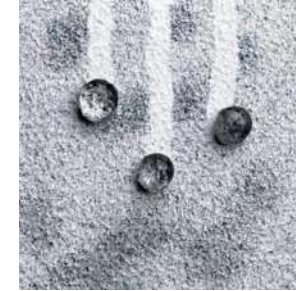
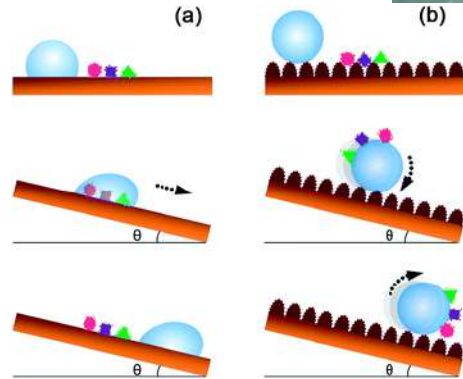
$$\theta_{rec} \ll 180^\circ$$

Applications

➤ Water repellency



➤ Self Cleaning



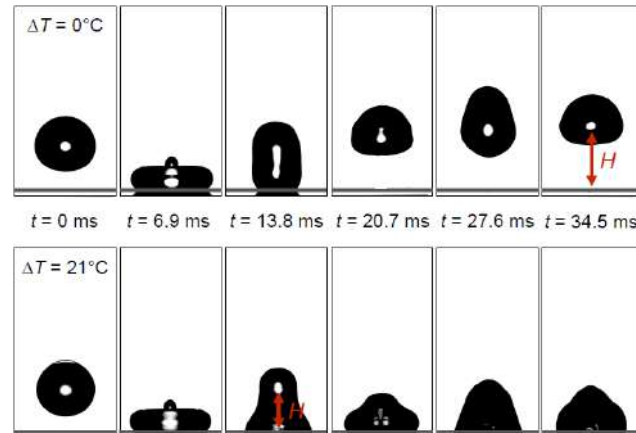
Nishimoto *et al.*, RSCAdv 2015

➤ Issues

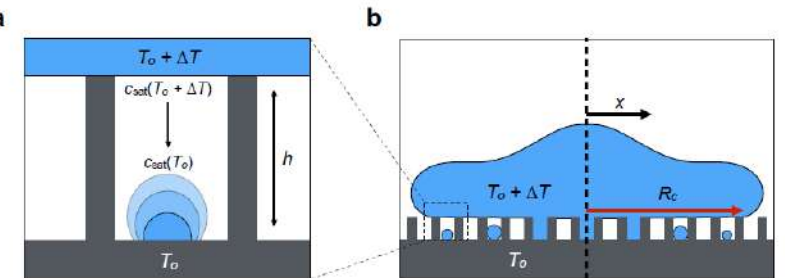
✓ Robustness

✓ Contamination

✓ Hot water

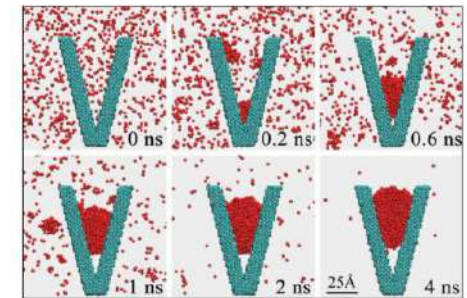
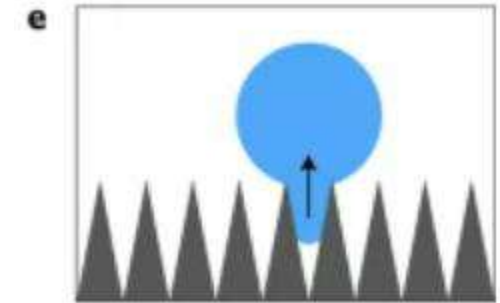
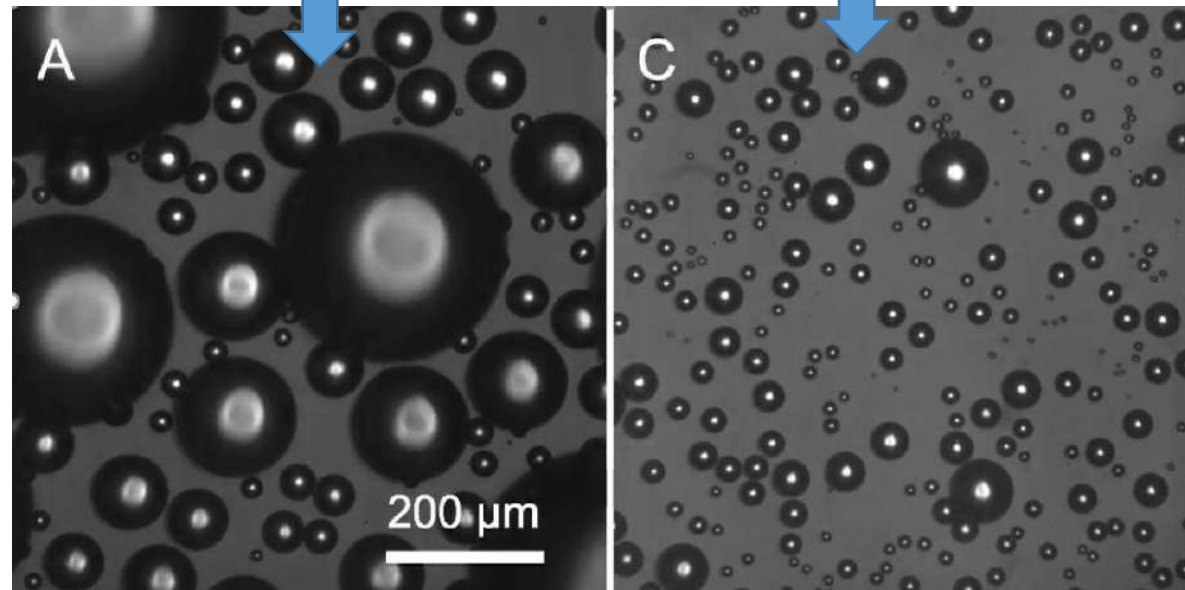
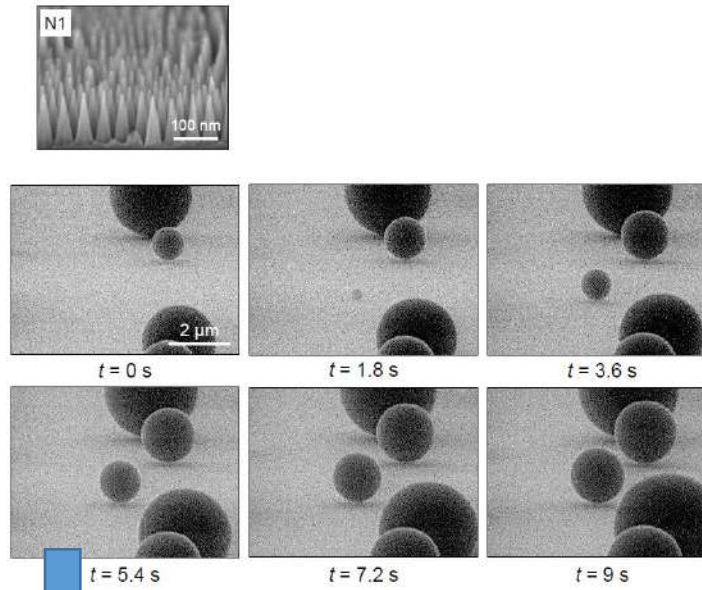
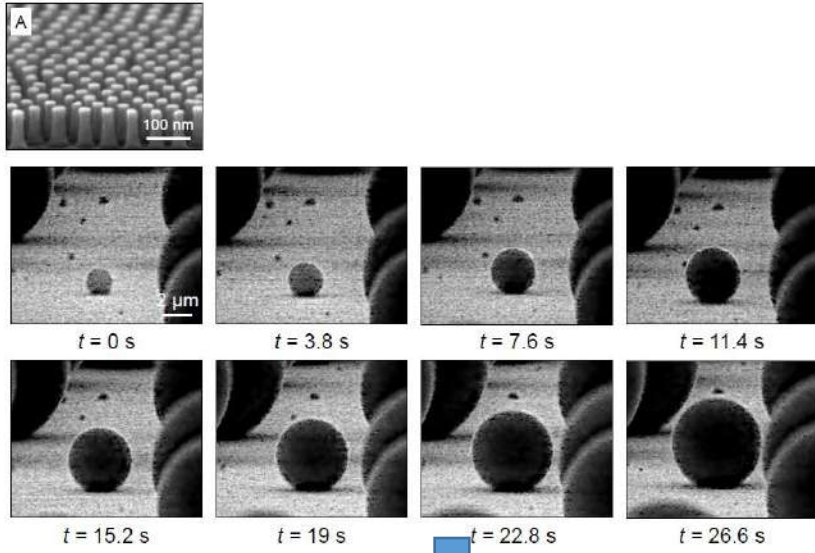


Mouterde *et al.*, Nature Comm. 2019



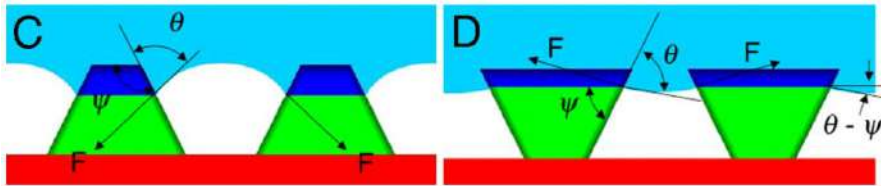
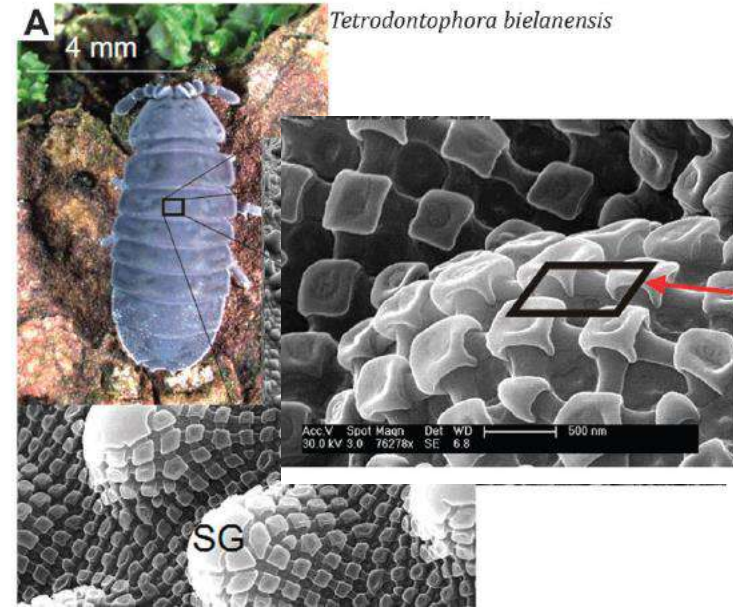
Anti-fogging

Lecointre *et al.*, *Nature Comm.* 2021

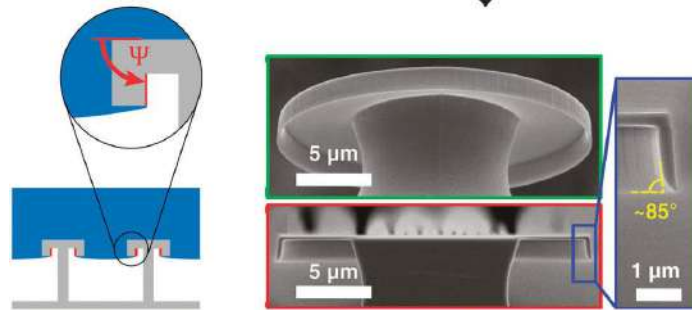


Omniphobic surfaces

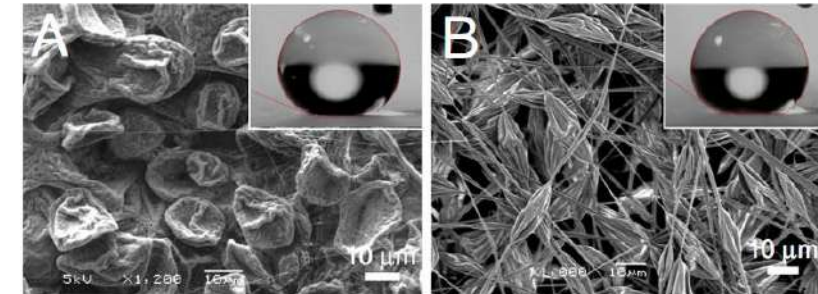
➤ Contact angle of oil on hydrophobic surface $< 90^\circ$



Tuteja et al., PNAS 2008



Hensel et al., Chem. Soc. Rev. 2016

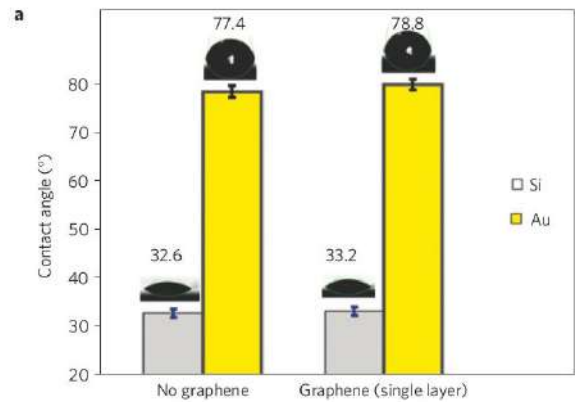
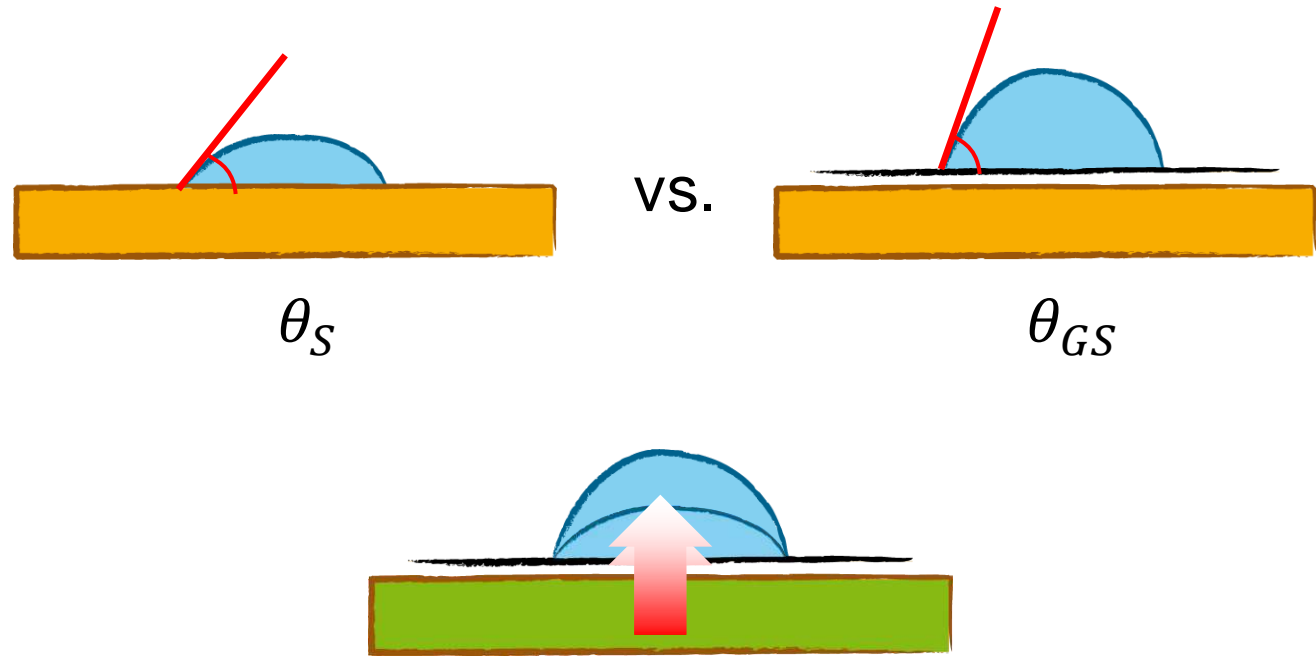
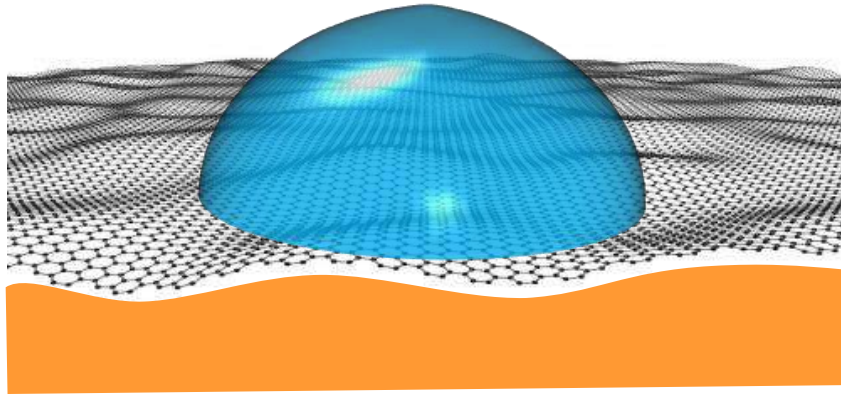


Tuteja et al., PNAS 2008

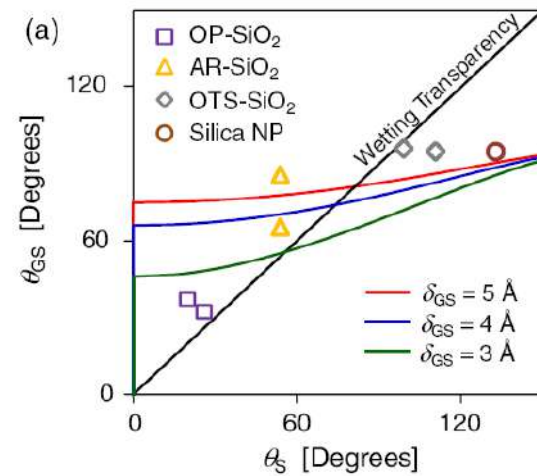
➤ Long range forces (vdW)

- ✓ Which interaction determines contact angle?
- ✓ How does it affect spreading regimes?
- ✓ How does it affect thin film stability?
- ✓ Shape of nanodroplets?

Wetting of graphene

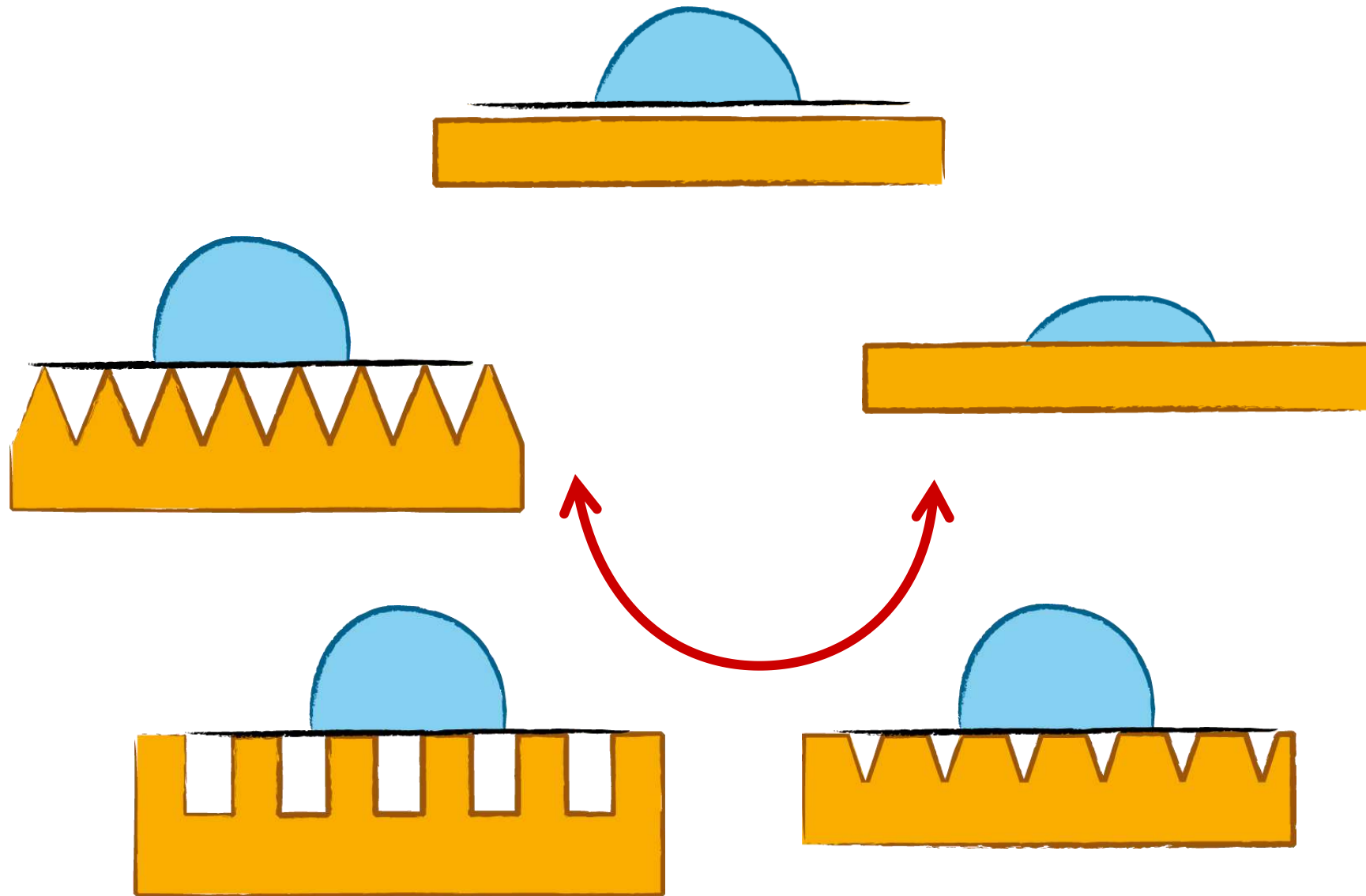


Rafee et al. Nature Mat. 2012

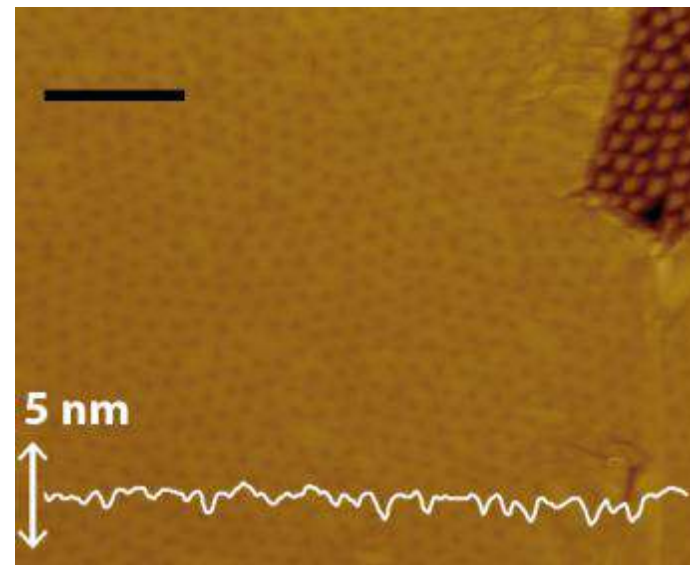
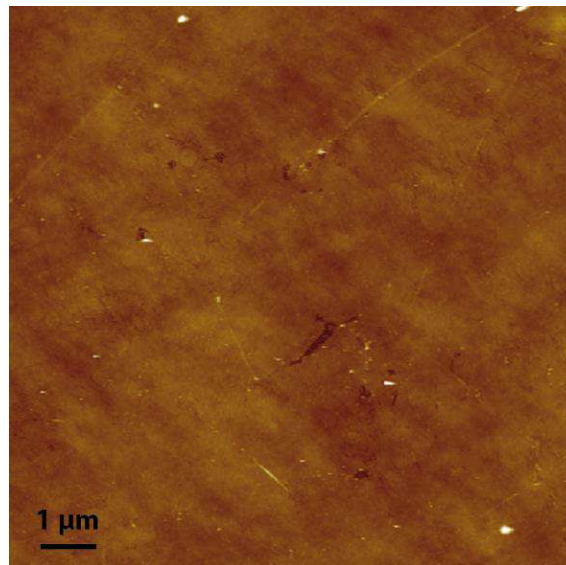
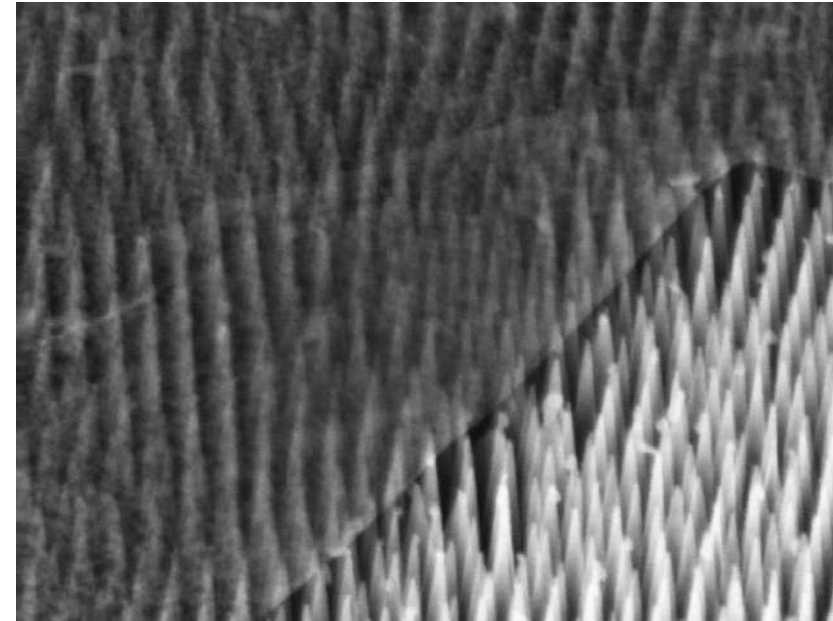
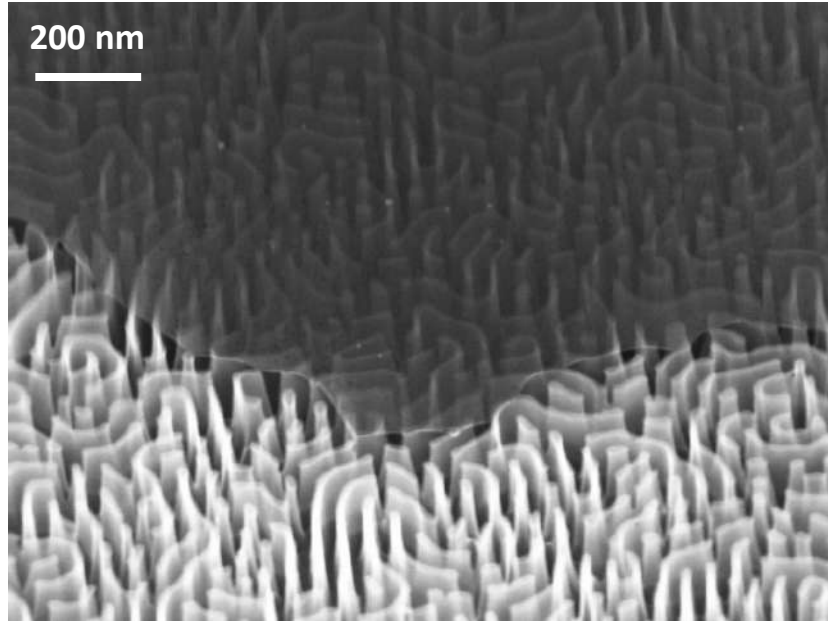


Shih et al. PRL 2012

Wetting of graphene

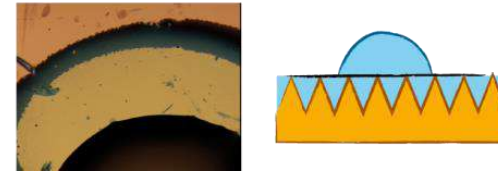
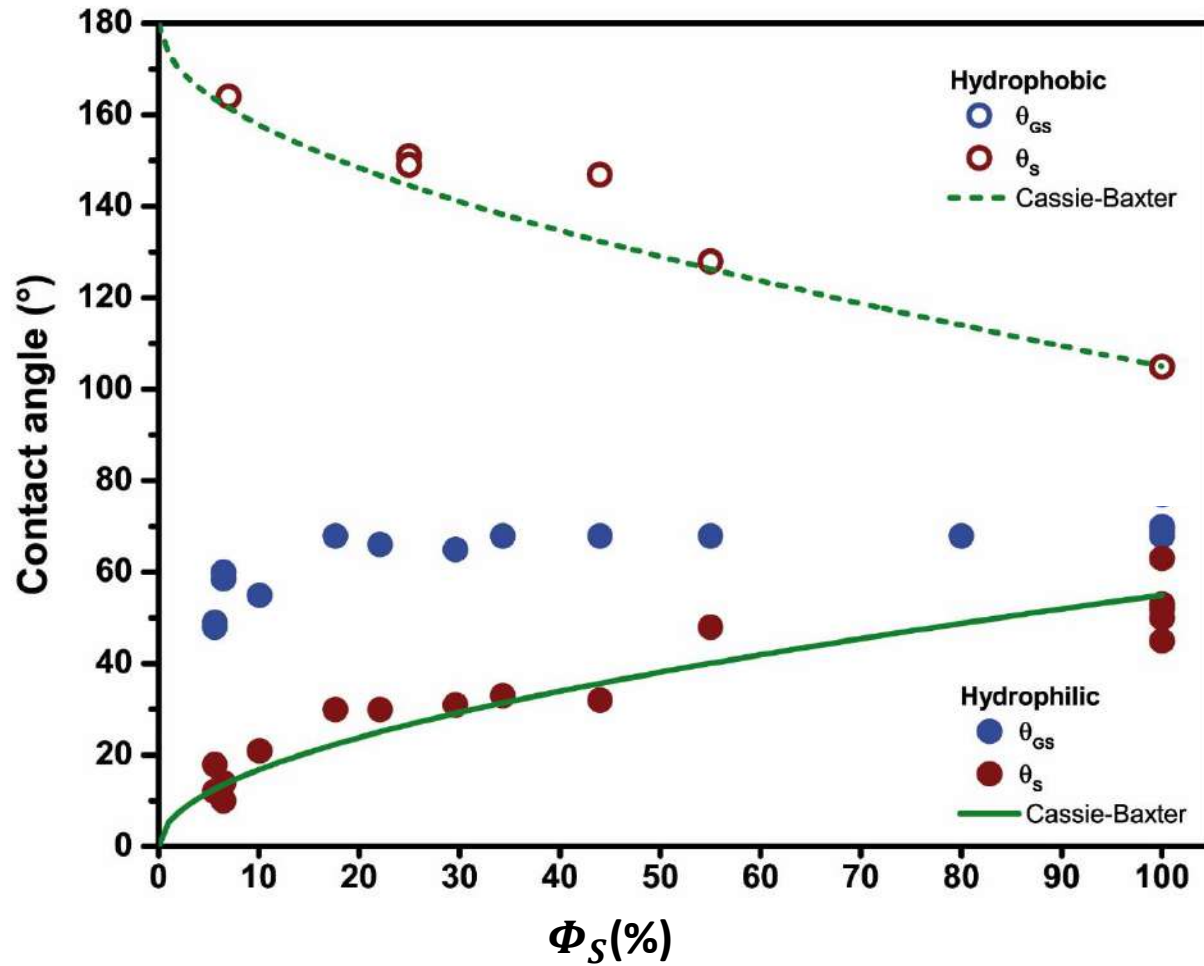


Wetting of graphene



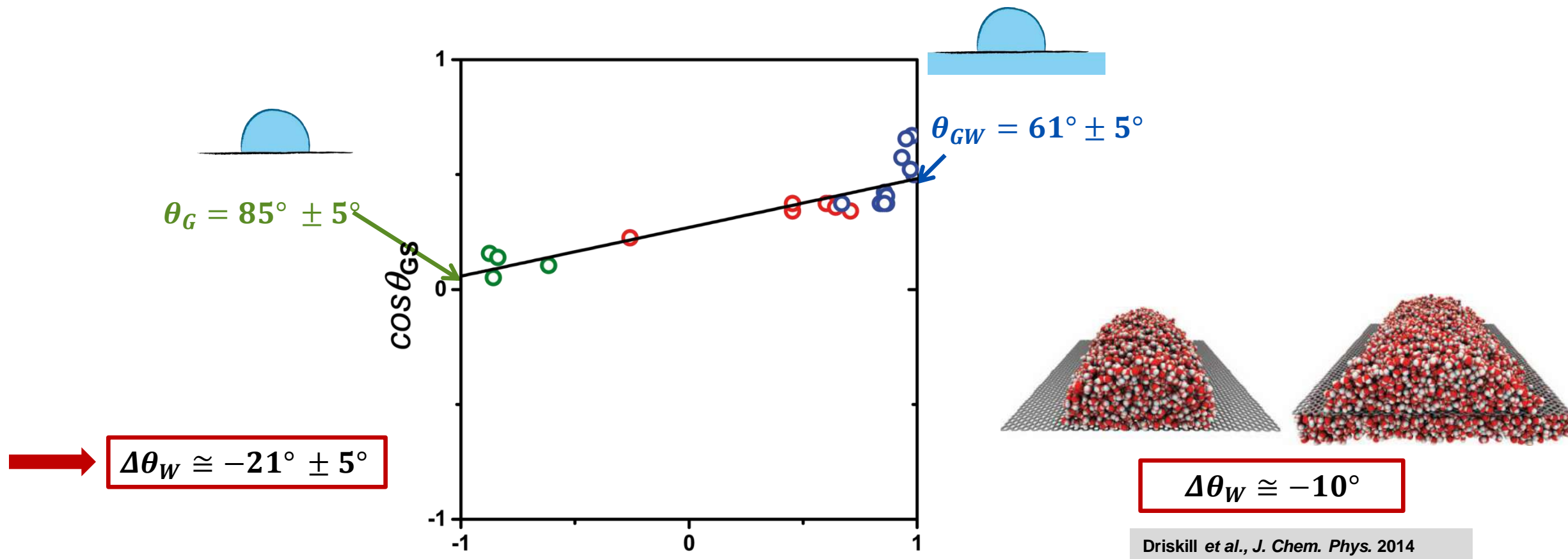
Rms roughness
0,4 nm

Wetting of graphene



➤ No wetting transparency

Wetting of graphene



Driskill et al., J. Chem. Phys. 2014

80 % of the interaction comes from the graphene
20 % from underlying substrate

$$W(h) = -\frac{A}{3\pi h^2}$$

Nano-wetting: effect of long range forces

➤ Effective interaction potential



$$E(h) = \gamma + \gamma_{SL} + W(h)$$

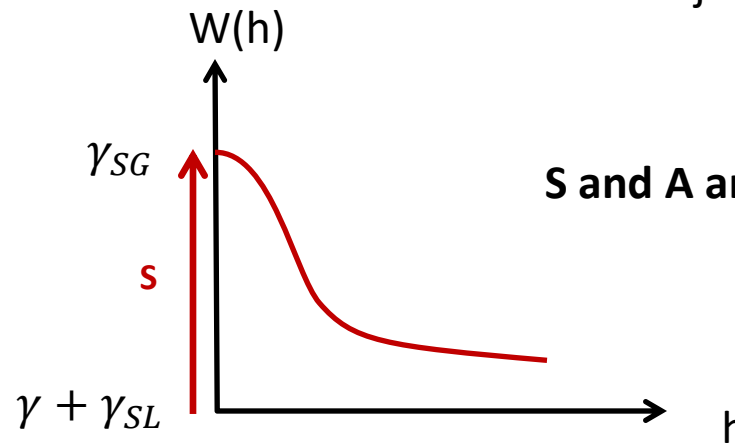
$$E(0) = \gamma_{SG} \quad \longrightarrow \quad W(0) = S$$

$$\mathcal{E}(h) = \mathcal{A}(\gamma + \gamma_{SL} + W(h))$$

At constant volume : $\frac{d\mathcal{A}}{\mathcal{A}} + \frac{dh}{h} = 0$

$$\frac{d\mathcal{E}}{d\mathcal{A}} = \gamma_{film} = \gamma + \gamma_{SL} + W(h) + h\Pi(h)$$

Brochard *et al.*, Langmuir 1991



S and A are rather independant

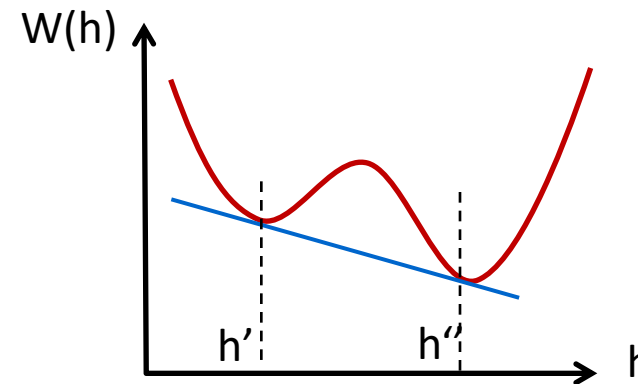
van der Waals interaction :

$$W(h) = \frac{A}{3\pi h^2}$$

Hamaker constant

$$A = \pi^2 k \bar{\alpha}_L (\bar{\alpha}_S - \bar{\alpha}_L)$$

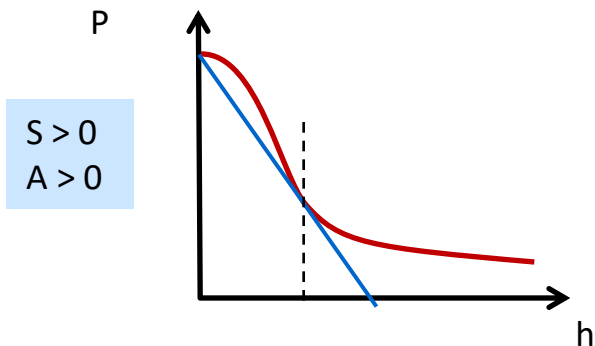
Disjoining pressure : $\Pi(h) = -\frac{dW}{dh} = \frac{A}{6\pi h^3}$



Nano-wetting: effect of long range forces

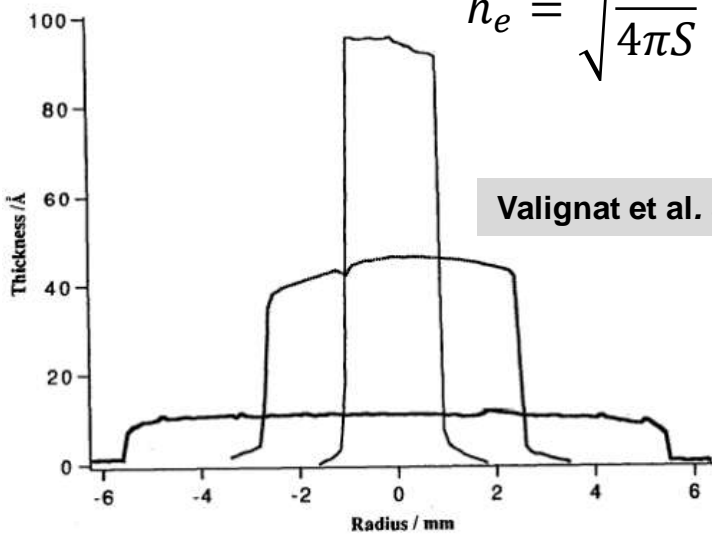


Total wetting

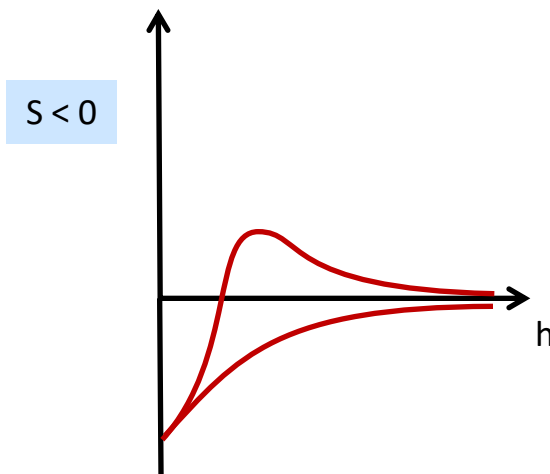


$$h_e = \sqrt{\frac{A}{4\pi S}}$$

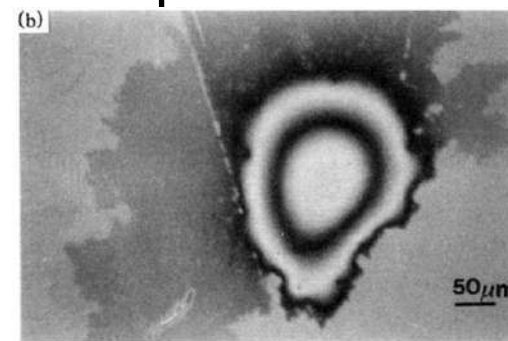
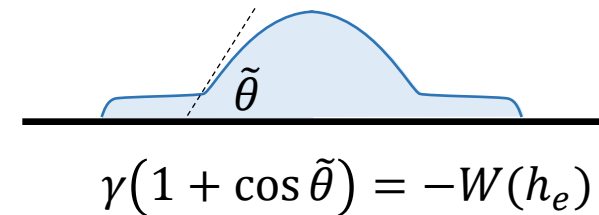
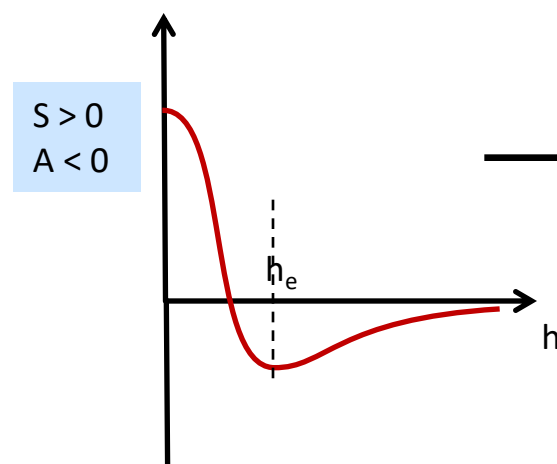
Valignat et al. *JCIS* 1994



Partial wetting



Pseudo-partial wetting



Silberzan, Léger, *PRL* 1991

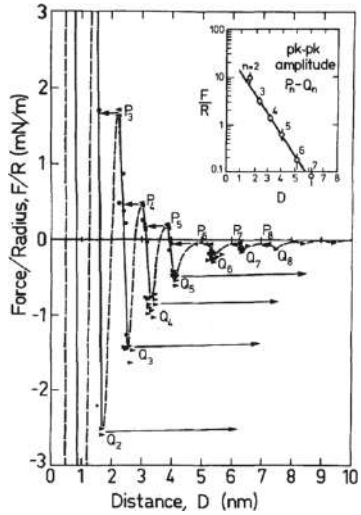
More frequent in L/L

Du et al.,
Chem Eng. Process. 2013

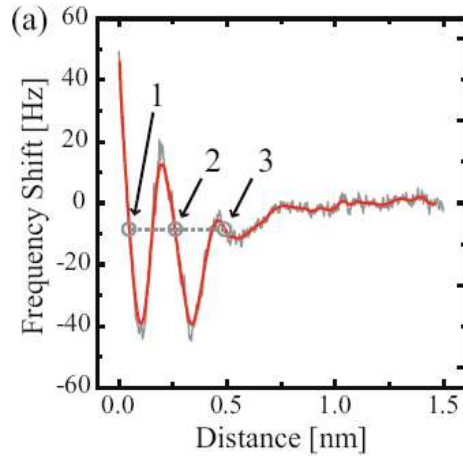
➤ Wetting criterion depends on :
- sign of spreading coefficient
- sign of Hamaker constant

Nano-wetting: effect of long range forces

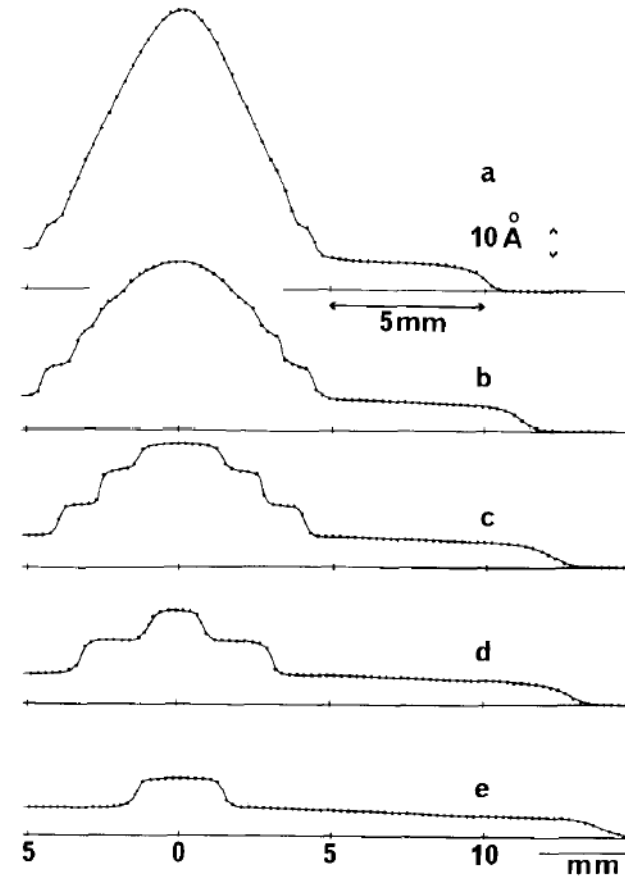
➤ Structural forces



Israelachvili (1968-1970)



T. Ando



Heslot *et al.*, Nature 1989

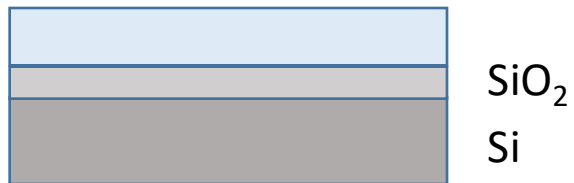
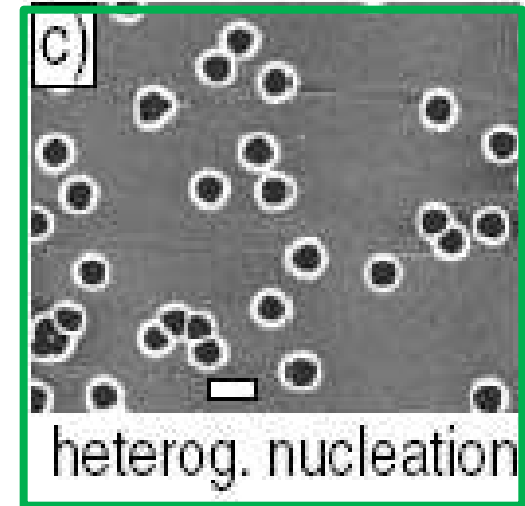
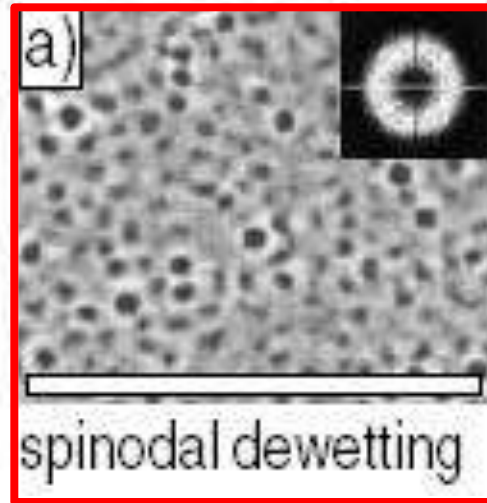
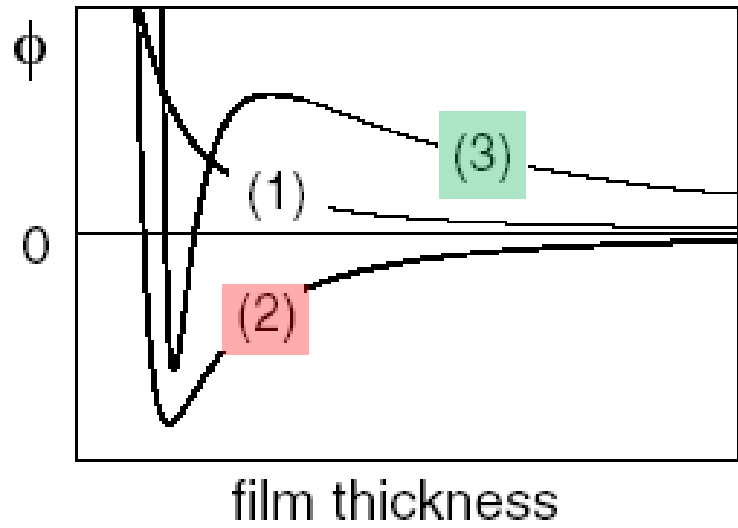
tetrakis(2-ethylhexoxy)-silane (TK)

➤ Wetting criterion depends on : - sign of spreading coefficient
- sign of Hamaker constant

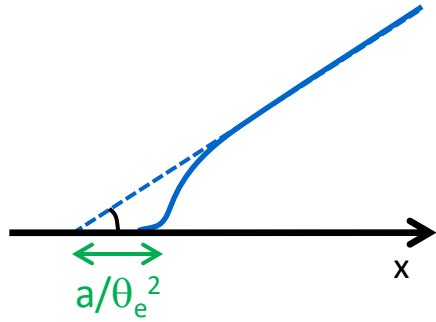
Effect on wetting films

➤ Dewetting

$\Phi''(h) < 0$ unstable



Influence on drops

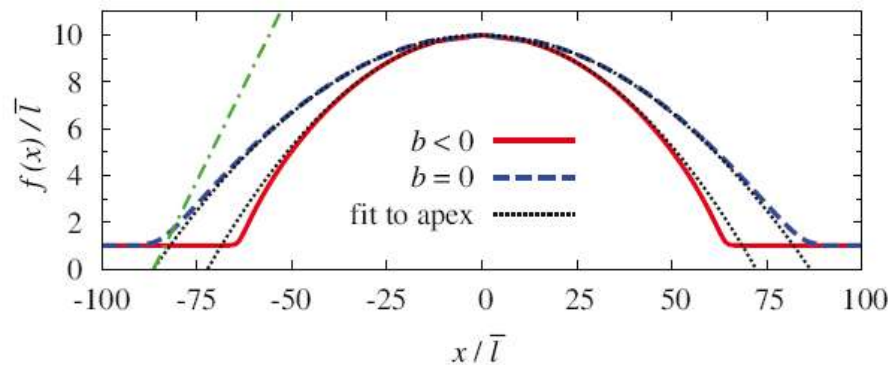


$$\gamma \frac{d^2 h}{dx^2} + \Pi(h) = 0$$

$$\frac{1}{2} \gamma \left(\frac{dh}{dx} \right)^2 - W(h) = \frac{1}{2} \gamma \theta_e^2$$

$$h^2 = \theta_e^2 x^2 \pm \frac{a^2}{\theta_e^2} \quad \text{with} \quad a^2 = \frac{|A|}{6\pi\gamma}$$

$$\mathcal{H}[f] = \int dx \int dy \left[\frac{1}{2} \sigma_{lv} (\nabla f)^2 + \omega(f) \right]$$



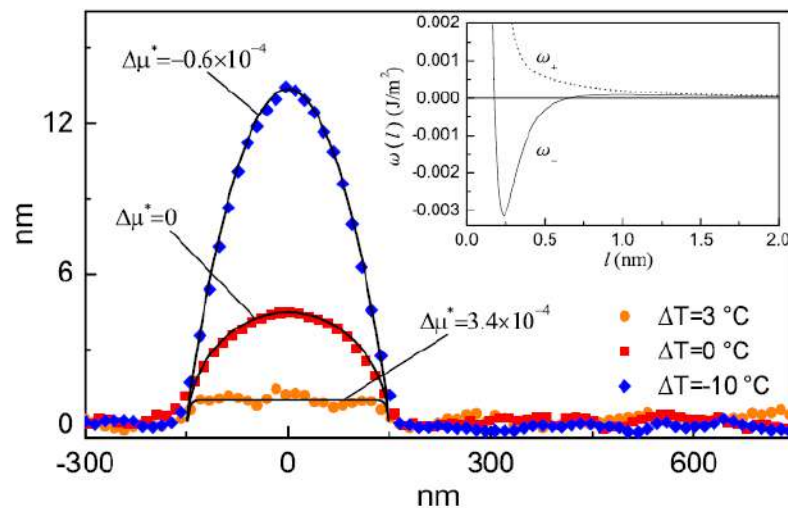
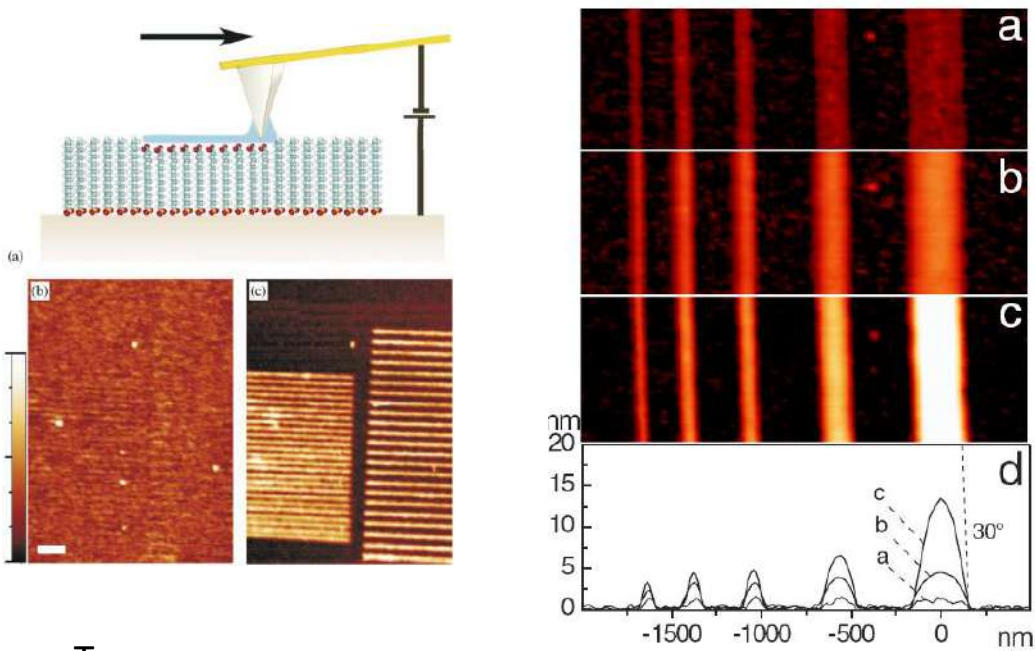
Getta and Dietrich. Phys. Rev. E 1998

Influence on drops

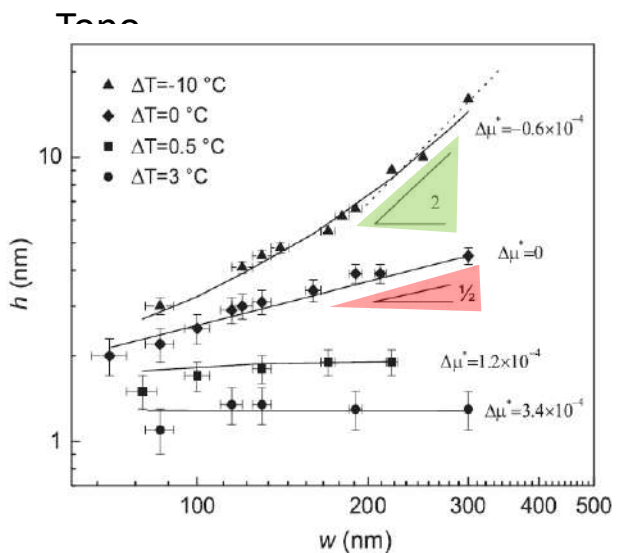
Condensation on hydrophilic nanostripes

Checco *et al.*,

Phys. Rev. Lett. 2003
Phys. Rev. E 2008



$$\gamma \frac{l''(x)}{\sqrt{(1+l'^2(x))^3}} = \rho \Delta\mu + \frac{\partial \omega_{CS}(x, l(x))}{\partial l}$$



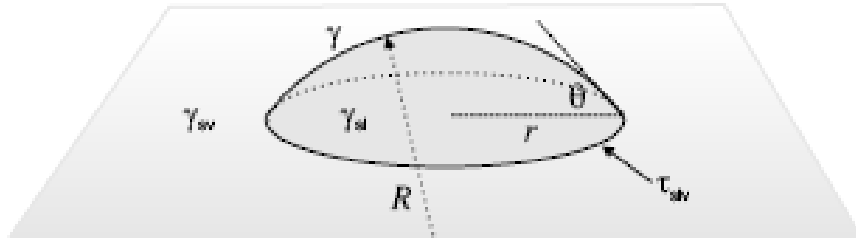
- $\Delta\mu = 0$ (saturation)

$$\gamma \frac{d^2 h}{dx^2} = \frac{-2A}{h^3} \Rightarrow \frac{h}{w^2} \propto \frac{1}{h^3} \Rightarrow h \propto w^{1/2}$$

- Spherical cap

$$h \propto w^2$$

➤ Phenomenological description



$$dE = \gamma dA_{LV} + \gamma_{SL} dA_{SL} + \gamma_{LV} dA_{LV} + \tau dL$$

$$= \left(\gamma \cos \theta + \gamma_{SL} - \gamma_{LV} + \frac{\tau}{r} \right) dA_{LV}$$

$$\cos \theta = \cos \theta_e - \frac{\tau}{\gamma r}$$

$$\gamma \approx \frac{kT}{a^2} = 10^{-2} \text{ N/m}$$

$$\tau \approx \frac{kT}{a} = 10^{-12} \text{ N}$$



$$\frac{\tau}{\gamma} \approx a$$

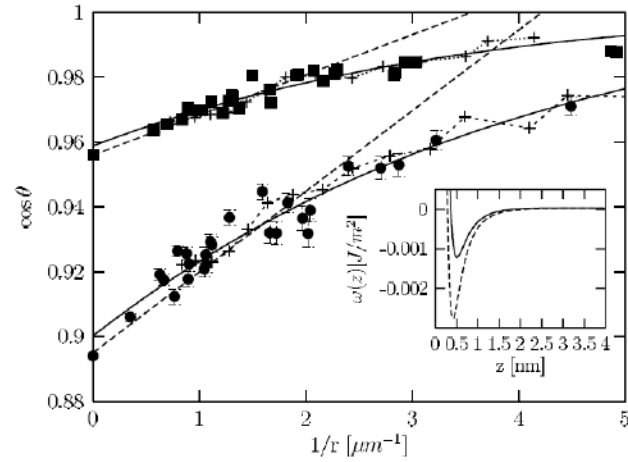
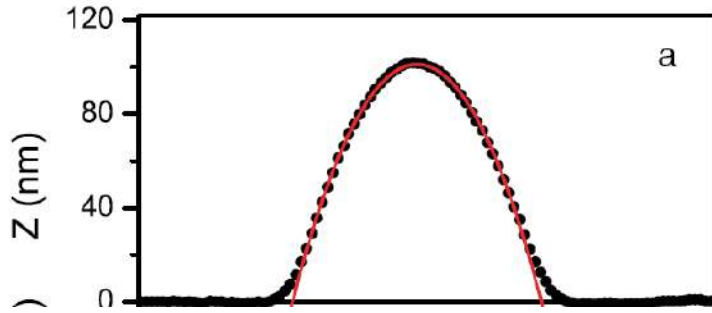
$$\tau > 0 \text{ or } \tau < 0$$

➤ Theoretical description

Multiples contributions

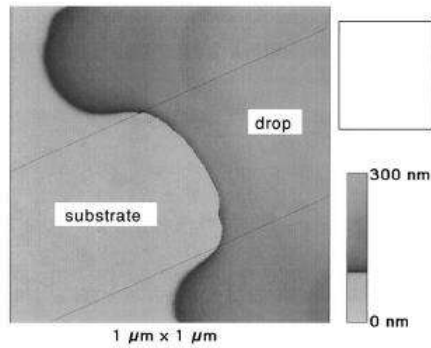
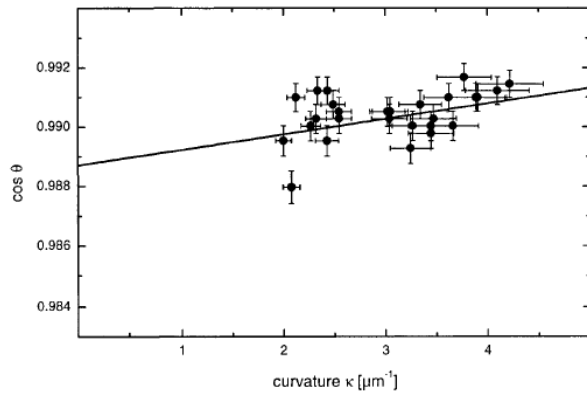
$$\tau_{vs} = \sqrt{2 \sigma_V} \int_2^{\infty} \left[\sqrt{\omega(\ell) - \omega(\tilde{\ell})} - \sqrt{-\omega(\tilde{\ell})} \right] d\ell, \quad + \text{Tolman length} + \text{line rigidity} \dots$$

Line tension: experimental determination



$$\tau = -500 \text{ pN}$$

Checco et al Phys. Rev.Lett. 2001

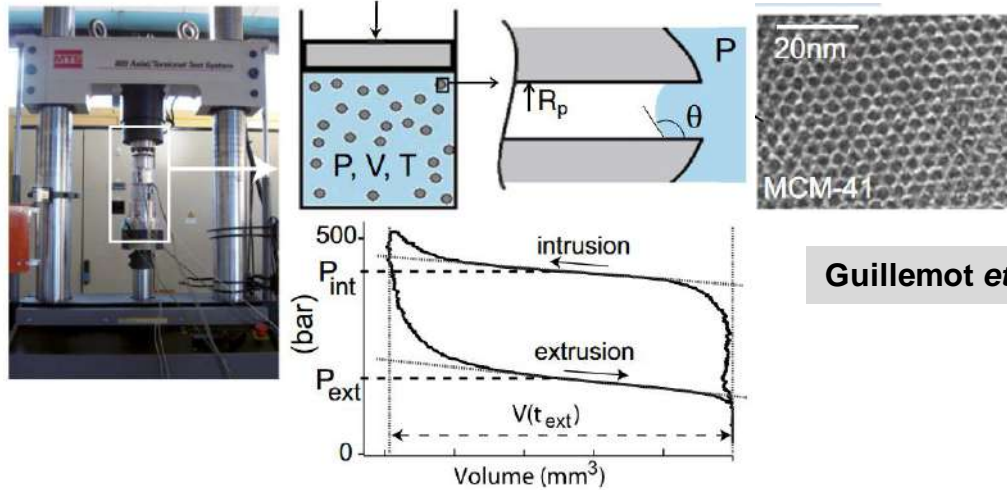


$$\tau = -300 \text{ pN}$$

Pompe, Herminghaus Phys. Rev.Lett. 2000

Line tension: experimental determination

➤ Nucleation of vapor bubble



Guillemot *et al.* PNAS 2012

Critical nucleus energy :

$$\Delta\Omega_c \simeq P_L K_1(\theta) R_p^3 + \gamma_{lv} K_2(\theta) R_p^2 + \tau K_3(\theta) R_p$$

$$P_{ext}^o = \frac{k_B T}{V_c} \ln \frac{L v t_0}{b} - \frac{K_2(\theta) \gamma_{lv}}{K_1(\theta) R_p} - \frac{\tau K_3(\theta)}{R_p^2 K_1(\theta)}$$

$$\tau = -23 \text{ to } -35 \text{ pN}$$

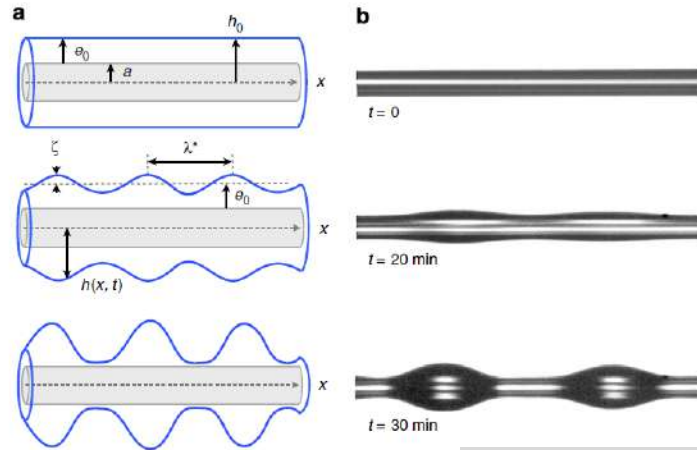
$$P_{ext} = -200 \text{ bars} \rightarrow 80 \text{ bars}$$



Huge impact on bubble nucleation (energy barrier)

Wetting of fibers

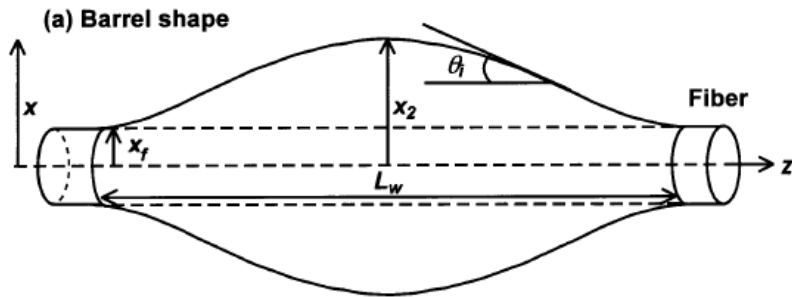
➤ Rayleigh-Plateau instability



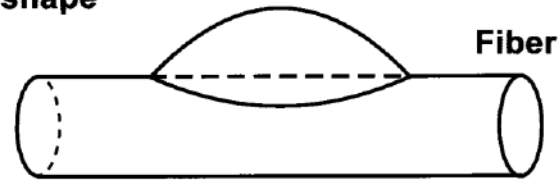
$$\lambda^* = 2\sqrt{2}a$$

Haefner *Nature Comm.* 2015

➤ Clampshell/barrel shapes



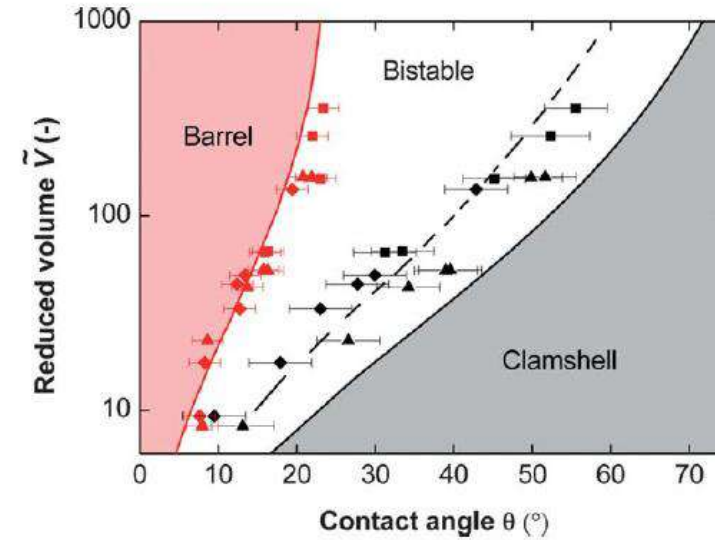
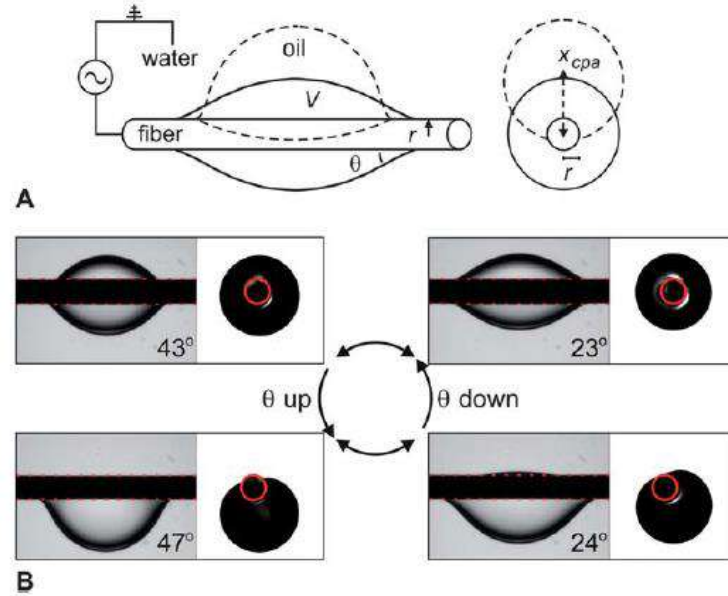
(b) Clam-shell shape



Mc Hale et al. *Coll. Surf. A.* 2002

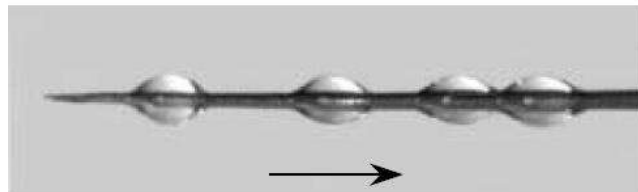
Wetting of fibers

➤ Barrell-clamshell transition

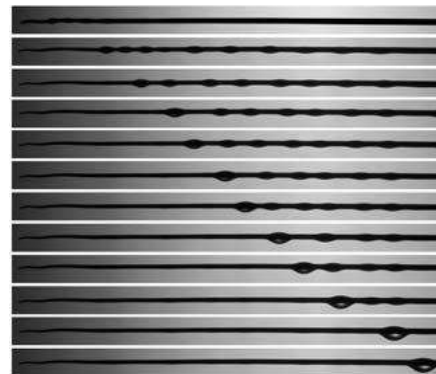


Eral et al. *Soft. Matt* 2011

➤ Tapered fibers

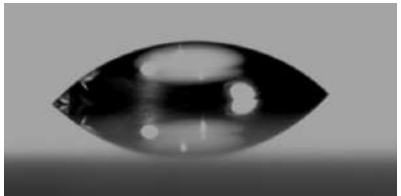


Lorenceanu et al. *JFM* 2004

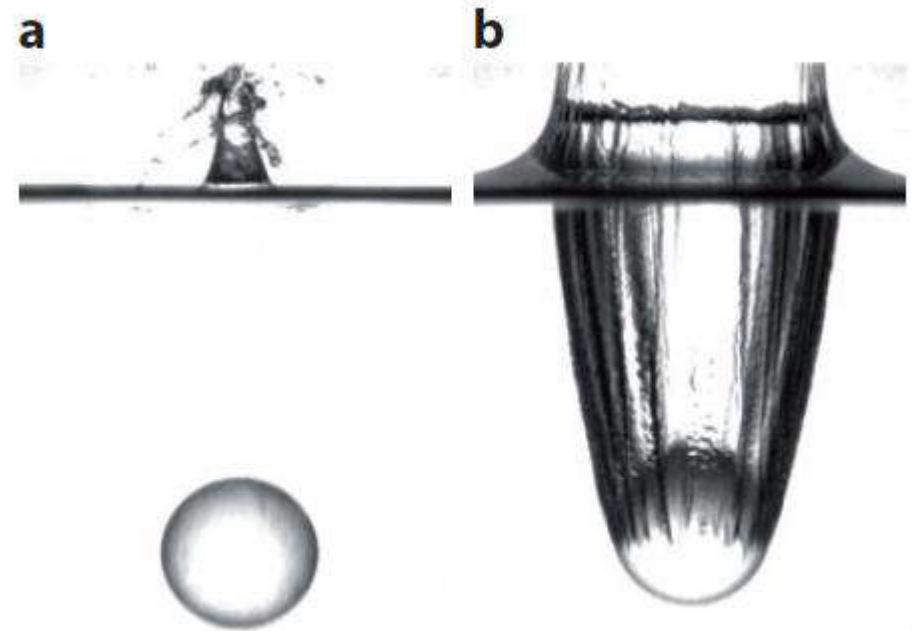


Coalescence filters, water collection, fog harvesting...

Wetting dynamics



=



Duez et al. *Nature Phys.* 2007

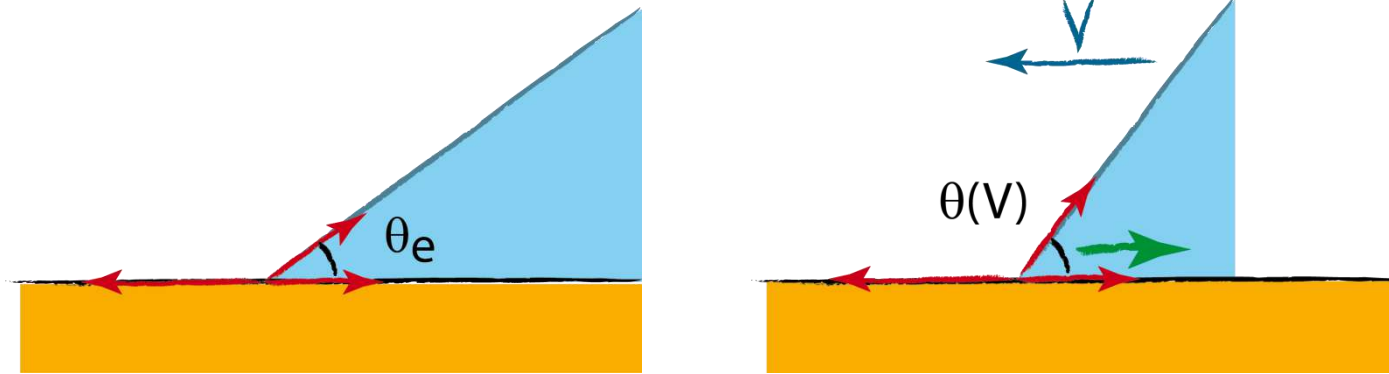
➤ Dynamic contact angle

- ✓ Dissipation at the contact line
- ✓ Dynamic instabilities

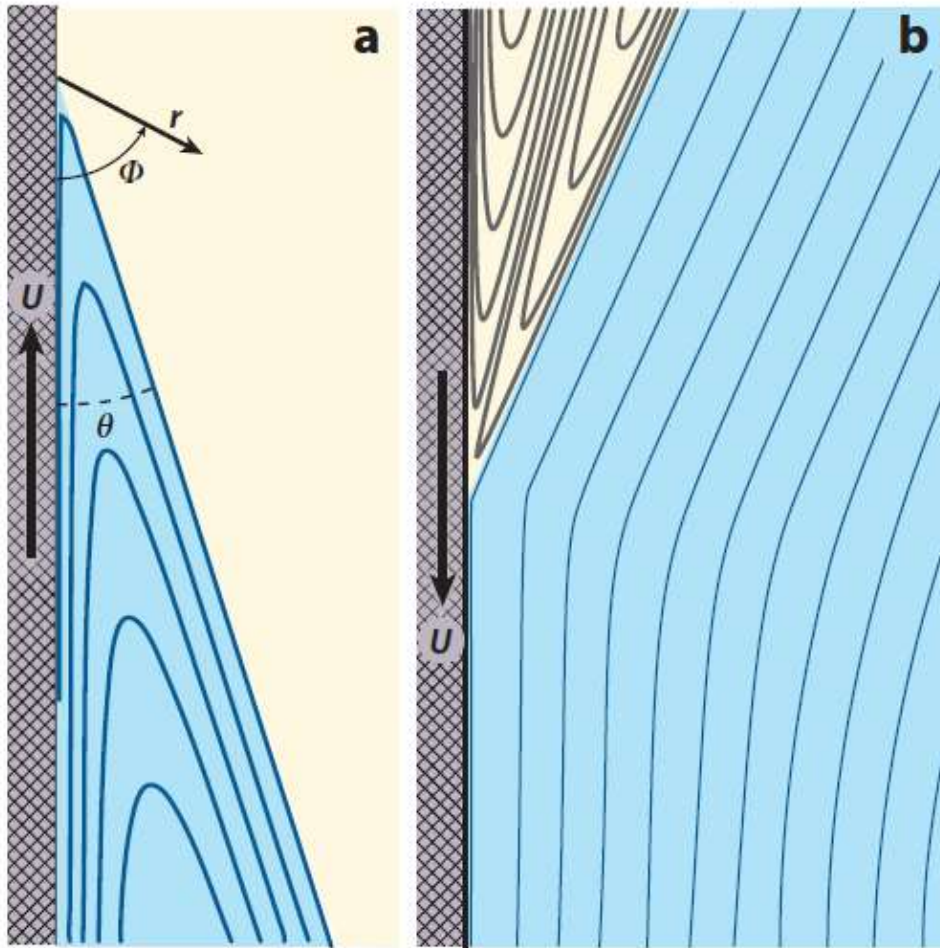
➤ Spreading regimes

- ✓ Droplet spreading
- ✓ Film dewetting

➤ Dynamic contact angle



Wetting dynamics

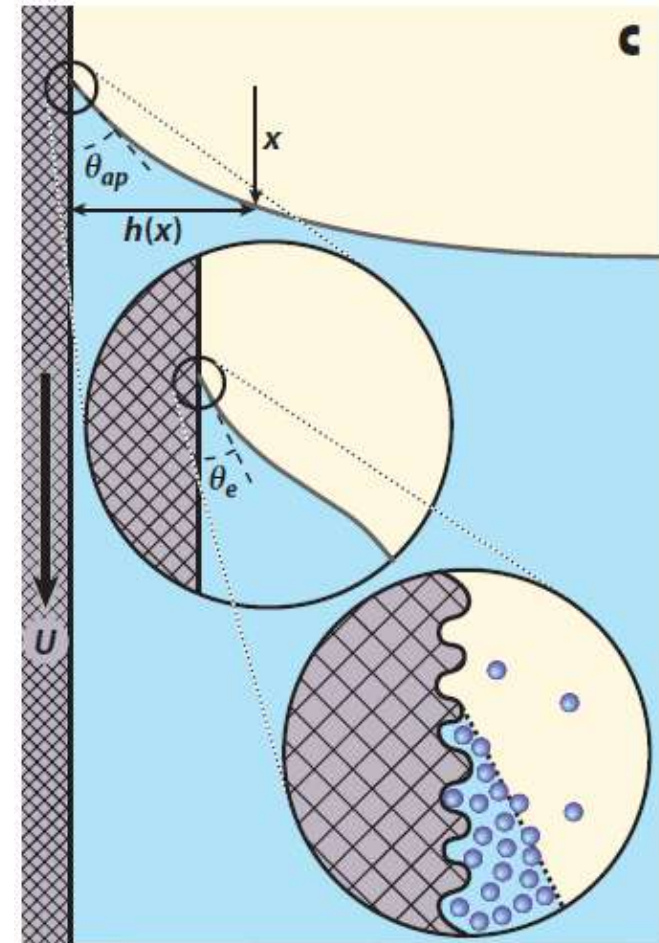


Huh, Scriven *JCIS* 1971

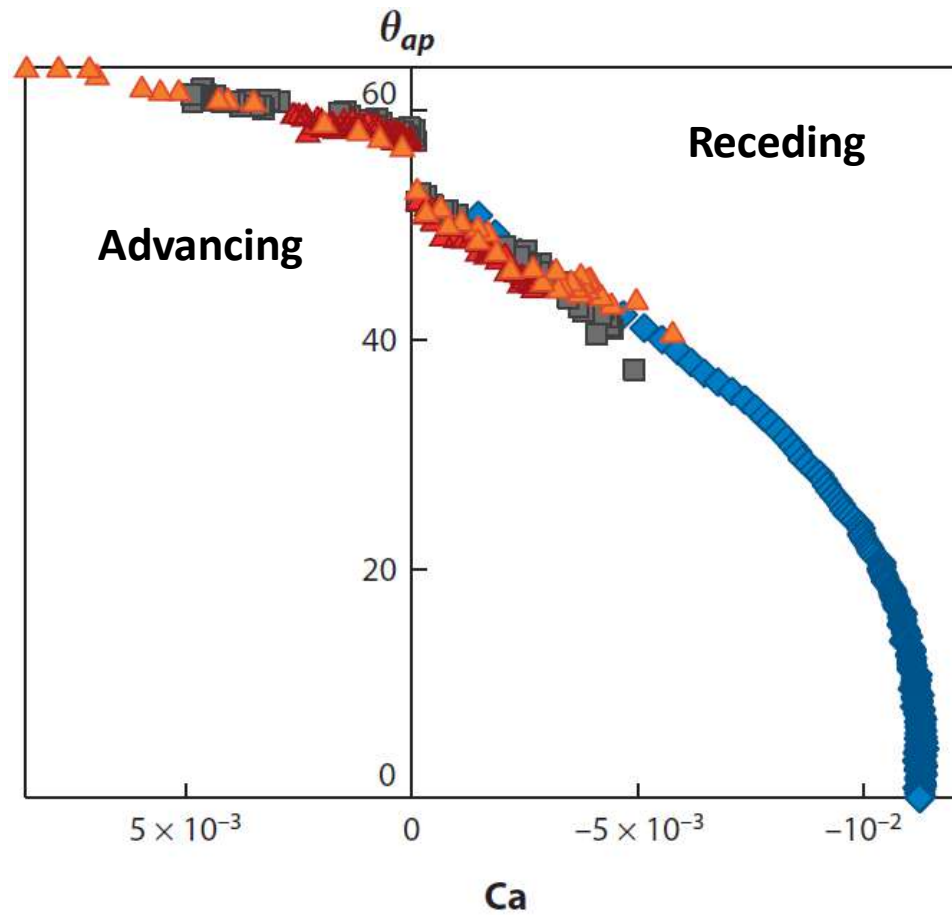
$$\sigma_{vis} \approx \eta \frac{U}{r}$$

diverges at CL

**Three regions:
macroscopic,
mesoscopic,
Molecular**



Wetting dynamics: macroscopic scale

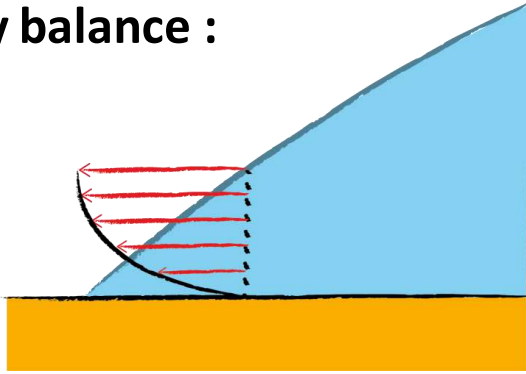


Capillary number

$$Ca = \frac{\eta U}{\gamma}$$

Wetting dynamics: mesoscopic scale

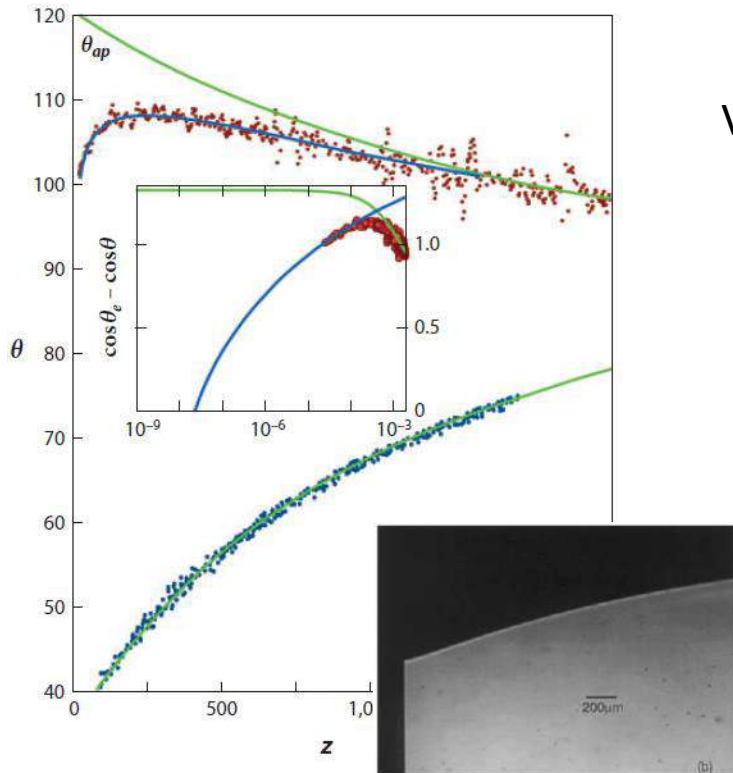
Visco-capillary balance :



Viscous stress vs. Capillary pressure

$$\frac{\partial P_{vis}}{\partial x} = \gamma(h_{xx})_x$$

$$h_{xxx} = -\frac{\pm 3Ca}{h^2}$$



Voinov solution: $h_x \approx \theta(x) \cong \left[9Ca \cdot \ln\left(\frac{x}{l}\right) \right]^{1/3}$

Voinov *Fluid Dyn.* (1976)

Matching :
Cox-Voinov equation:

$$\theta_{ap}^3 = \theta_e^3 \pm 9Ca \cdot \ln\left(\frac{L}{l}\right)$$

Cox *J. Fluid. Mech* (1986)

Ramé *et al. JCIS* (1971)

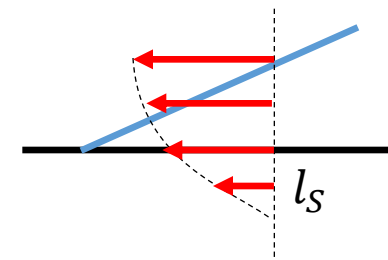
Wetting dynamics: microscopic scale

➤ Regularization of contact line singularity

mechanism	reference
mesoscopic precursor film	Hervet and de Gennes (1984)
molecular film	Eres <i>et al.</i> (2000)
Navier slip	Huh and Scriven (1971)
nonlinear slip	Thompson and Troian (1997)
surface roughness	Hocking (1976)
shear thinning	Weidner and Schwartz (1993)
evaporation/condensation	Wayner (1993)
diffuse interface	Seppecher (1996)
normal stresses	Boudaoud (2007)

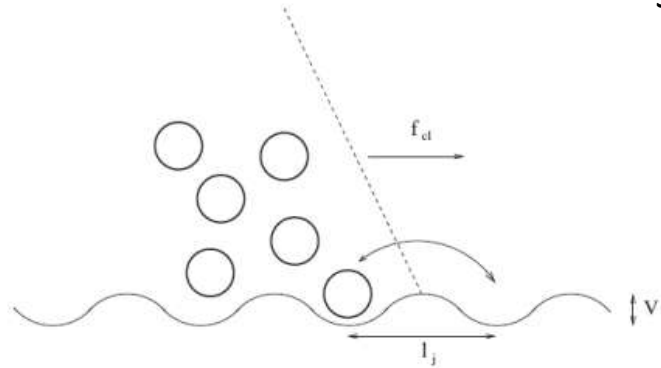


$$\sigma_{vis} \approx \frac{3\eta U}{h + 3l_s}$$



Wetting dynamics: Molecular Kinetic Theory

➤ Thermally activated process



Jump frequency

T.D. Blake, J.M. Haynes, JCIS (1969)

$$v_{\pm} = v_0 \exp\left(\frac{-V \pm Fl_j/2}{kT}\right)$$

with $v_0 \sim \frac{kT}{h} \sim 10^{13} \text{ Hz}$ $V \sim \gamma l_j^2 (1 + \cos \theta_e)$

$$F \sim \gamma l_j (\cos \theta_e - \cos \theta)$$

$$U = l_j (v_+ - v_-) = 2v_0 l_j \exp\left(-\frac{V}{kT}\right) \sinh\left(\frac{\gamma l_j^2 (\cos \theta_e - \cos \theta)}{2kT}\right)$$

Eyring model $\eta = \frac{kT}{v_0 l_j^3} \exp\left(-\frac{V_l}{kT}\right)$ with $V_l \sim 2\gamma l_j^2$

$$U = \frac{\gamma}{\eta} \exp\left(-\frac{V - V_l}{kT}\right) (\cos \theta_e - \cos \theta)$$

$$\theta^2 - \theta_e^2 \sim Ca$$

= line friction

l_j, V as adjustable parameter

Values from 0,1 to 10 nm

➤ Mixed models

- ✓ spatial scales

MKT $\rightarrow \theta_\mu(U)$ introduced in Cox-Voinov

- ✓ Velocity

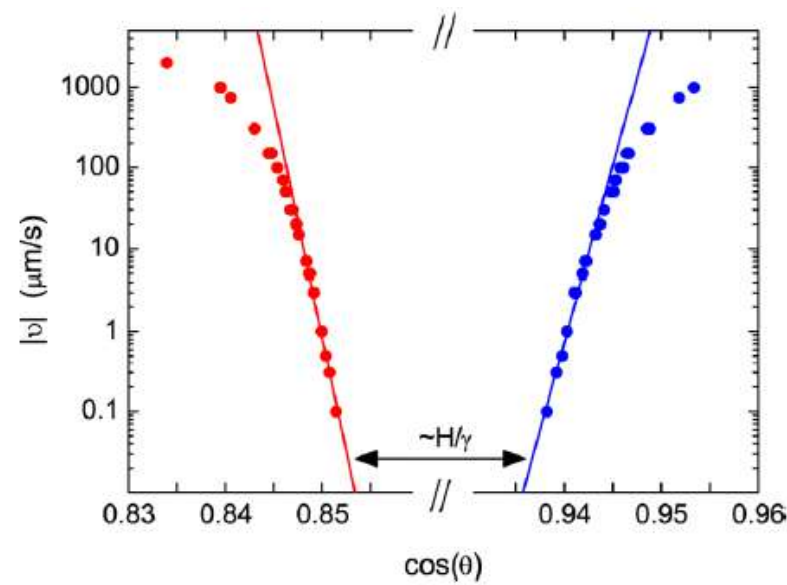
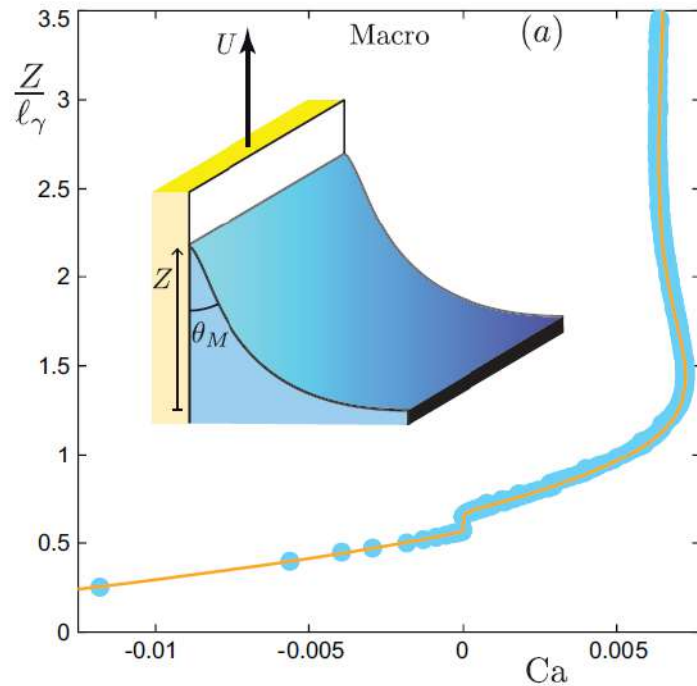
Small velocities: MKT

Large velocities: Cox-Voinov

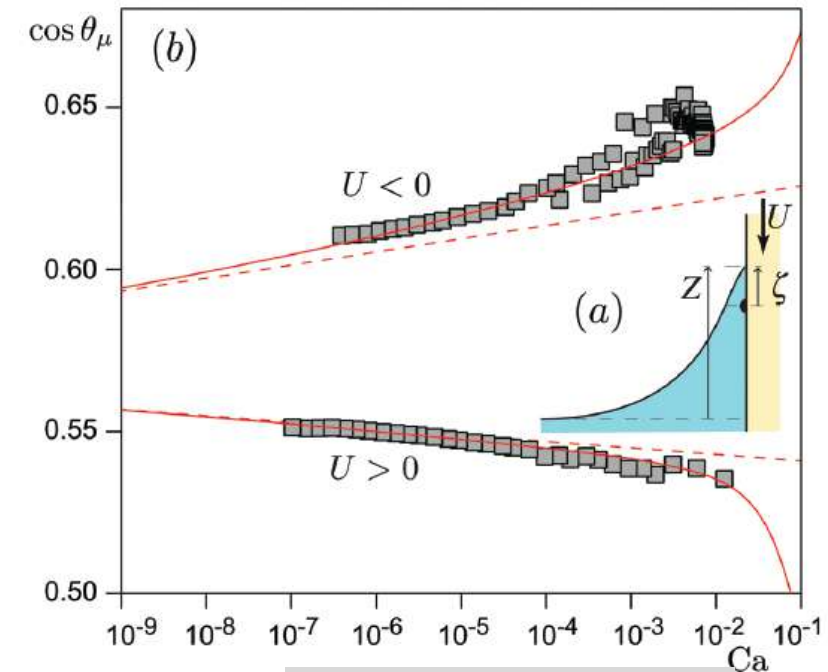
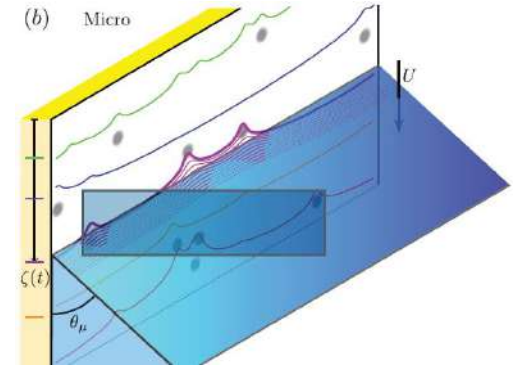
Wetting dynamics on heterogeneous surfaces

➤ Use of Cox-Voinov with static advancing or receding contact angles

➤ MKT with defects



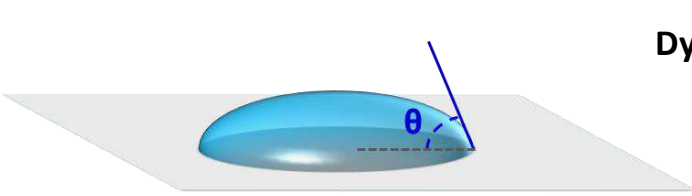
Davitt et al. Langmuir 2013



Perrin et al. PRL 2013

✓ Link with surface defect energy and size

Spreading dynamics: Tanner's law



Dynamics $\theta^3 = \frac{9\dot{R}\eta}{\gamma} \ln(L/l)$

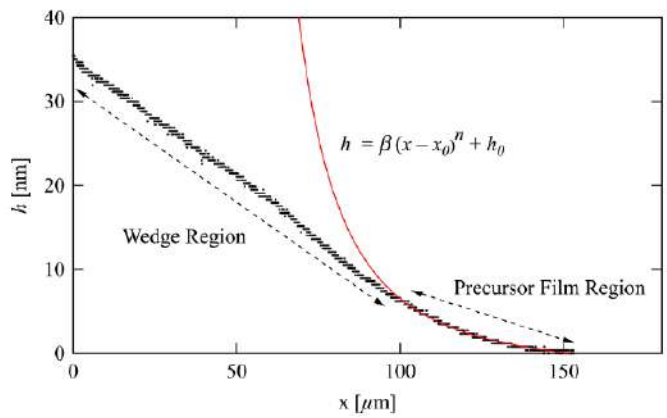
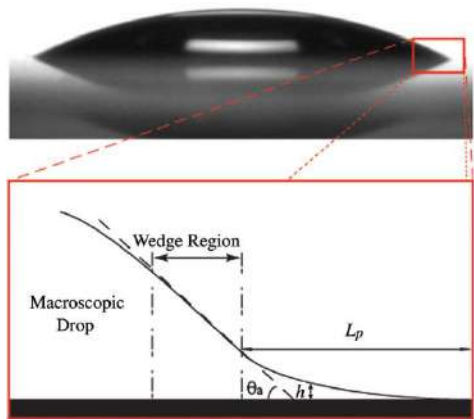
Geometry $\theta \propto \frac{1}{R^3}$

$$R \propto t^{1/10}$$

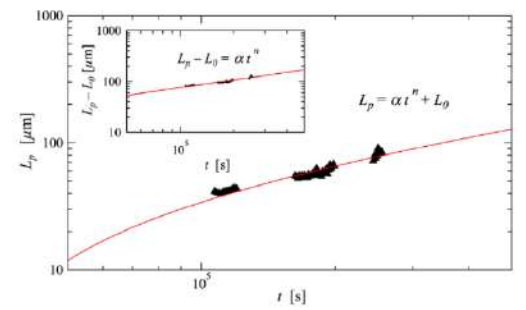
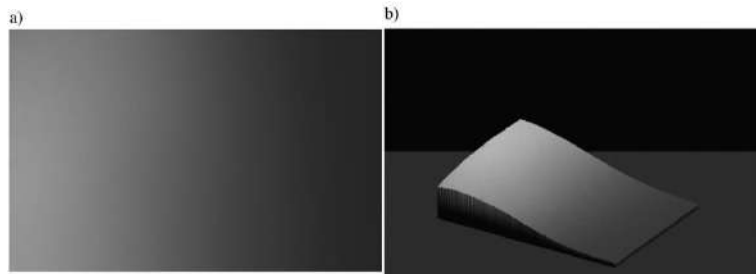
L. Tanner, *J. Phys. D* (1979)

➤ Precursor film

Huang et al., *Phys. Rev. Lett.* 2011



$$L_p = \sqrt{\frac{A}{3\pi\eta h_c}} t^{1/2}$$



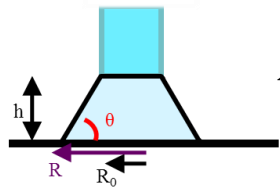
$$h(x, t) = \frac{A}{6\pi\eta} t \frac{1}{x^2}$$

Spreading dynamics: nanodispensing

L. Fabié, T.O. Soft Matt. 2012

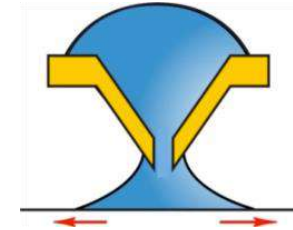
$$\theta^3 - \theta_m^3 = 9 \frac{\eta V}{\gamma} \ln\left(\frac{L}{l}\right)$$

Spreading at constant pressure



$$\tan \theta \approx \theta \approx \frac{h}{R - R_0}$$

$$\left(\frac{h}{R - R_0}\right)^3 = \theta_m^3 + 9\alpha \frac{dR}{dt}$$



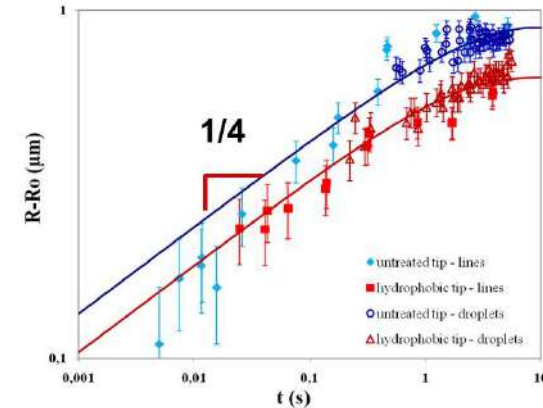
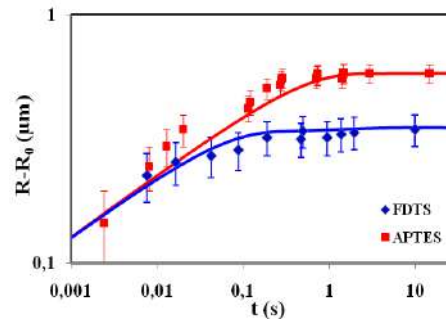
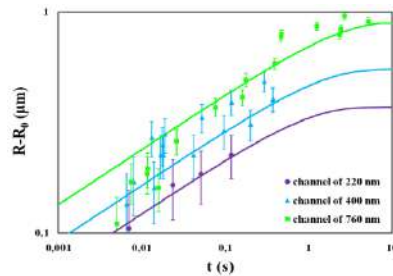
L. Fabié, T.O. Langmuir 2011

For $\vartheta_m = 0$, analytical solution :

$$R - R_0 = At^{1/4}$$

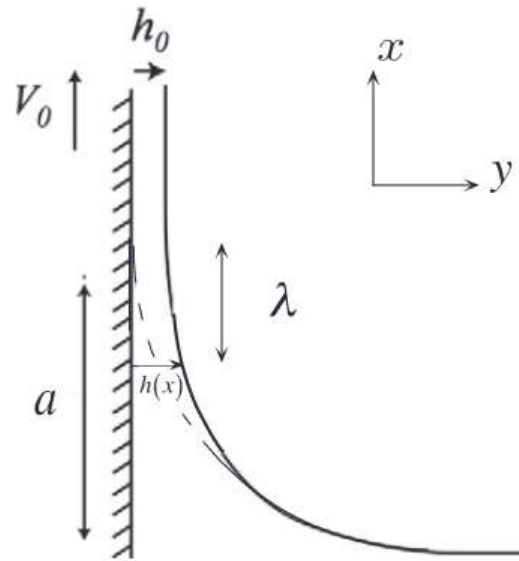
For $\vartheta_m \neq 0$, numerical solution

$$\ln(L/l) = 10$$

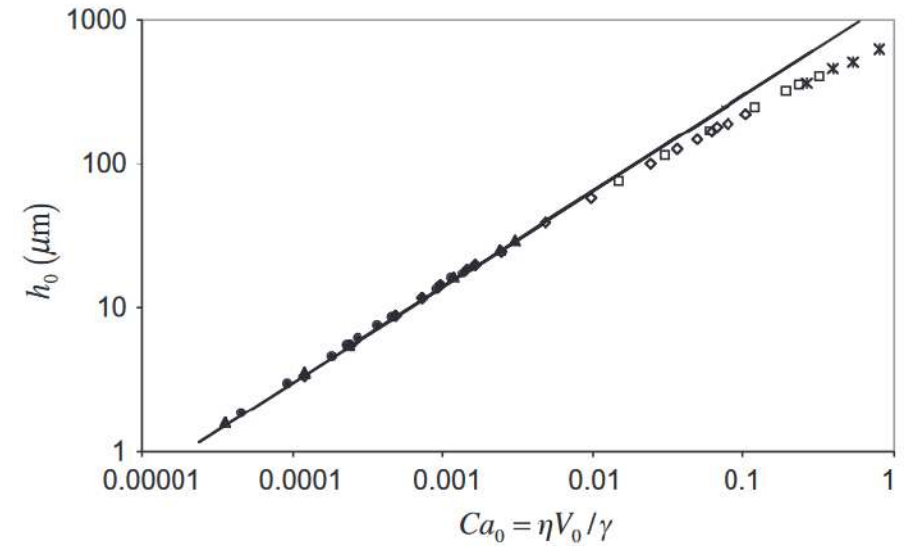


Coating: liquid entrainment

➤ The Landau-Levich film (mouillage total)



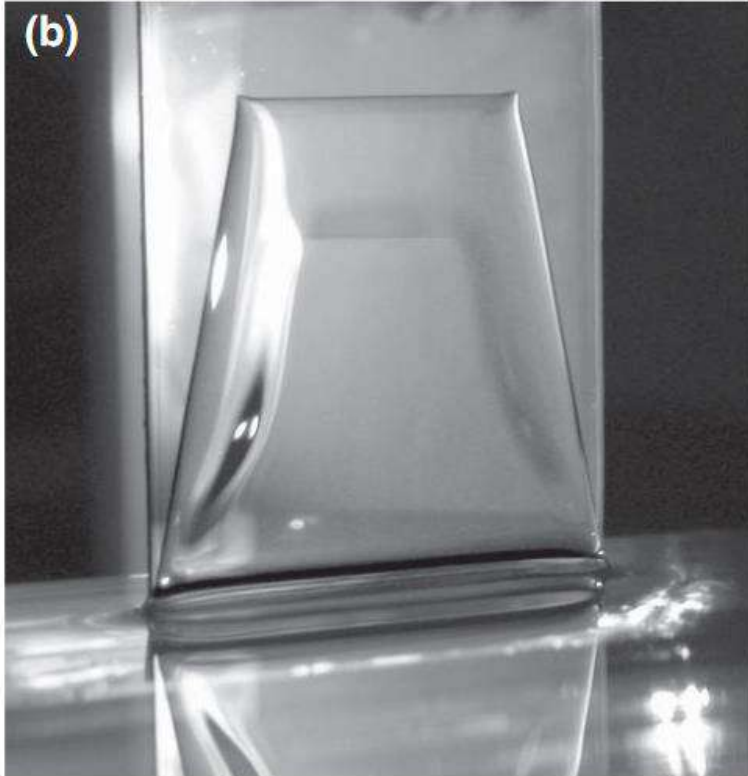
$$h_0 \approx 0.94aCa^{2/3}$$



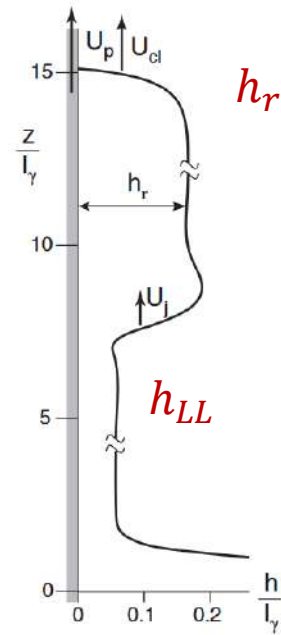
Maleki *et al.*, *JCIS* 2011

Dynamical instability

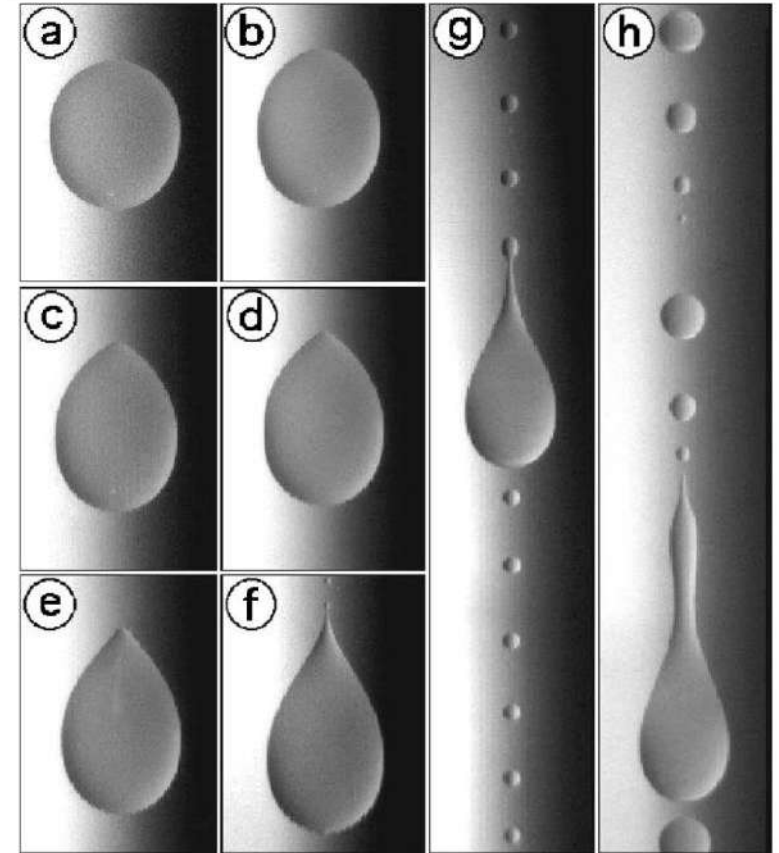
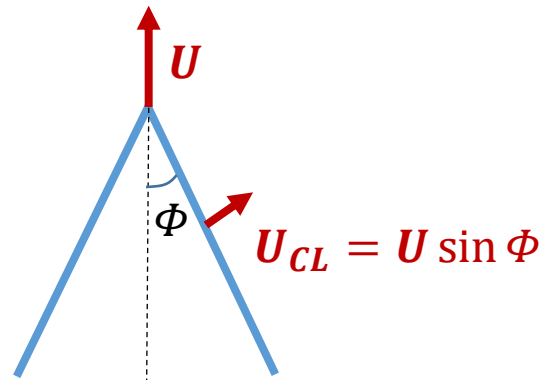
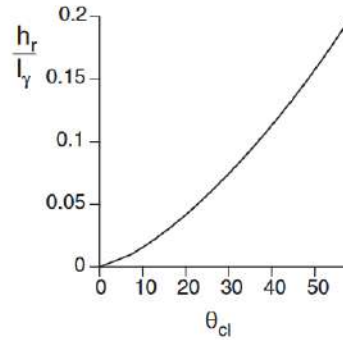
➤ Partial wetting



Snoeijer et al., PRL 2006



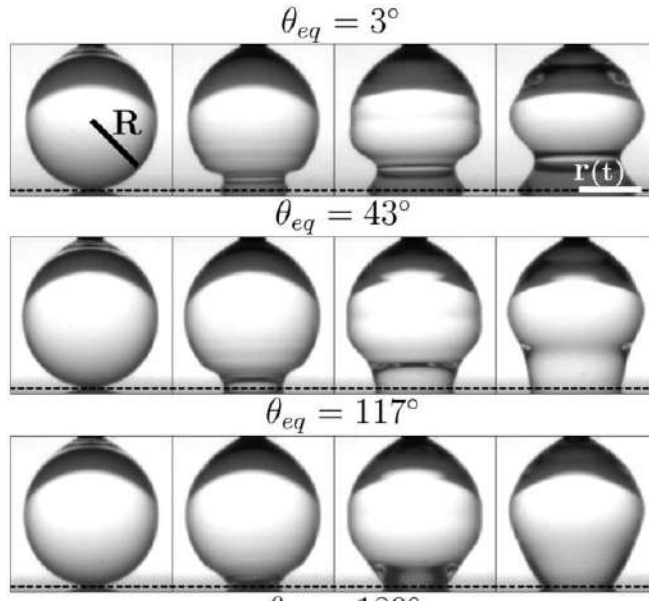
h_r determined by contact line dynamics



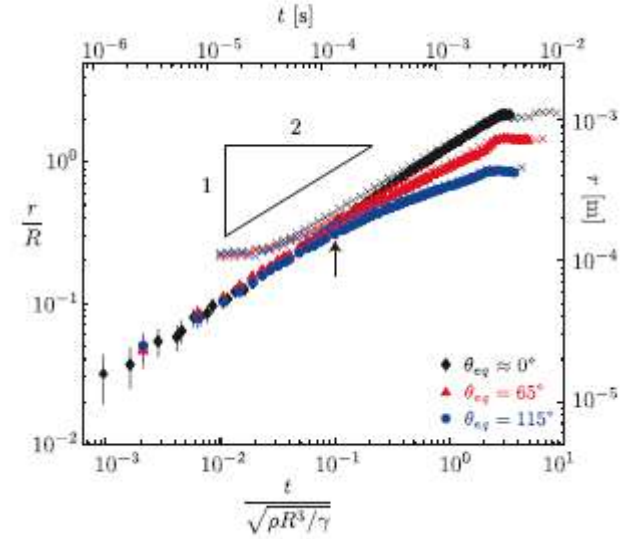
Le Grand et al., JFM (2005)

Snoeijer et al., Phys. Fluids 2007

Short time dynamics: inertial spreading



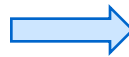
Bird, et al. *Phys. Rev. Lett* (2008)



Winkels et al., *Phys. Rev. E* (2012)

$$\rho \left(\frac{dr}{dt} \right)^2 \sim \gamma \frac{R}{r^2}$$

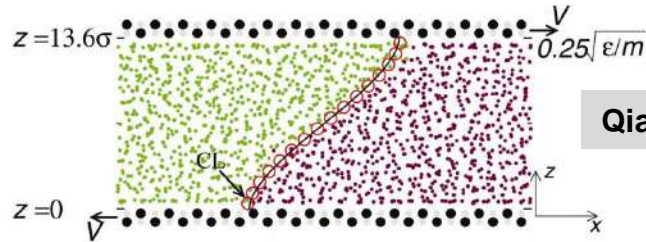
Inertia capillarity



$$r \sim t^{1/2}$$

Simulations

➤ Molecular dynamics



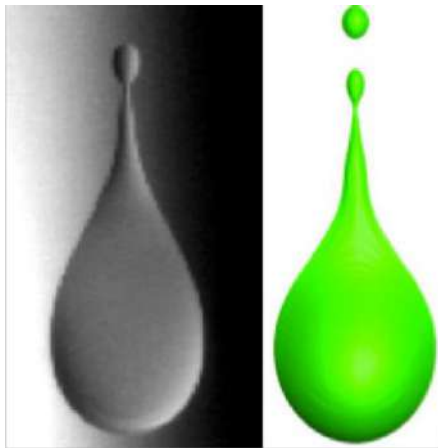
Qian et al., *Phys. Rev. E* (2003)

Sui et al., *ARFM* 2014

➤ Direct numerical simulations

- ✓ DNS solver with sub-grid model (slip or dynamics contact angle)

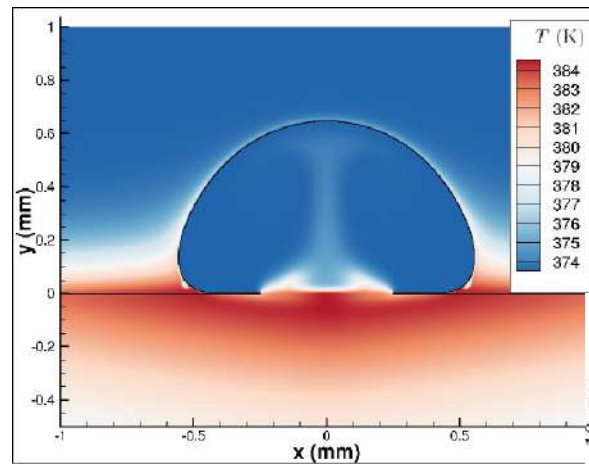
Code JADIM (IMFT)



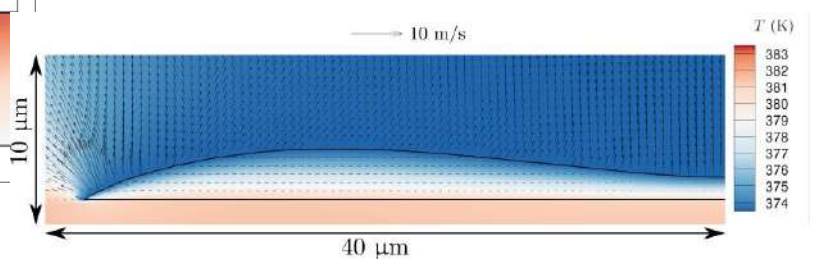
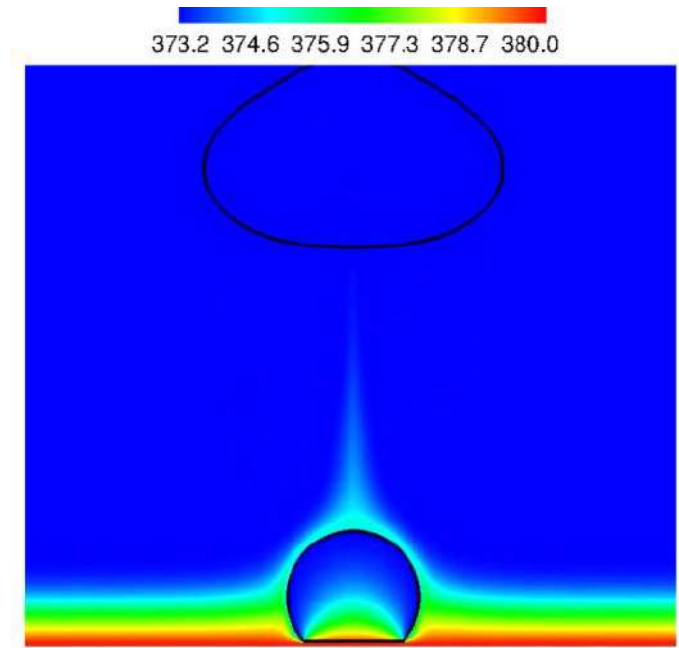
Maglio, Legendre, *Comp. Exp. Fluid Mech.* (2014)

Le Grand et al., *JFM* (2005)

Code DIVA (IMFT)

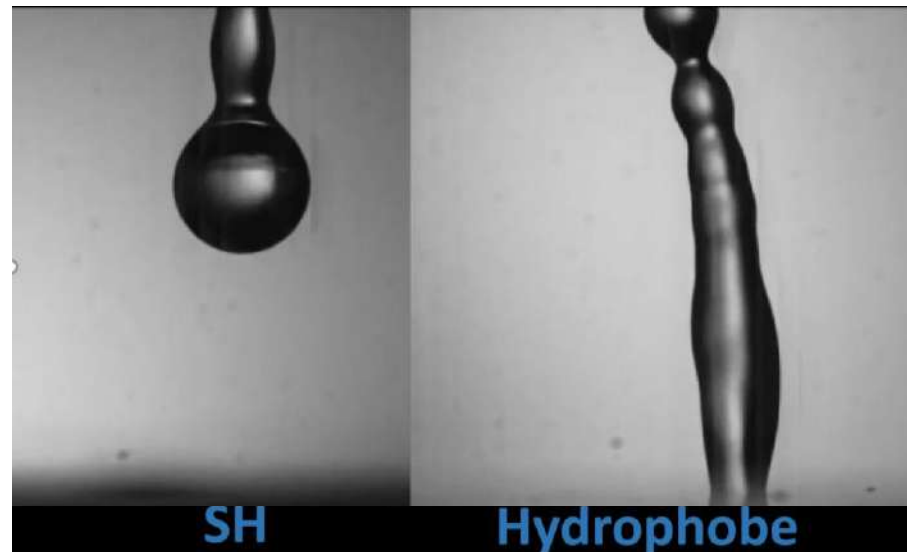
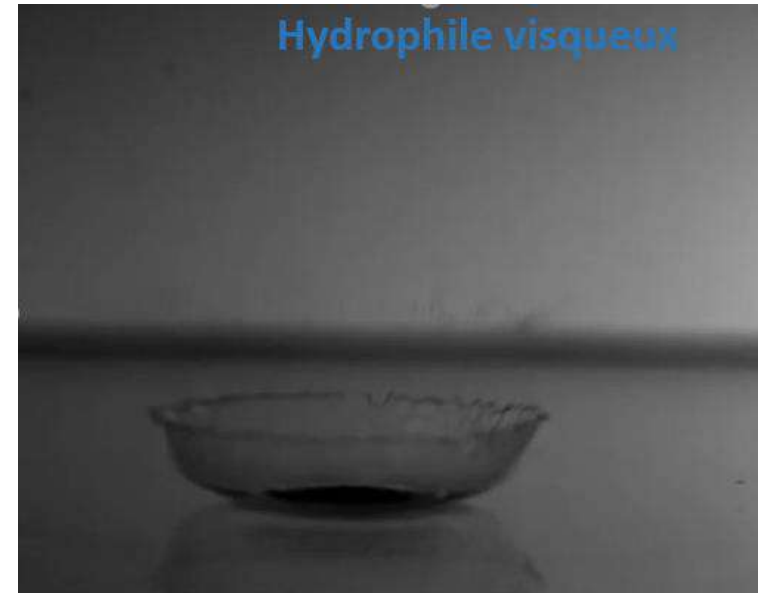
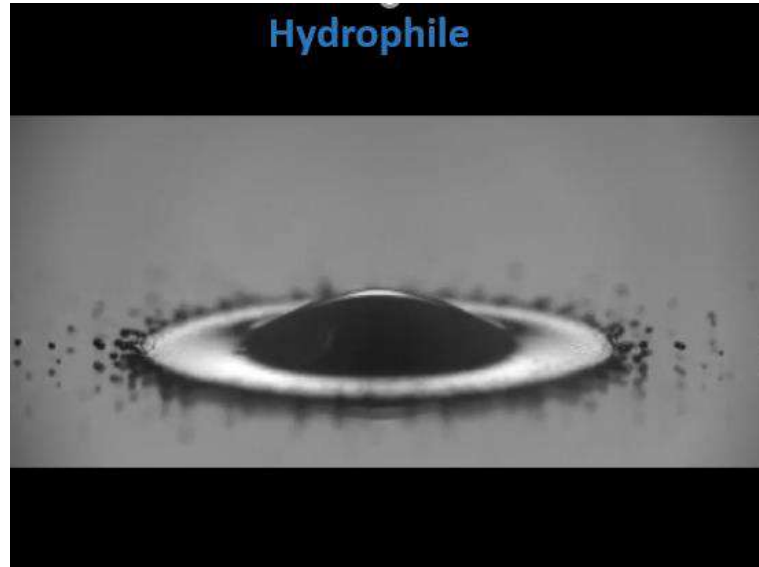


Microlayer regime



Urbano et al., *IJHMT* (2018)

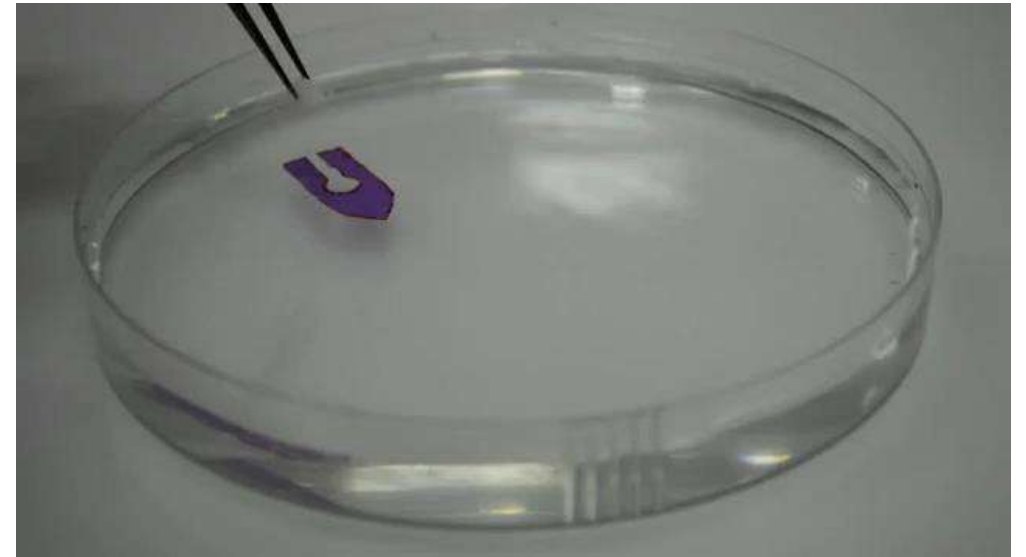
Impact



Wetting with complex liquids

- Marangoni effect
- Wetting with volatile liquids
- Reactive wetting
- Wetting with surfactants solutions

Marangoni effect



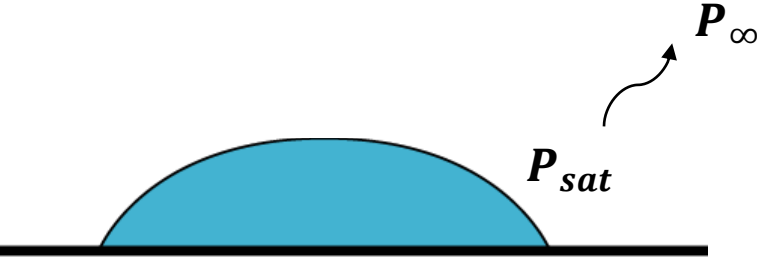
Ex : camphor boat

Solutal Marangoni effect $\gamma(C)$: surfactants $\gamma(C) \downarrow$, binary solutions

Thermal Marangoni effect $\gamma(T) \downarrow$

Wetting with volatile liquids

➤ Droplet evaporation



$$m = \frac{1}{3} \pi r^3 \rho (1 - \cos \theta)^2 (2 + \cos \theta)$$

Diffusion controlled

$$\frac{dm}{dt} = -4\pi r^2 D \left(\frac{dc}{dr} \right) = -4\pi r D (c_0 - c_\infty)$$

$$\frac{dm}{dt} \propto r$$

Picknett, Bexon *J.C.I.S.* 1977

Tapez une équation ici.

Constant contact angle mode:



$$\frac{dm}{dt} = -\alpha m^{1/3}$$



$$m^{2/3} = m_0^{2/3} - \frac{2}{3} \alpha t$$

$m^{2/3}(t)$ linear

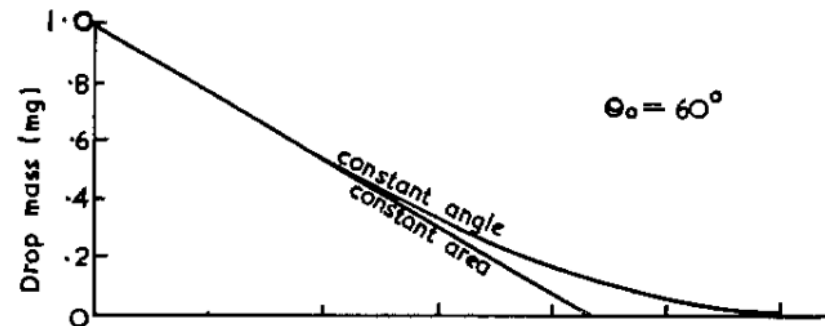
$t_{evap} \propto m_0^{2/3}$

Constant radius R mode:



$$r \sin \theta = R$$

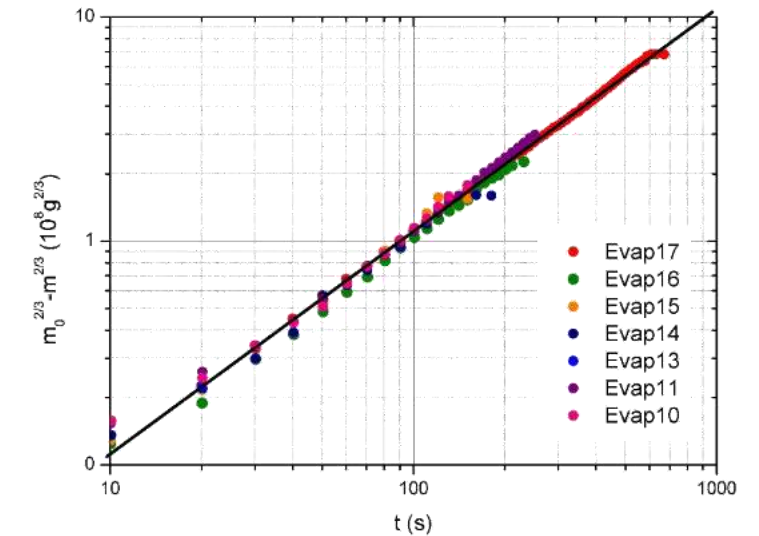
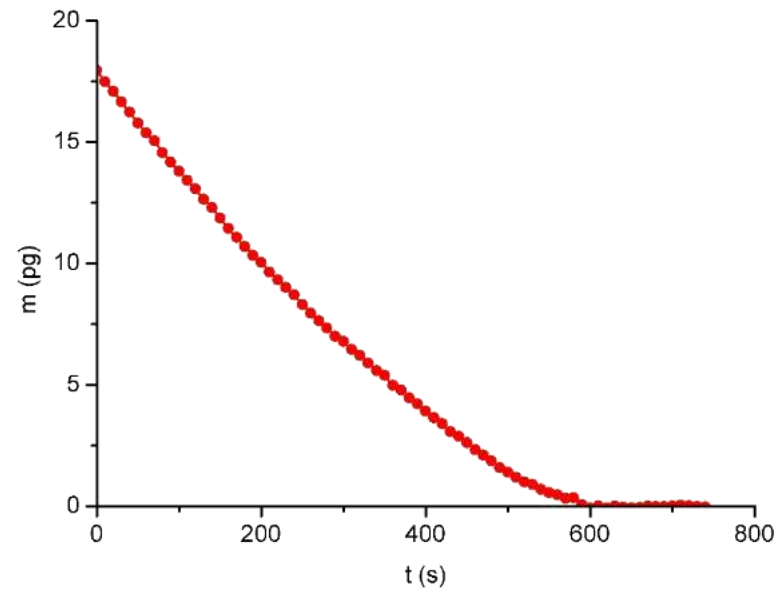
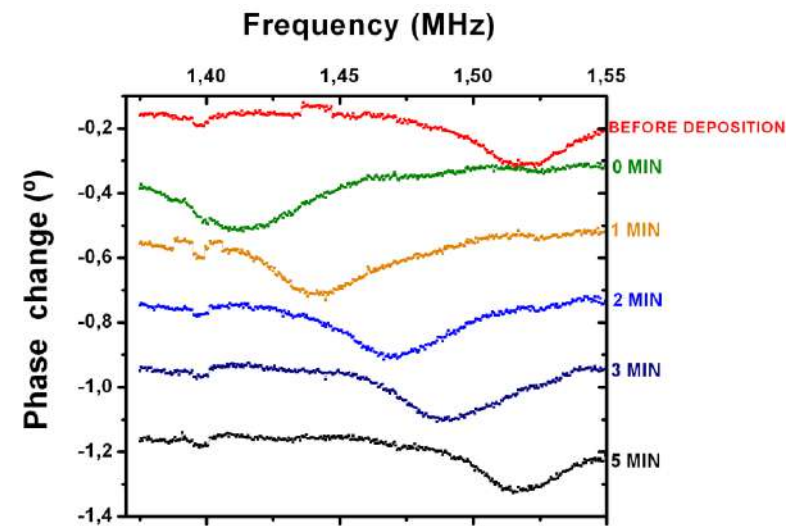
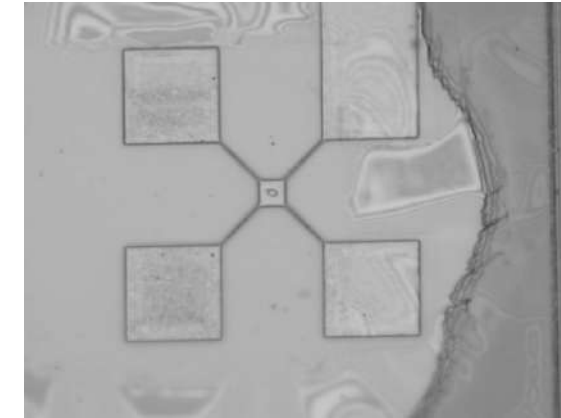
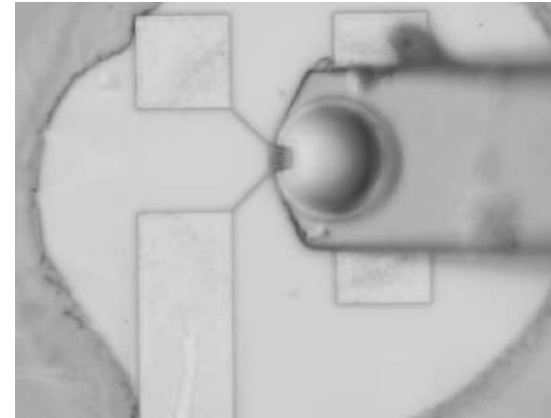
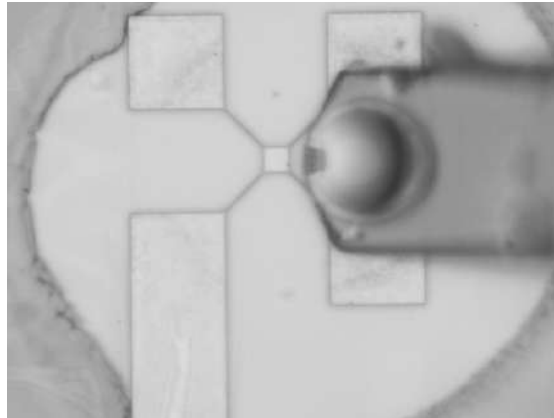
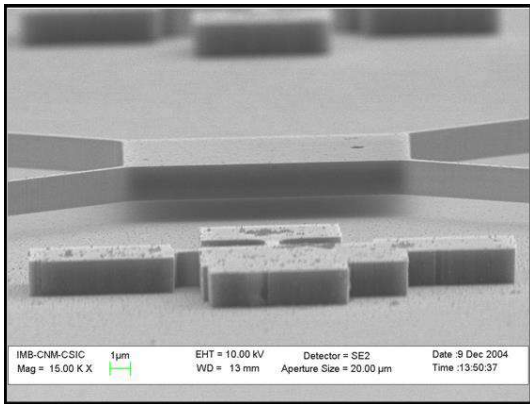
$$\frac{dm}{dt} = -Cte$$



Wetting with volatile liquids

Arcamone, *J.Phys. Chem. B* 2007

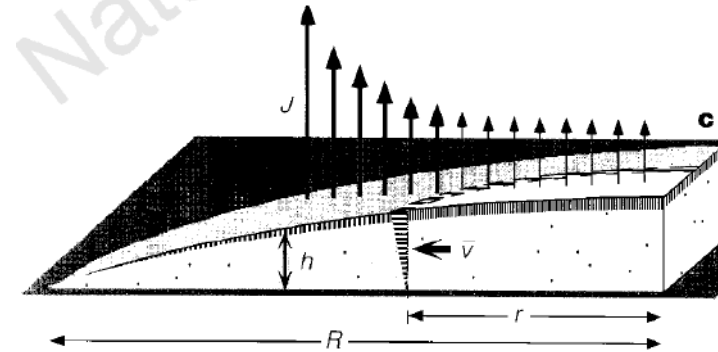
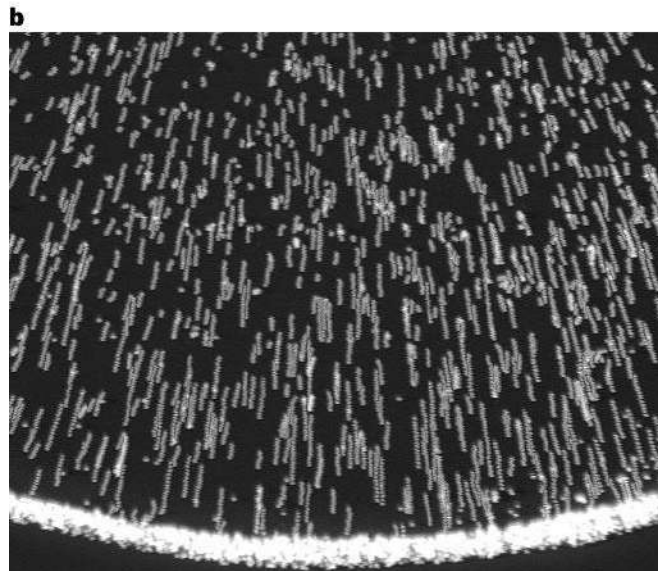
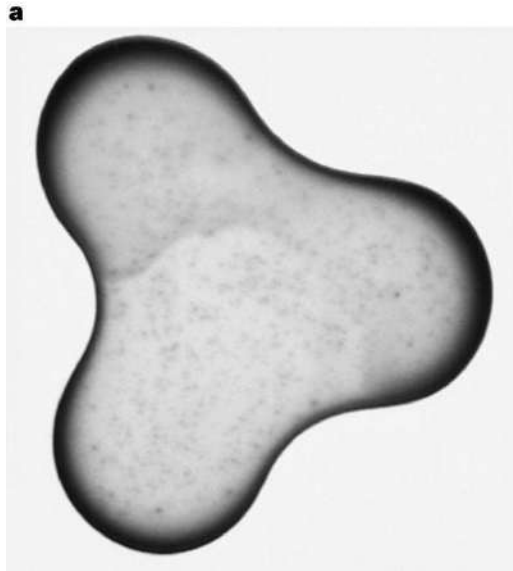
➤ Evaporation of femtoliter (10^{-15} l) droplets



Wetting with volatile liquids

➤ Deposition pattern: the coffee stain

Deegan *et al.*, *Nature* 1994



$$J(r) \propto (R - r)^{-\lambda} \quad \text{with} \quad \lambda = (\pi - 2\theta)/(2\pi - 2\theta)$$

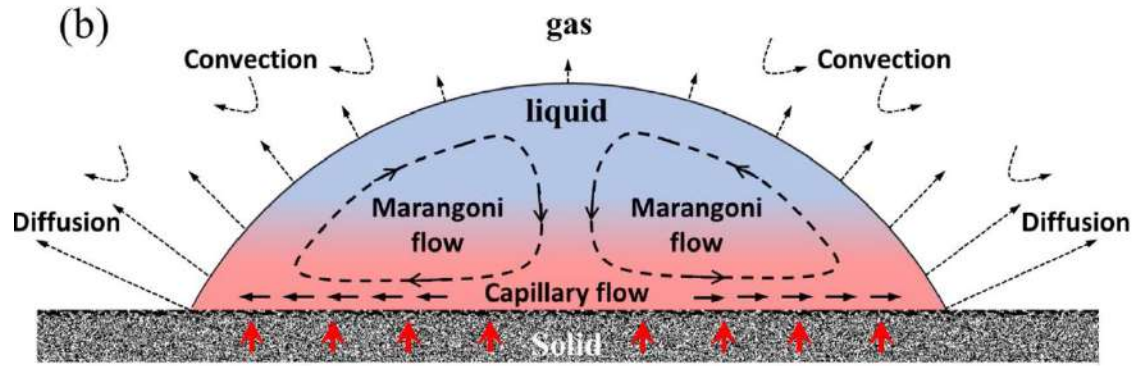
Pinned contact line



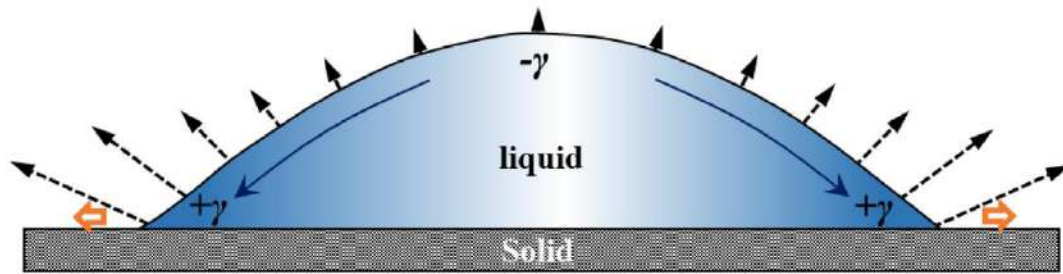
➡ **Capillary flow**
towards the contact line

Wetting with volatile liquids

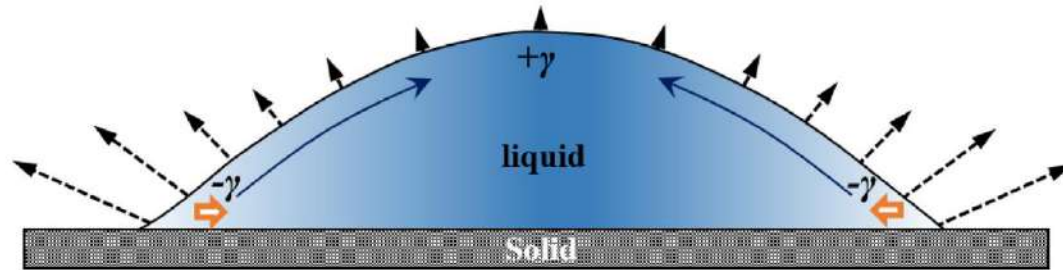
Wang et al., Phys. Rep. 2022



Binary mixture

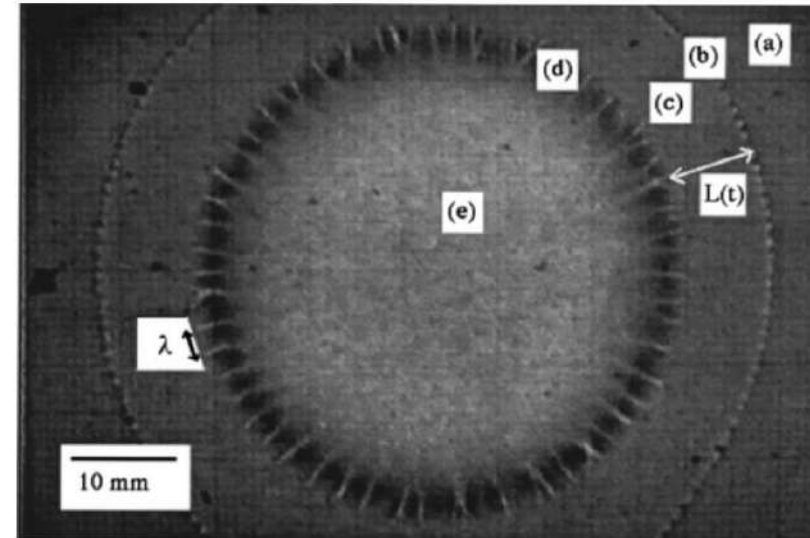


(a) Marangoni Spreading



(b) Marangoni Contraction

Water/alcohol mixture



Fanton et al., Langmuir 1998

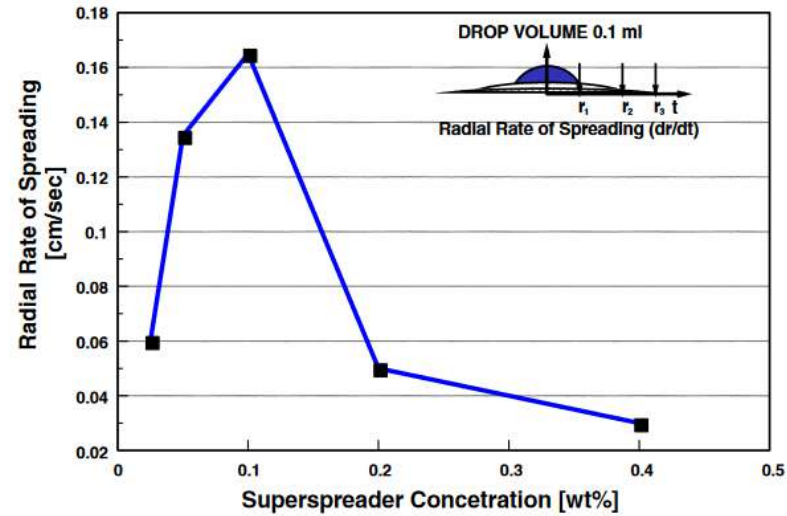


Wetting with complex liquids: surfactants

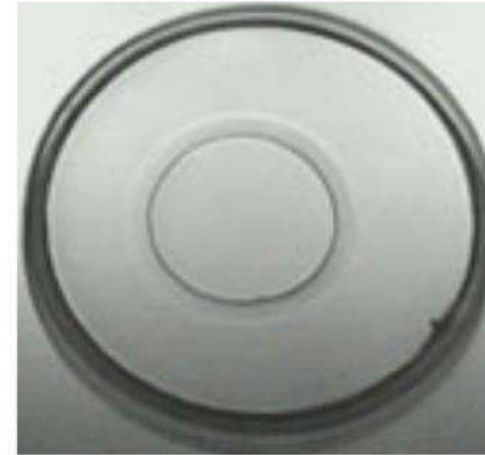
➤ A particular case: superspreading on hydrophobic surfaces

Nikolov et al., *Current Opinion Coll Interf Sci.* 2020
EPJST 2011

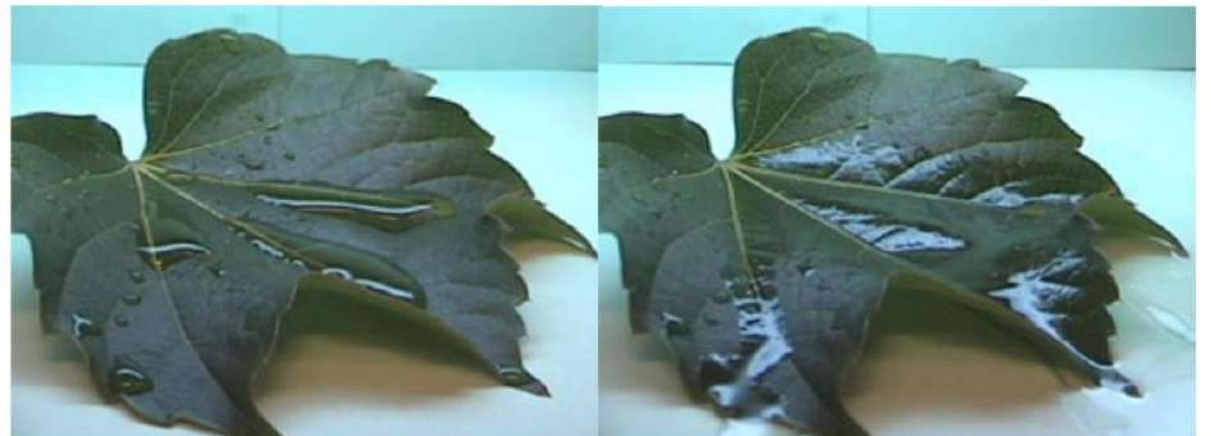
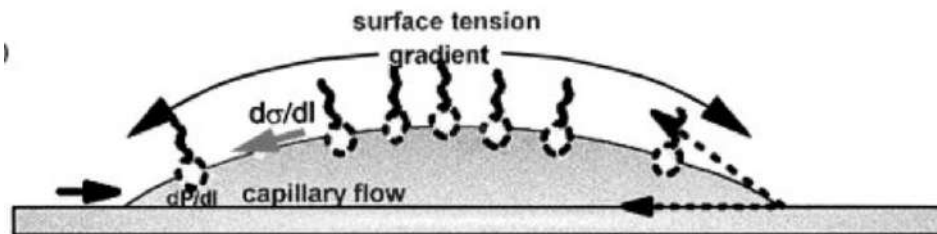
1. Initial spreading



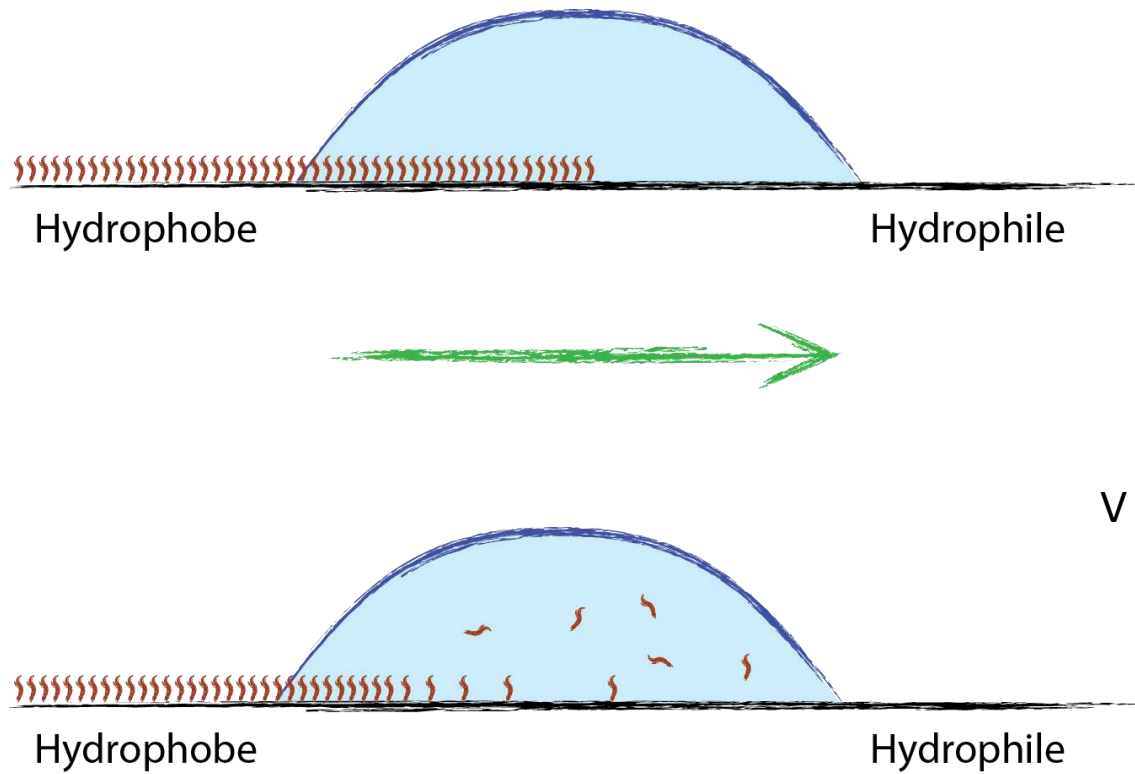
2. Bump



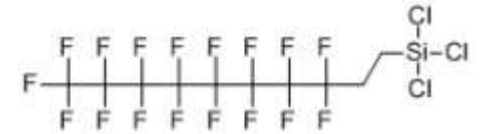
3. Instability



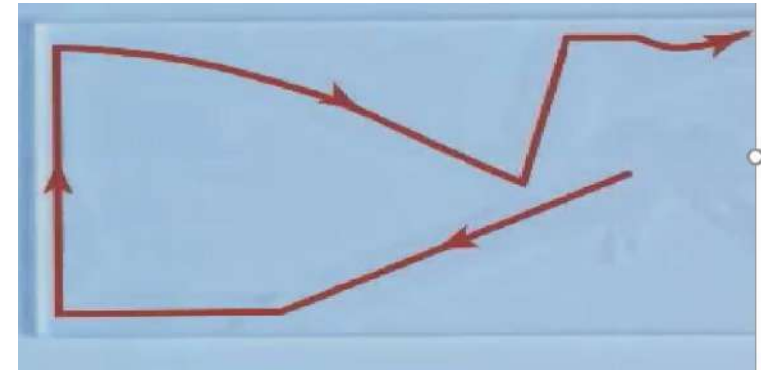
Reactive wetting: self-propelled droplets



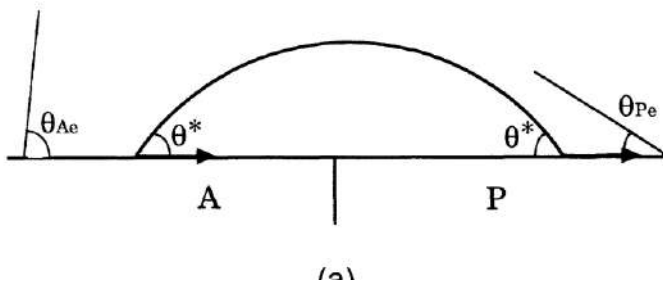
1H,1H,2H,2H, perfluorodecyltrichlorosilane
in octane



$V \sim 10 \text{ cm/s}$



Reactive wetting: self-propelled droplets



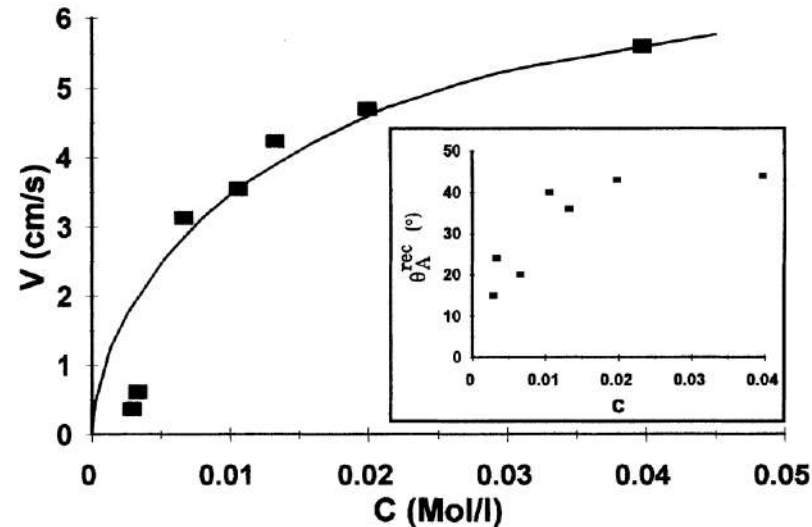
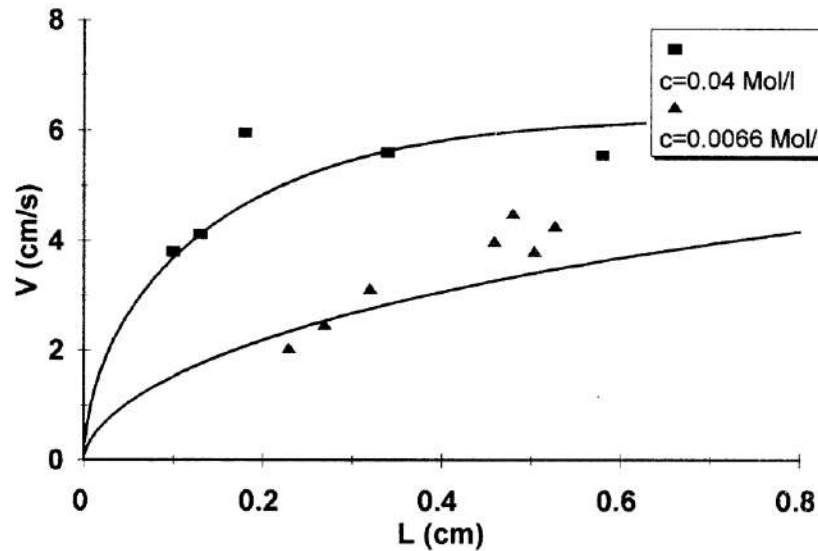
$$\begin{cases} V_A(\theta^*) = \frac{\gamma}{9\eta l} (\theta_A^3 - \theta^{*3}) \\ V_P(\theta^*) = \frac{\gamma}{9\eta l} (\theta^{*3} - \theta_P^3) \end{cases} \quad V = \frac{\gamma}{18\eta l} (\theta_A^3 - \theta_P^3)$$

$$\gamma \cos \theta_A = \gamma \cos \theta_P - \gamma_1 \Phi_S$$

$$\Phi_S = 1 - \exp\left(-\frac{t}{\tau}\right) = 1 - \exp\left(-\frac{L}{V\tau}\right)$$



$V(C)$ and $V(L)$

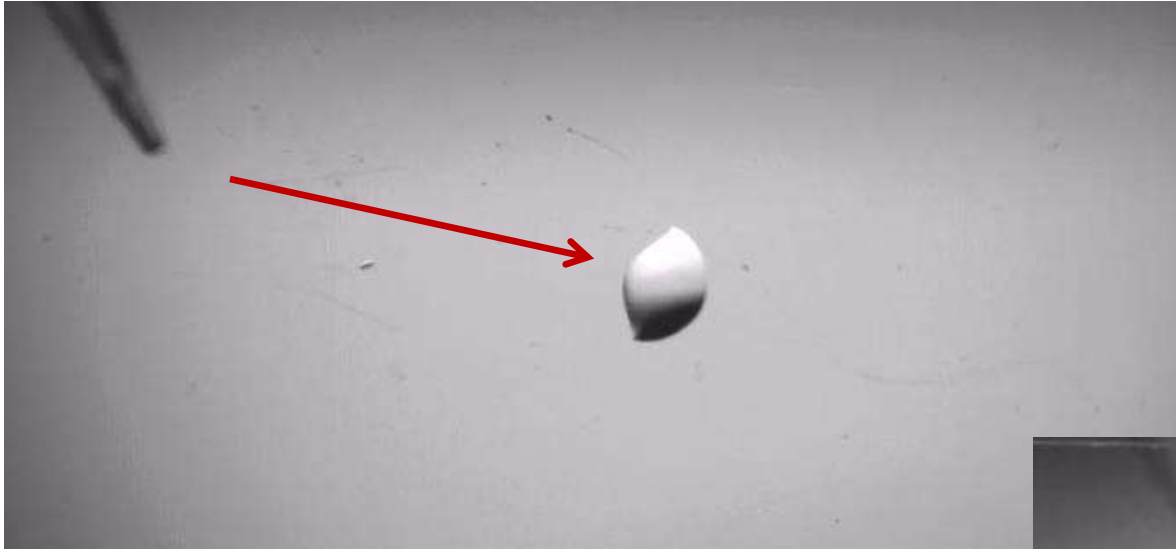


$\tau \in [0,01 \text{ s} - 0,2 \text{ s}]$

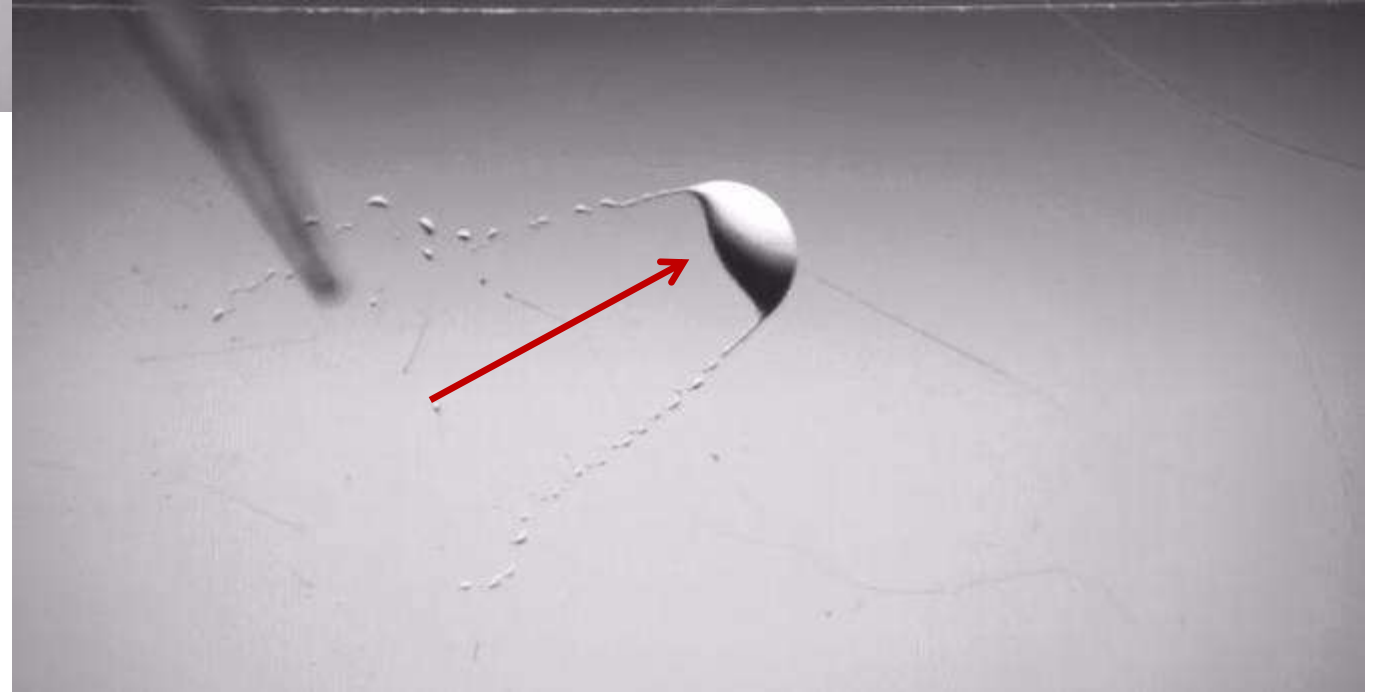
Reactive wetting: self-propelled droplets

Perfluorosilane 0,02 M sur **lame de verre brute**

Beaune *et al.*, *PNAS*. 2018

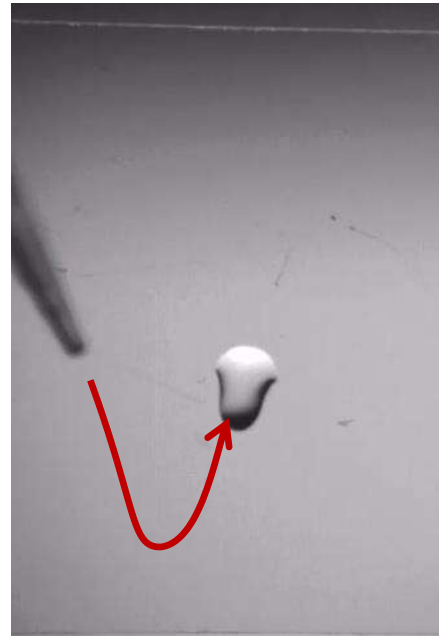


sur **lame de verre traitée corona (mouillage total)**

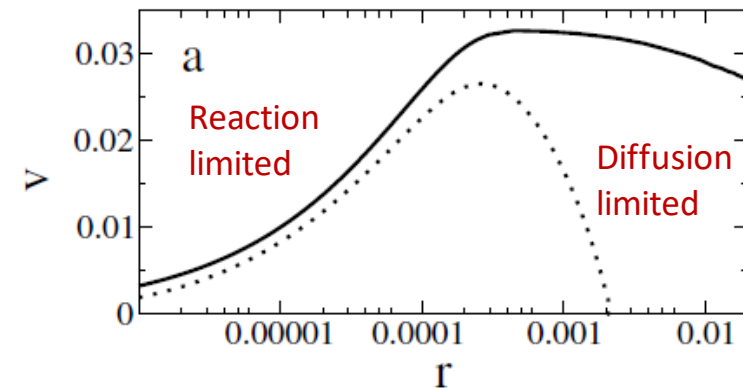
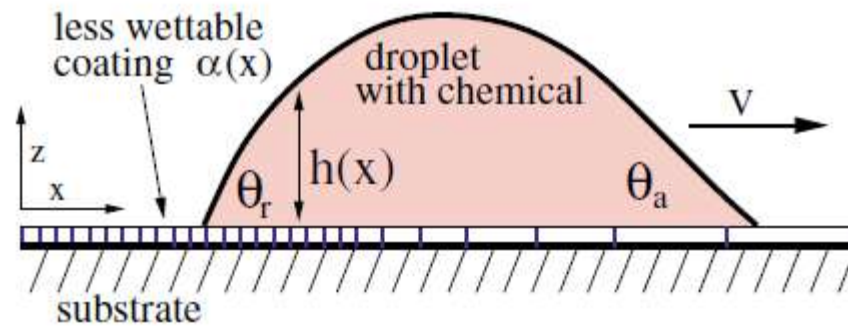
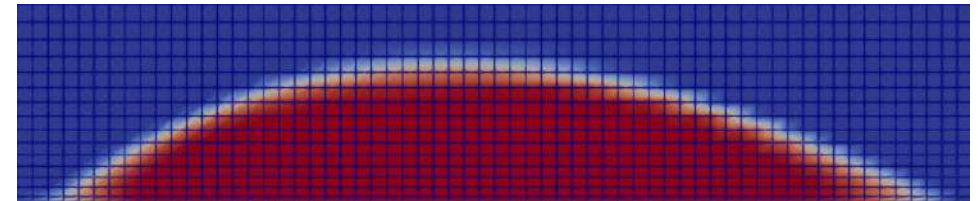


Reactive wetting: self-propelled droplets

Perfluorosilane 0,05 M sur **lame de verre brute**



Simulations JADIM



Wetting and adsorption: contamination and cleaning

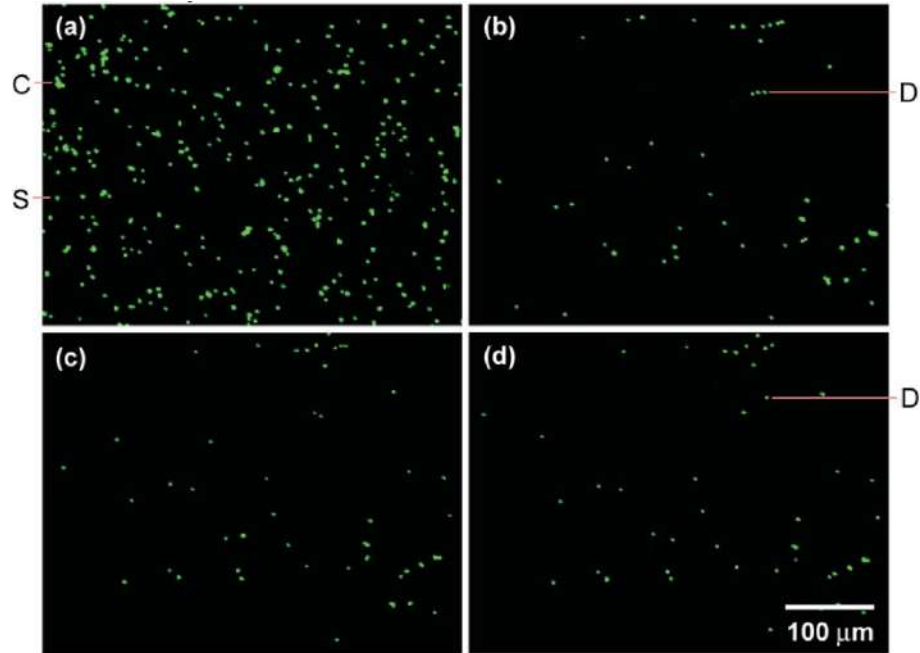
- Water is totally wetting clean glass surface
 - ✓ Cleaning can be achieved by UV/Ozone cleaning, O₂ plasma, piranha solution...
- Left in air, glass becomes hydrophobic due to airborne contaminant adsorption
 - ✓ Short hydrocarbons (butane, propane) are the main source of contamination
 - ✓ Water contact angle : 30-40 °

Millet et al. *J. Geophys. Reseach Atm.* 2005



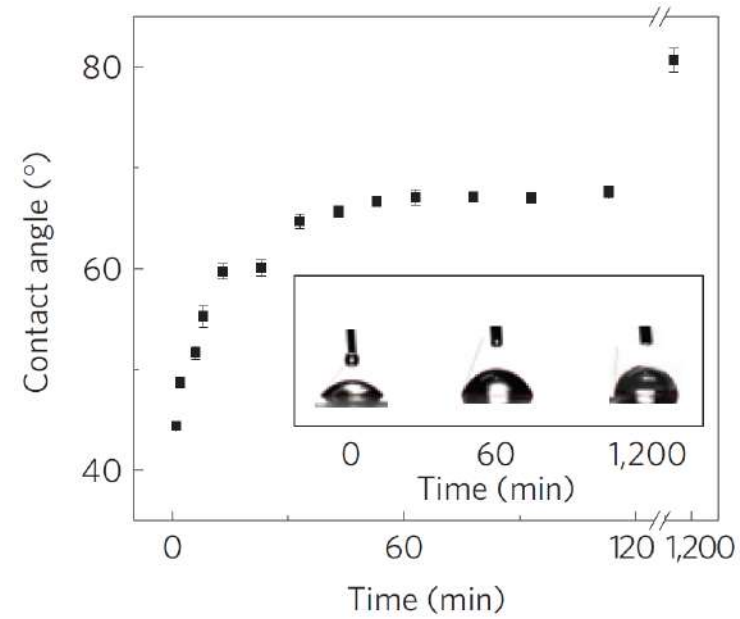
Wetting and adsorption: contamination and cleaning

➤ Contamination by particles

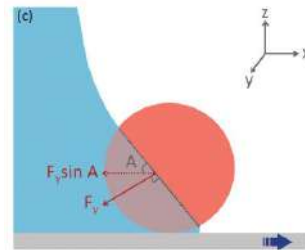


Sharma et al. JCIS 2008

➤ Chemical contamination by airborne molecules

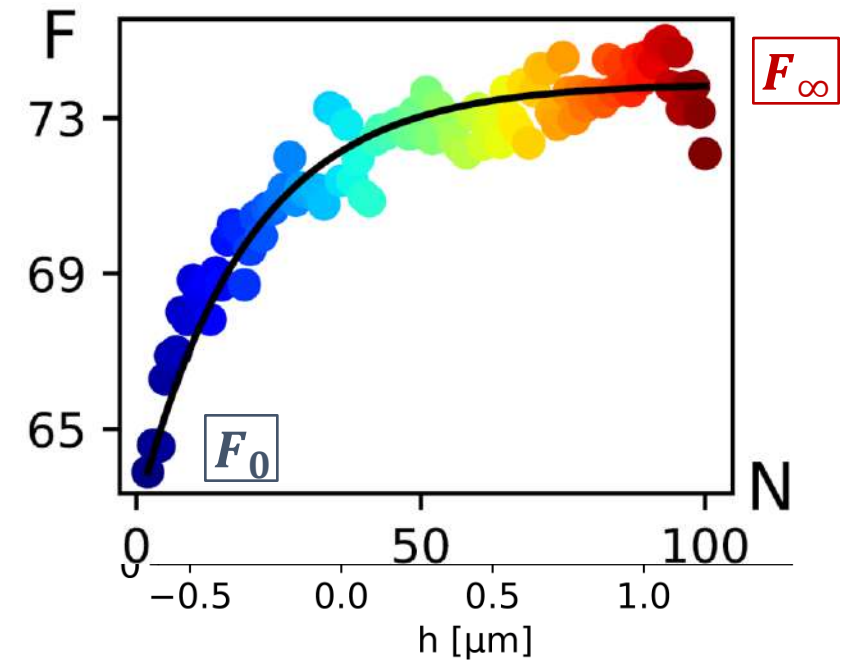
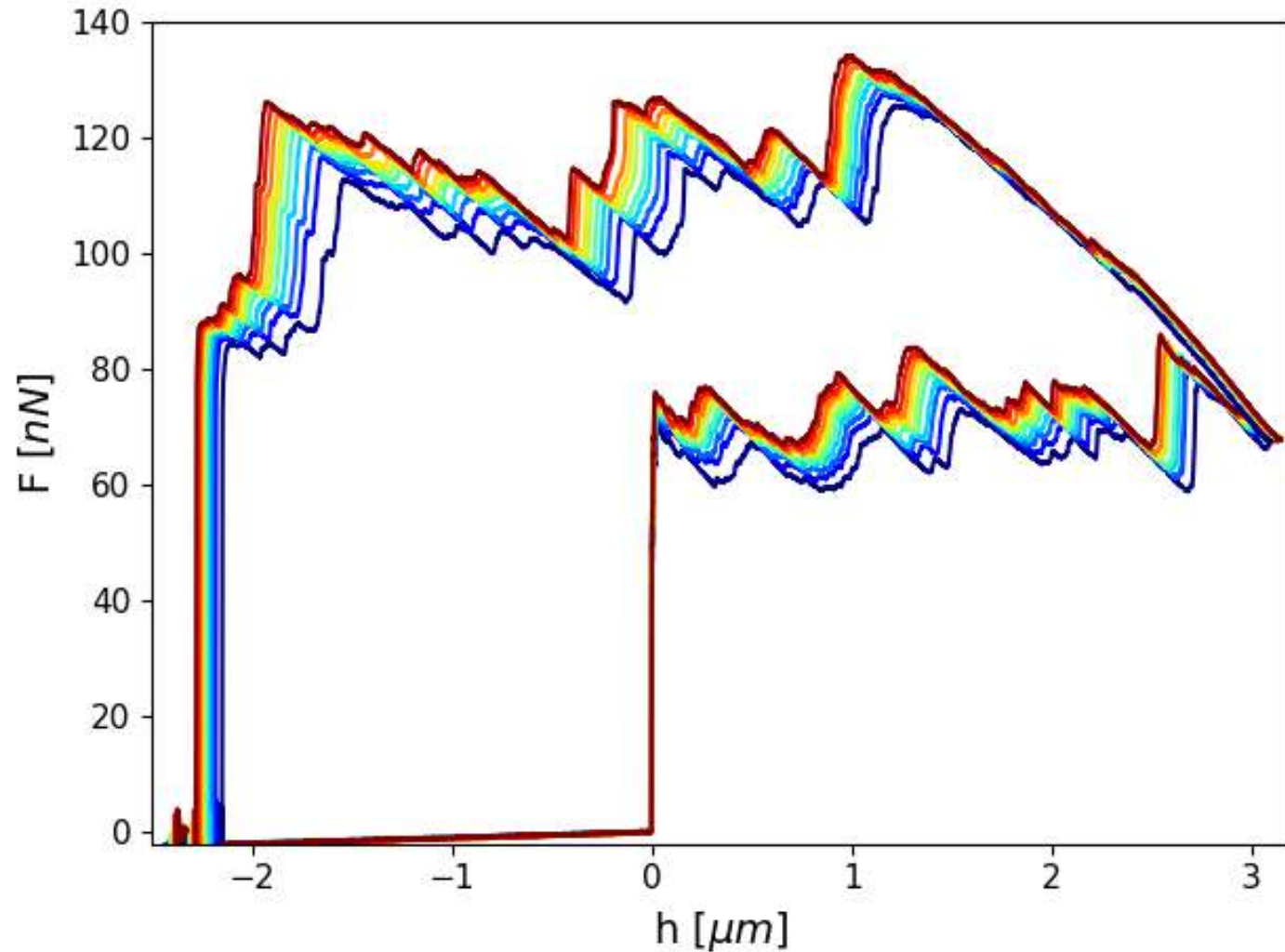


Liu et al. Nature Mat 2013



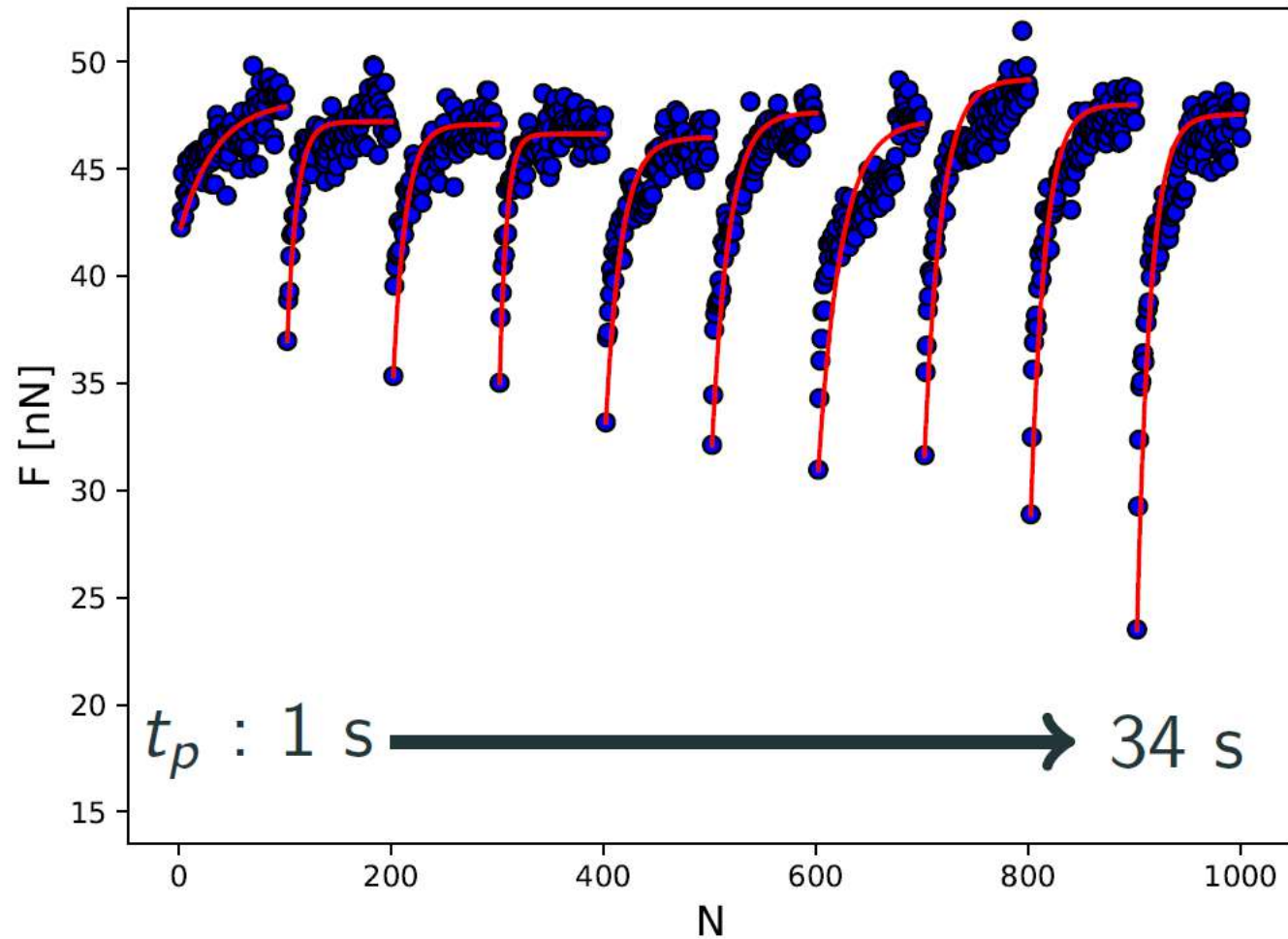
Aramrak et al. Soft Matt 2021

Wetting and adsorption: contamination and cleaning



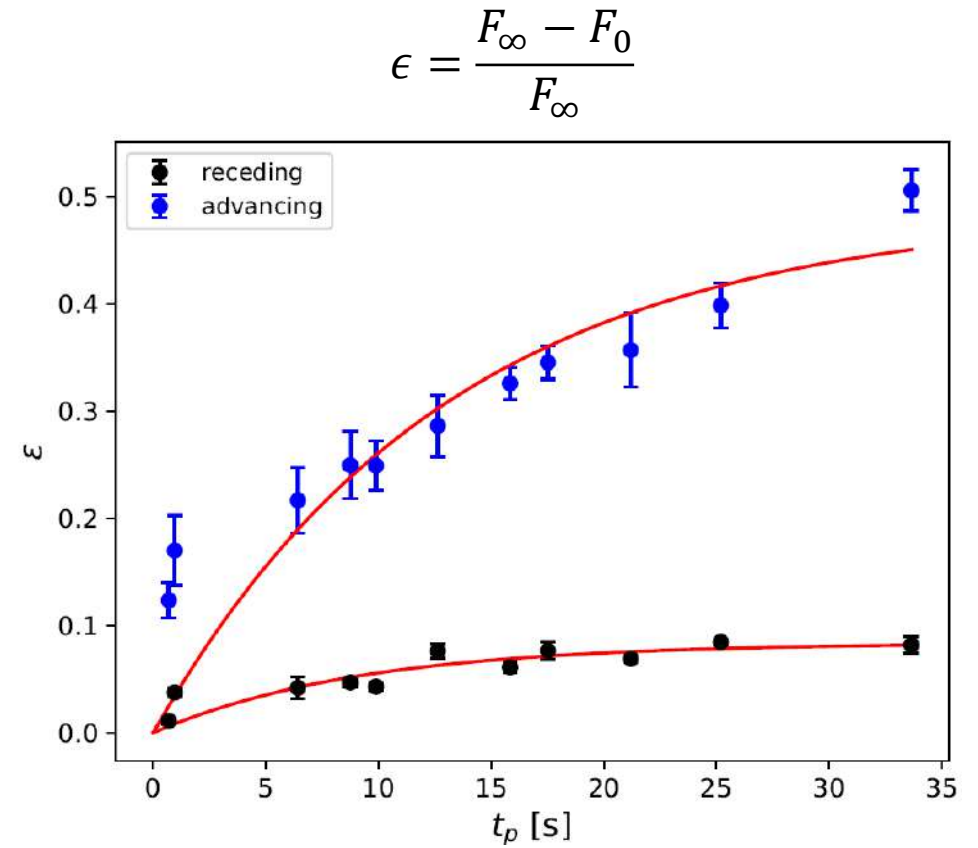
$$F = F_0 + (F_\infty - F_0) \left(1 - e^{-\frac{N}{N^*}} \right)$$

Wetting and adsorption: contamination and cleaning



$$F = F_0 + (F_\infty - F_0) \left(1 - e^{-N/N^*} \right)$$

Desorption



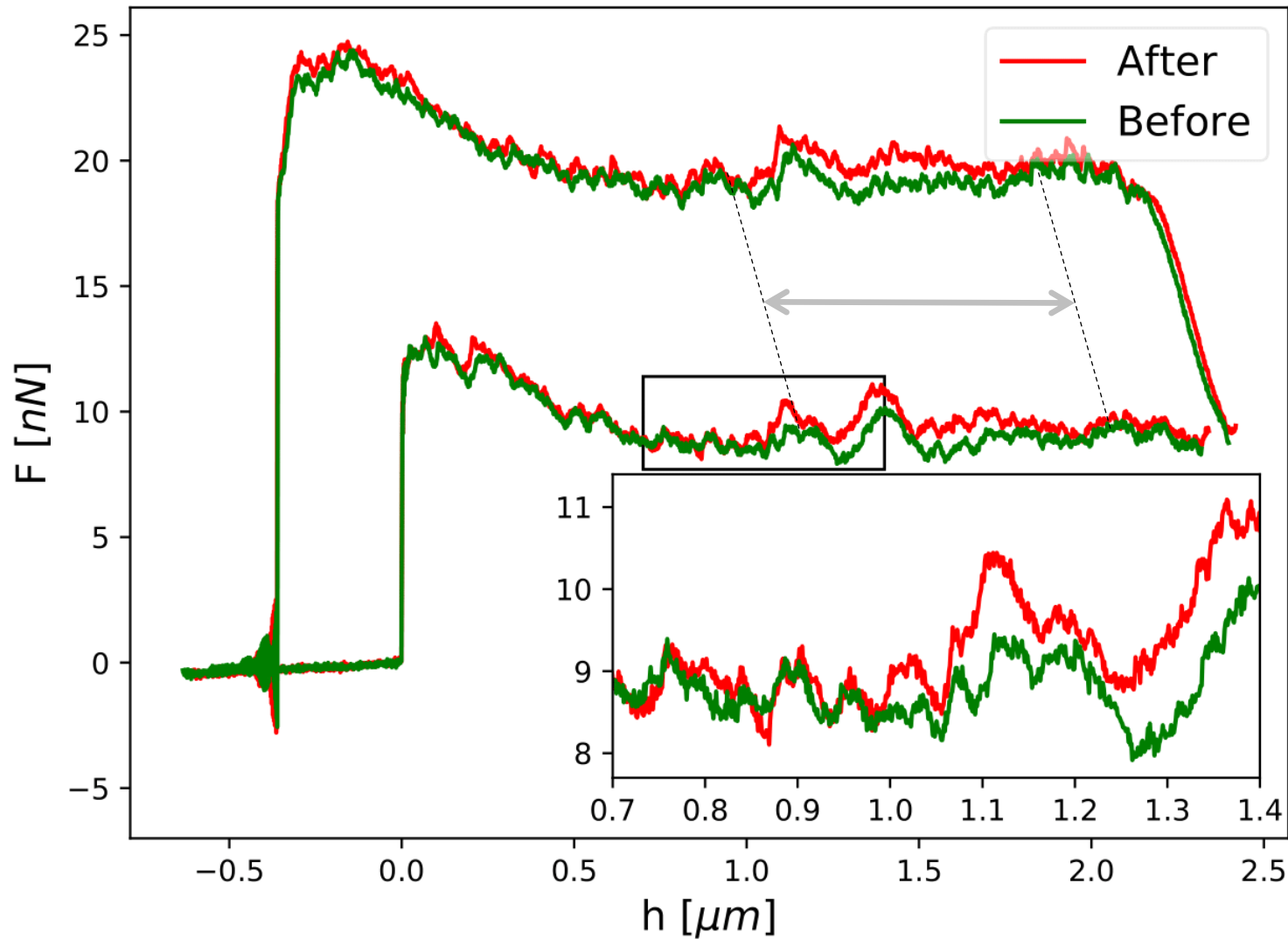
$$\epsilon = \frac{F_\infty - F_0}{F_\infty}$$

$$\epsilon = \epsilon_{max} \left(1 - e^{-t/\tau} \right)$$

Adsorption

Wetting and adsorption: contamination and cleaning

Franiatte *et al.* *PRL* 2021
Langmuir 2022



- Force curve until $h = 2.5$ μm
- 1000 cycles between $h = 1$ μm and 2 μm
- Force curve until $h = 2.5$ μm

→ Desorption at the contact line