



Structural Role of Elements in Glasses from Classical Concept to a Reflexion over Broad Composition Range

Cargèse

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Organic Glasses

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Glassforming molecular liquids and polymers

I- phenomenology of the glass formation
key features of interest

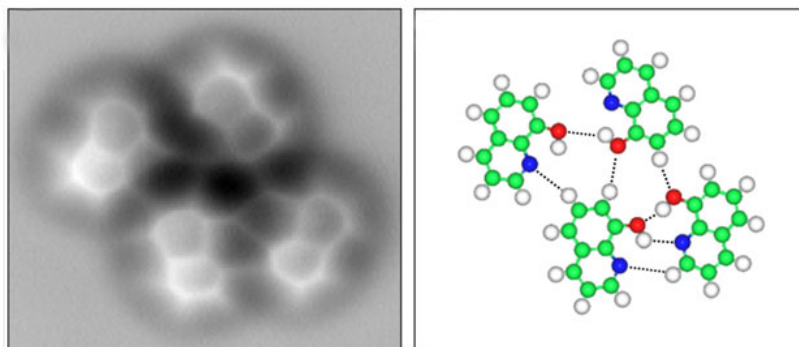
II- Notion of “Fragility”
role of control parameter temperature and density
effect of intermolecular interactions

III- local order and mesoscopic order ;
underlying length scales
high T activation energy
ultra stable glasses

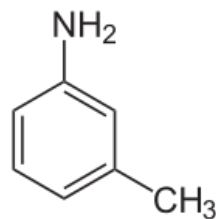
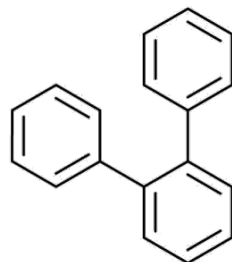
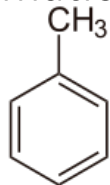
IV- molecular networks ?
Reflexion over Broad Composition Range
solutions and mixtures

Molecules and polymers

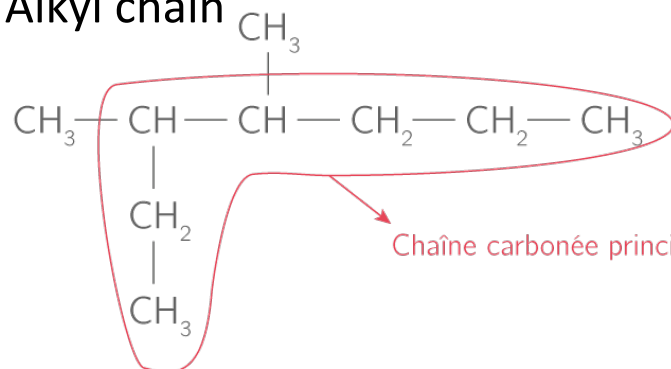
AFM 2009 IBM



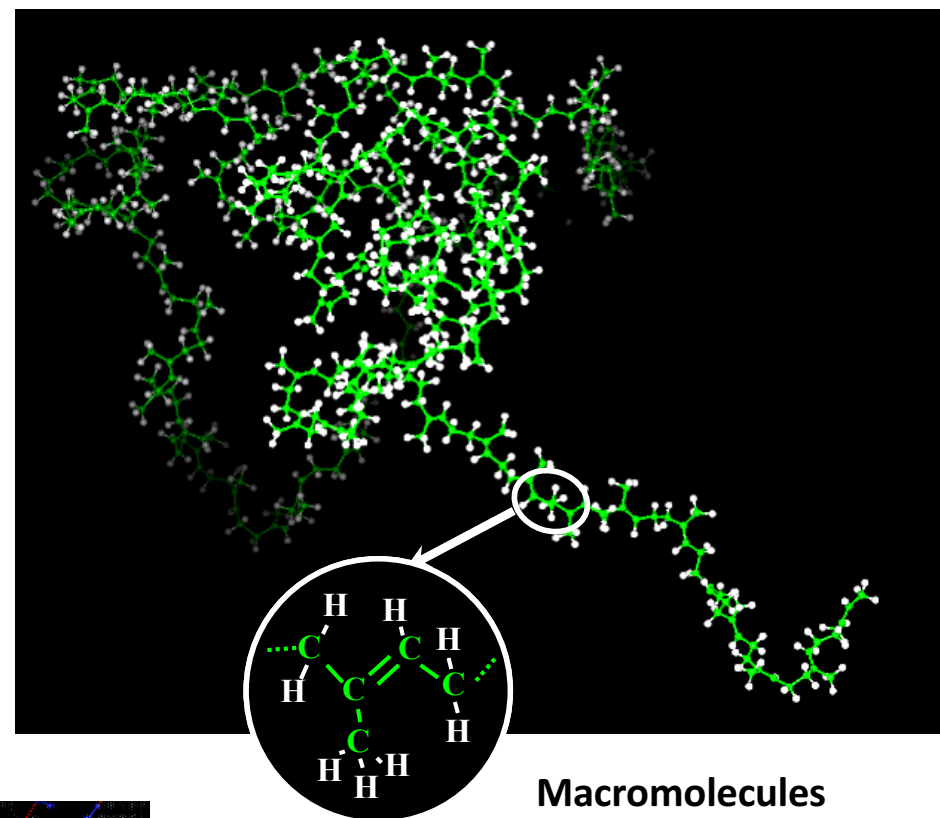
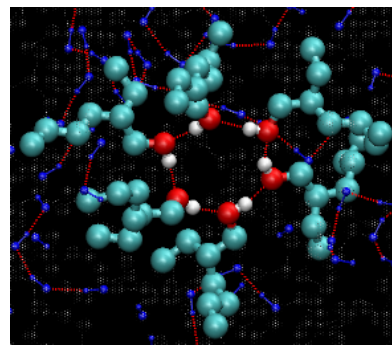
Aromatic



Alkyl chain



alcohol

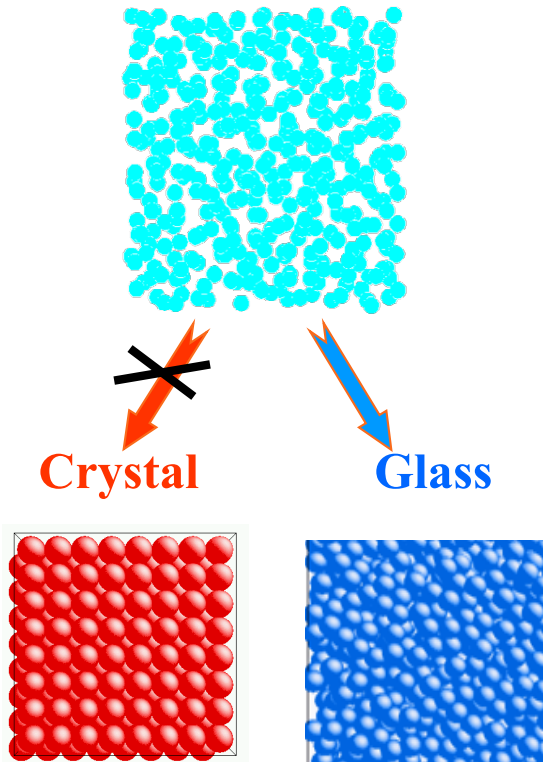


Macromolecules
Repetition of
N Monomers

A minimal complexity is required :

Argon does not form a glass , neither binary Lennard-Jones systems

avoiding the crystallisation
with locally preferred structures



➤ **Intramolecular** Covalent bonds
200-600 kJ/mole

➤ **Intermolecular** interactions :
Hydrogen bond 1-40kJ/mole
L-J ~10 kJ/mole
coulombic

No network

Molecules Low Tg = 56K- 350K

Polymers Tg = 150K 470K

Empirical rules :

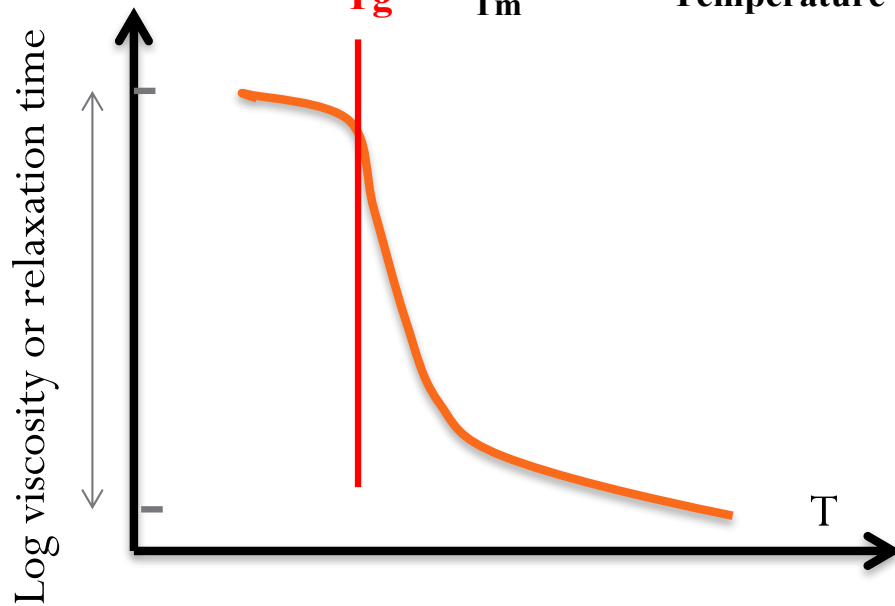
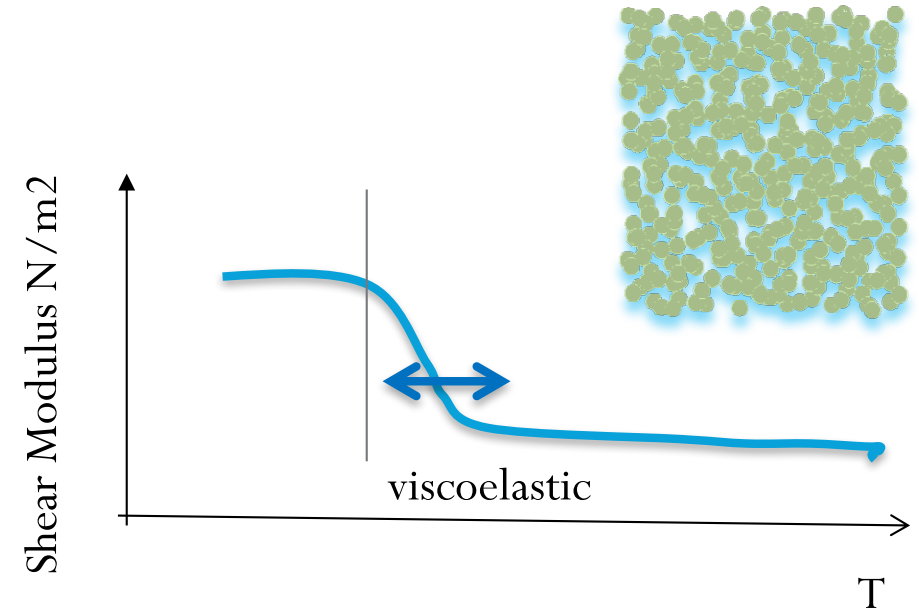
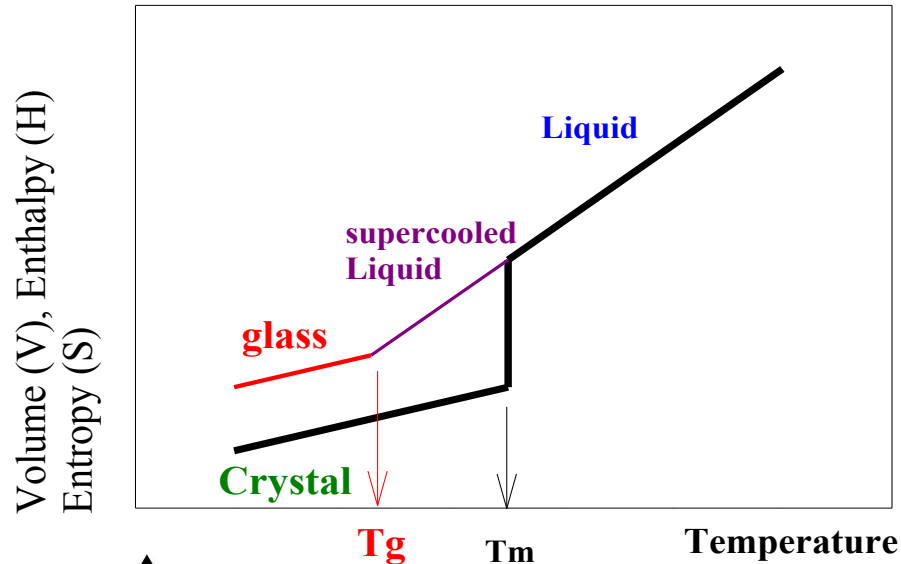
T_b / T_g > 2

T_g / T_m ~ 2/3

Glass formation = solidification without crystallization
Quench easy, sometimes in liq N₂

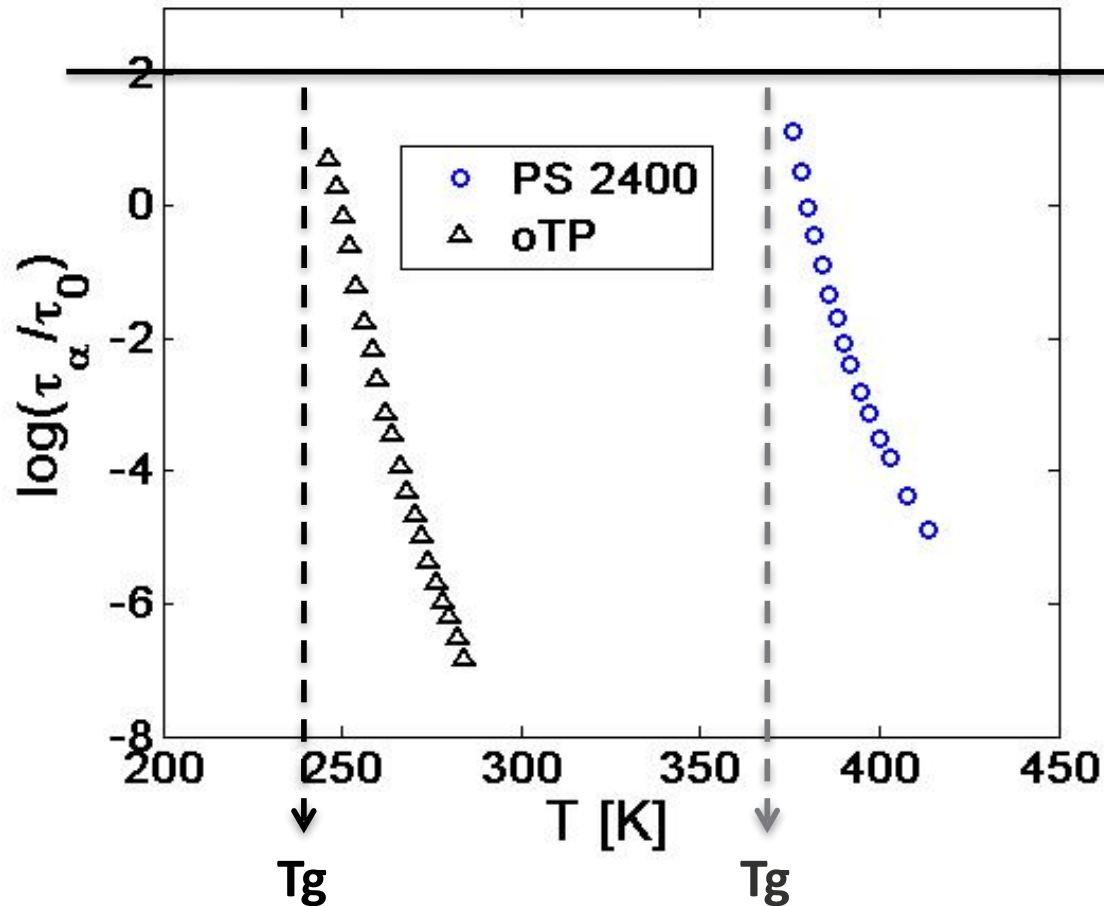
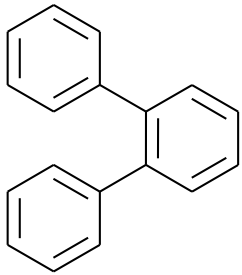
at atmospheric pressure

Looking for an universal relationship between thermodynamics, structure and dynamics

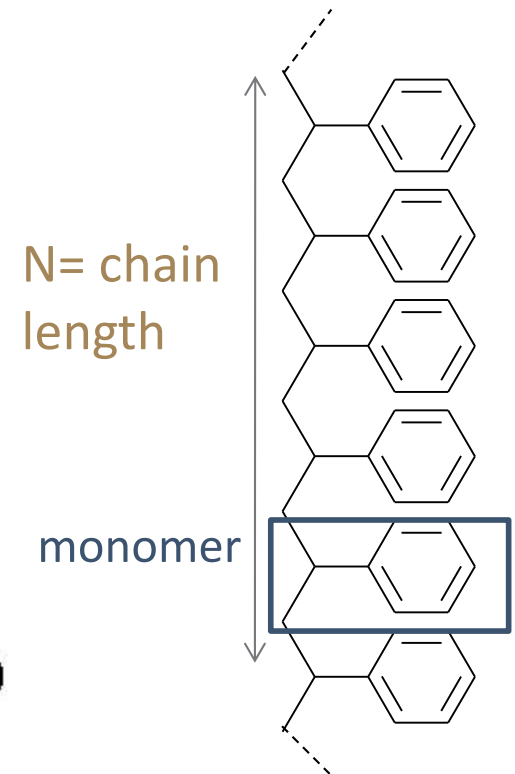


Dramatic viscous slowing down

o-terphenyl



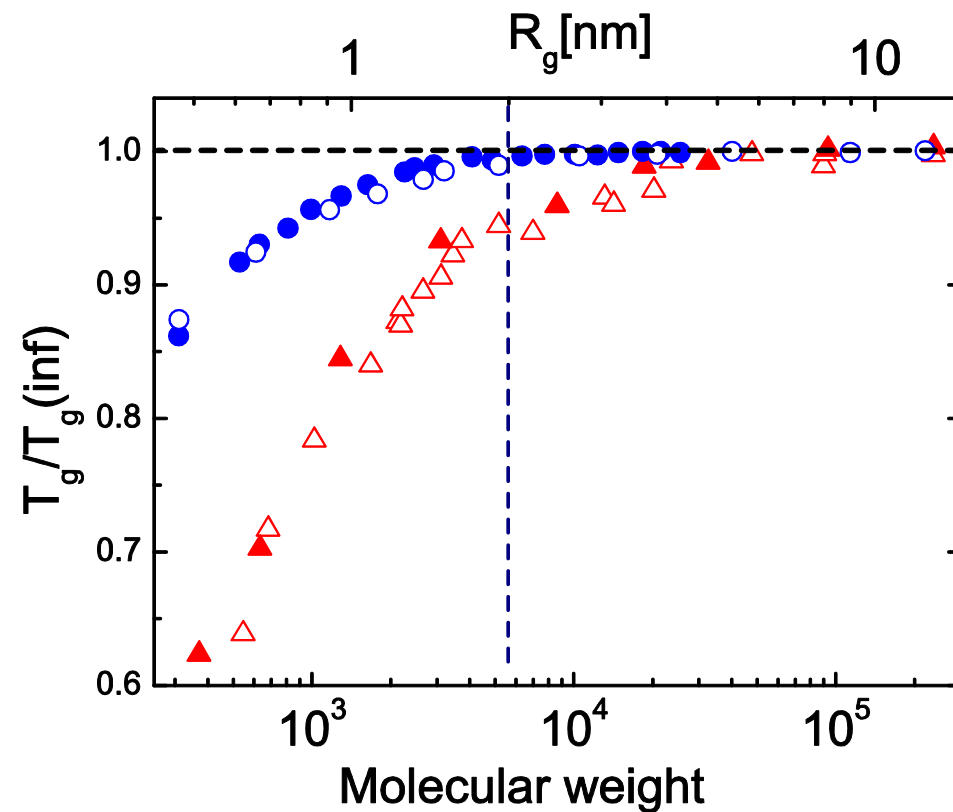
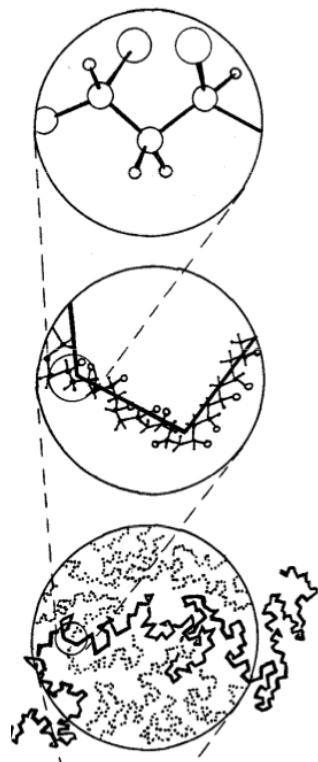
Polystyrene



Dramatic viscous slowing down approaching the glass transition
How sensitive to external parameters?

Polymers : Molecular Weight effect on Tg

$$T_g = A + \frac{B}{M_w} \quad (\text{Flory})$$

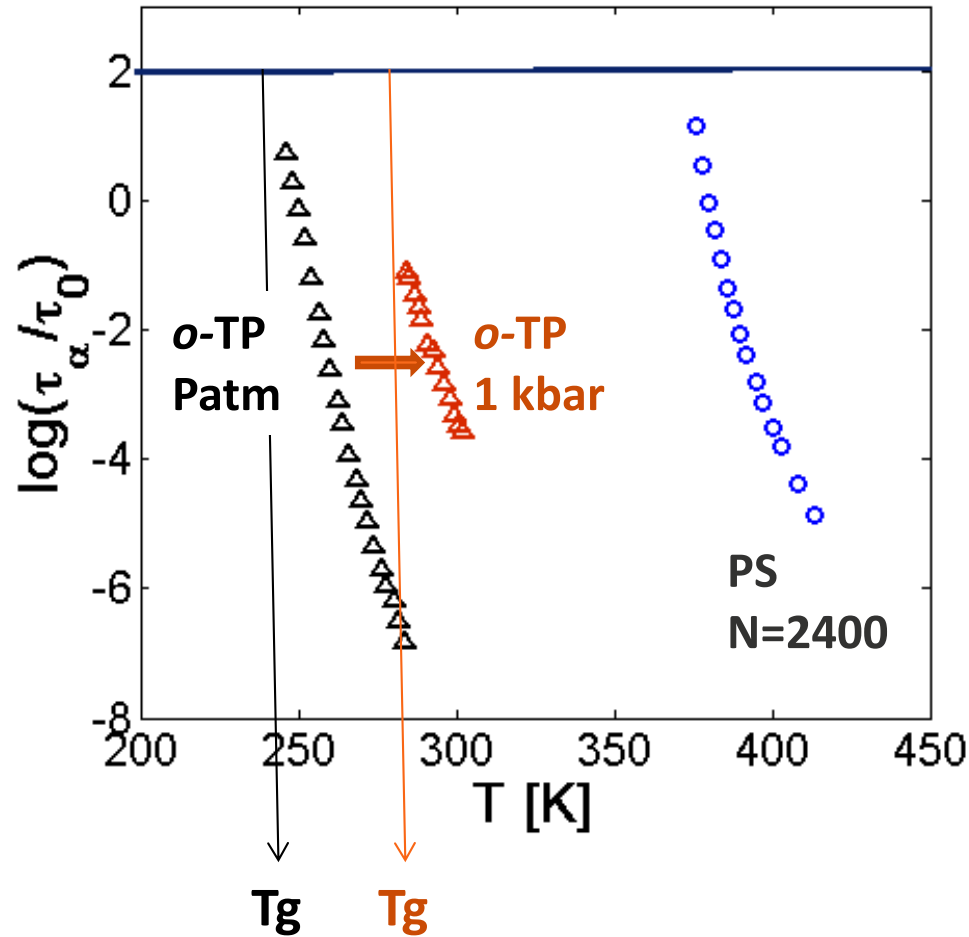


Flexible Polymer PDMS vs rigid polymer PS

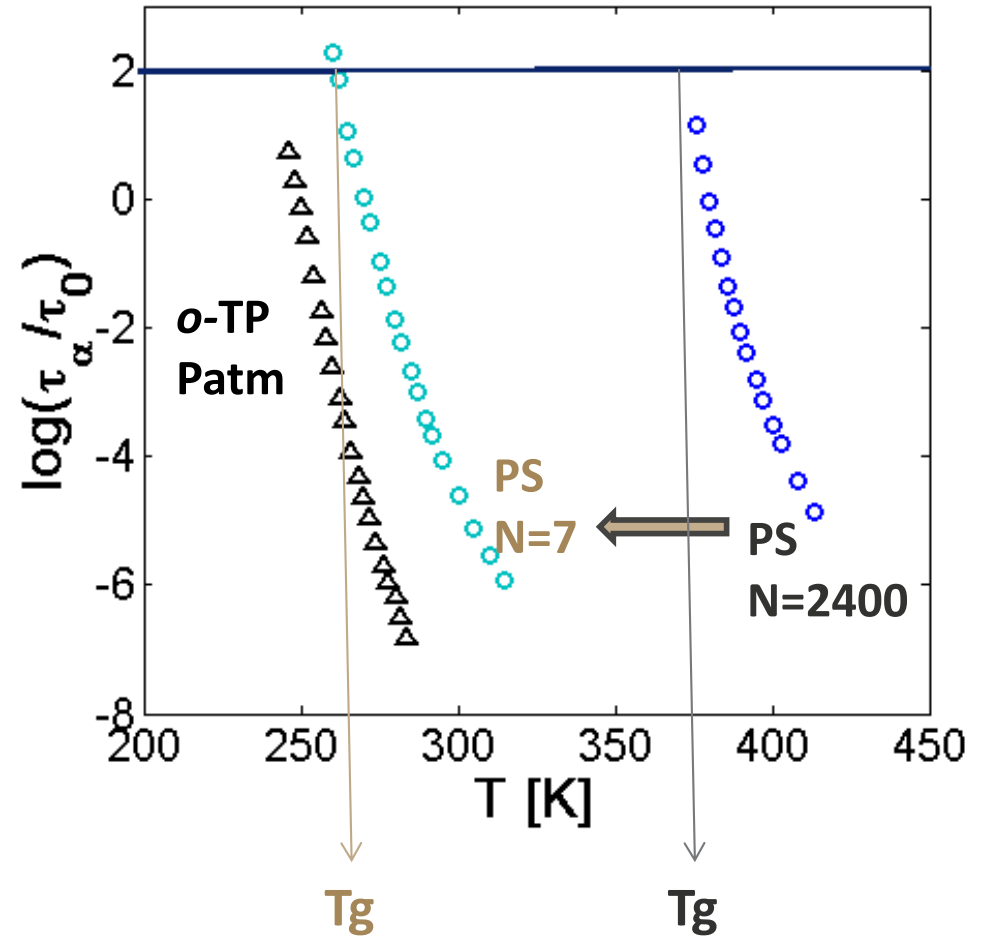
R_g versus M_w

Other control parameters : pressure and chain length

Effect of P



Effect of N

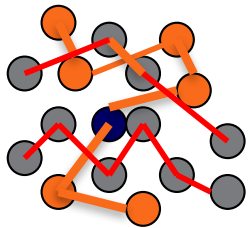
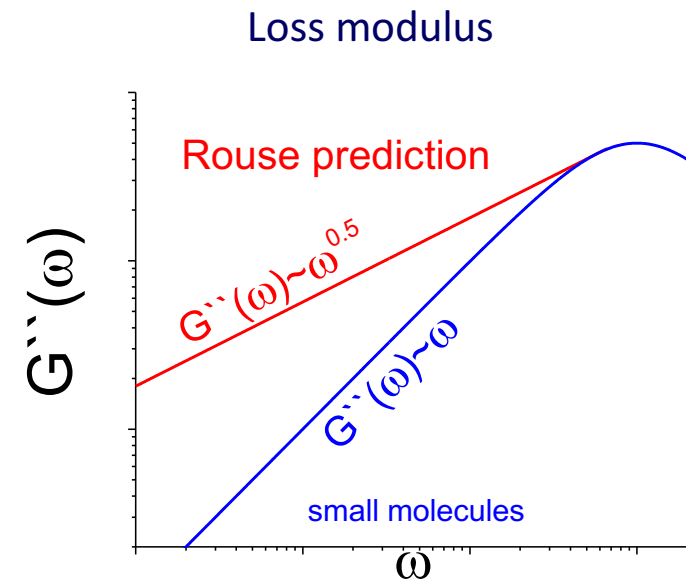
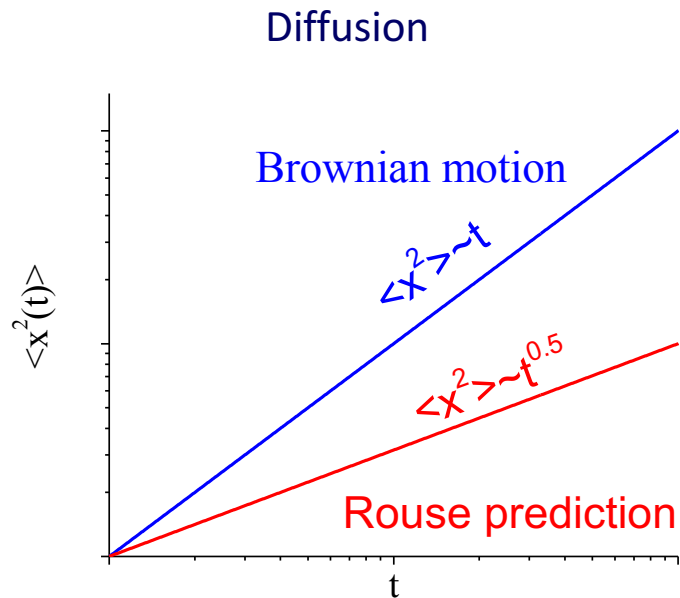
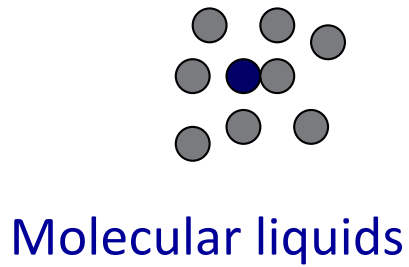


Tg changes with **pressure P** and chain length N



Differences between molecules and polymers

Brownian motion $\xi \frac{dx(t)}{dt} = f(t)$ ξ is a friction, $f(t)$ is a random force.



Polymeric system

Additional force appears in a polymer due to chain connectivity.

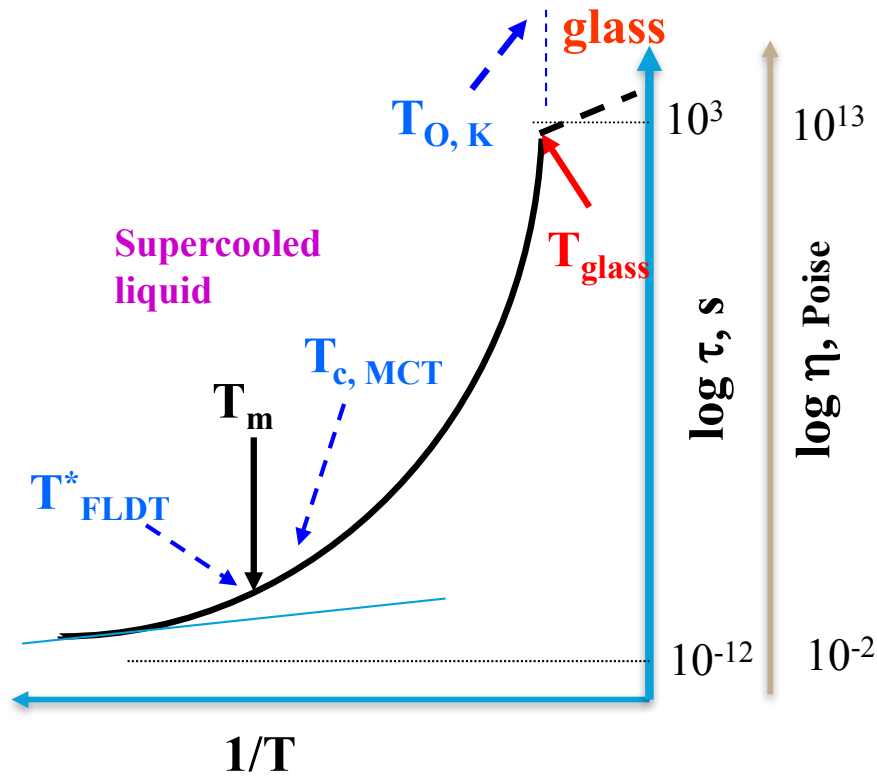
Rouse model assumes Gaussian statistics of the chain and:

It predicts: $\langle x^2 \rangle \sim t^{0.5}$ and $G''(\omega) \sim \omega^{0.5}$

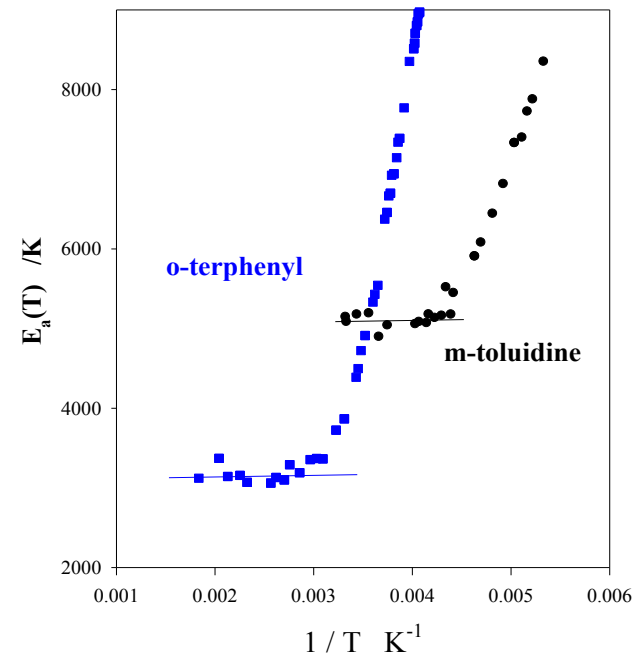
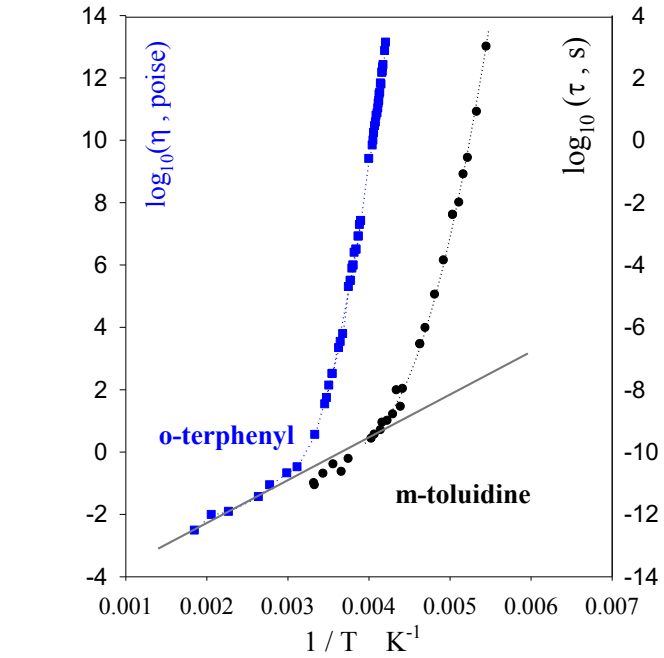
$$\xi \frac{dx(t, n)}{dt} = \frac{3T}{a^2} \frac{\partial^2 x(t, n)}{\partial^2 n} + f(t)$$

Viscous slowing down and dynamical arrest thermally activated kinetics

$$\tau = \tau_{\infty} \exp [E_a(T) / k_B T]$$



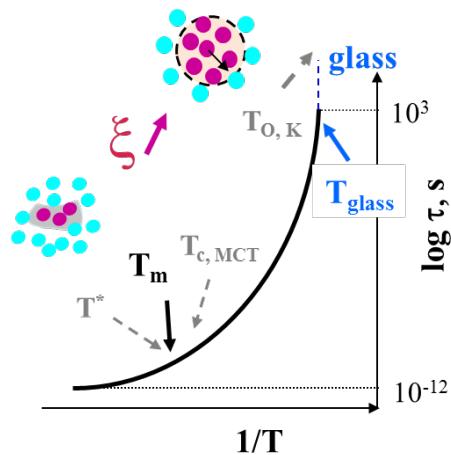
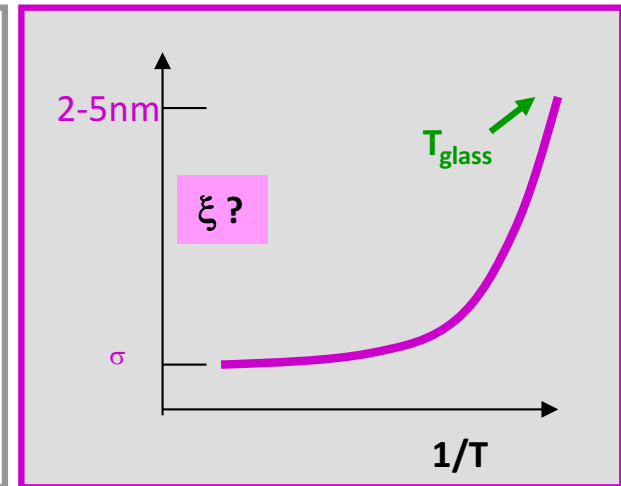
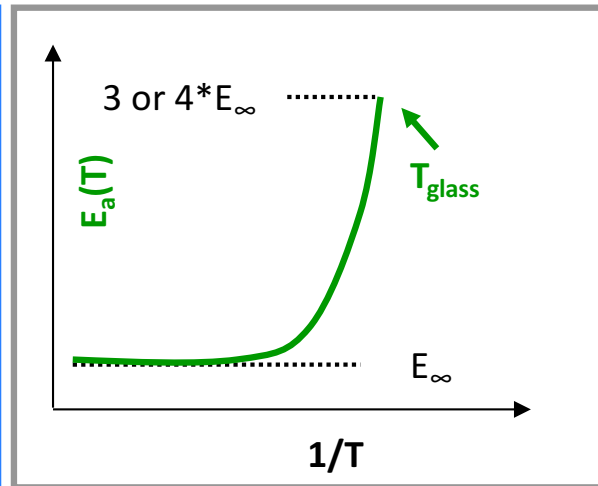
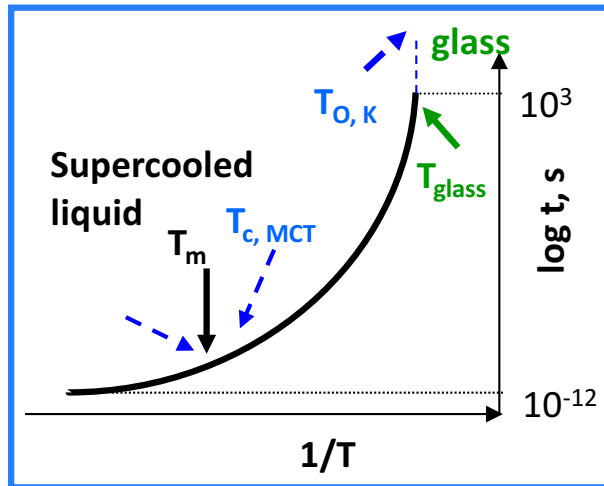
Increase of an effective activation energy
 → **collective behavior**
 suggests an **heterogeneous dynamics**



Central question :

Is the sharp increase of the viscosity related to the growth of a 'cooperative' length scale ?

$$\tau = \tau_{\infty} \exp [E_a(T) / k_B T] \quad \text{and} \quad E_a(T) \propto \xi(T)?$$



Glass Science : a jungle of phenomenology

Summarized by a **quest for characteristic lengthscale(s)**

provided by dynamical properties

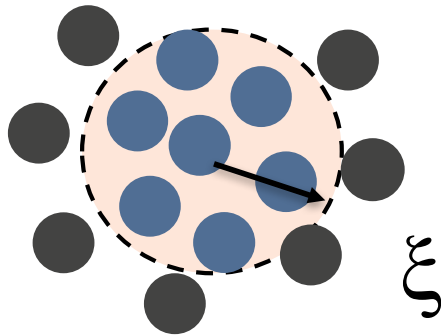


Cooperativity and correlation length

Cooperativity length static

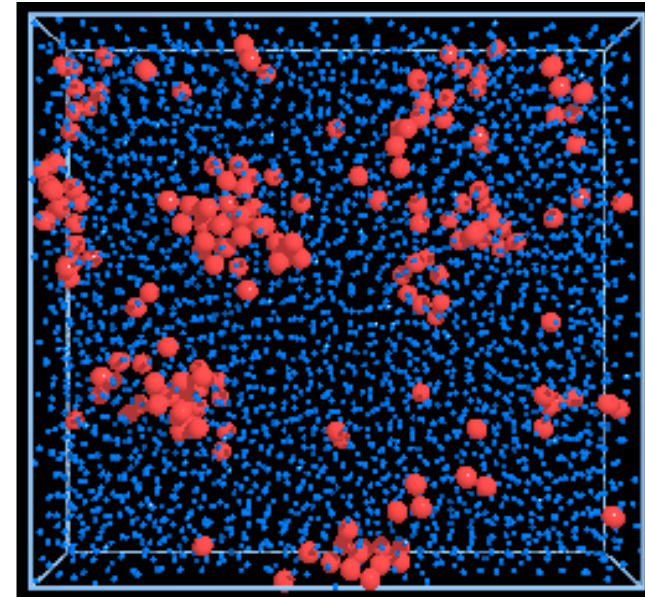
$$\tau(T) = \tau_0 \exp \left[N_{\text{coop}}(T) \frac{\Delta\mu}{k_B T} \right]$$

$N_{\text{coop}} \nearrow$ as $T \searrow$



Adam and Gibbs, 1965

Correlation length Dynamical



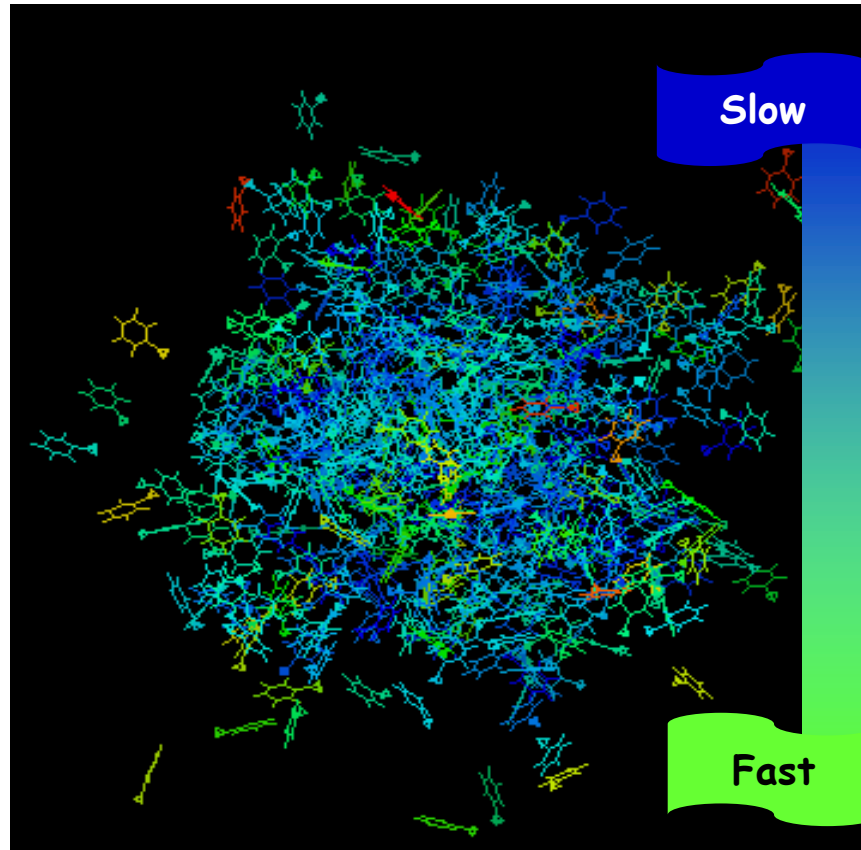
E. Weeks *et al.*, *Science* 287, 627 (2000)
www.physics.emory.edu/~weeks/lab/glass

$\Phi = 56\%$ supercooled fluid
 5% fastest of the particles in red

Large clusters of cooperative fast particles :
 spatially heterogeneous dynamics

Relevant length : dynamic or static ?

Toluene



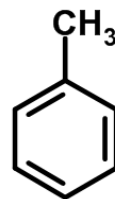
“**Dynamic heterogeneities**”: at low T, liquid made of regions with widely different mobilities

Spatial correlations in the dynamics

From Standard 2-time correlation function

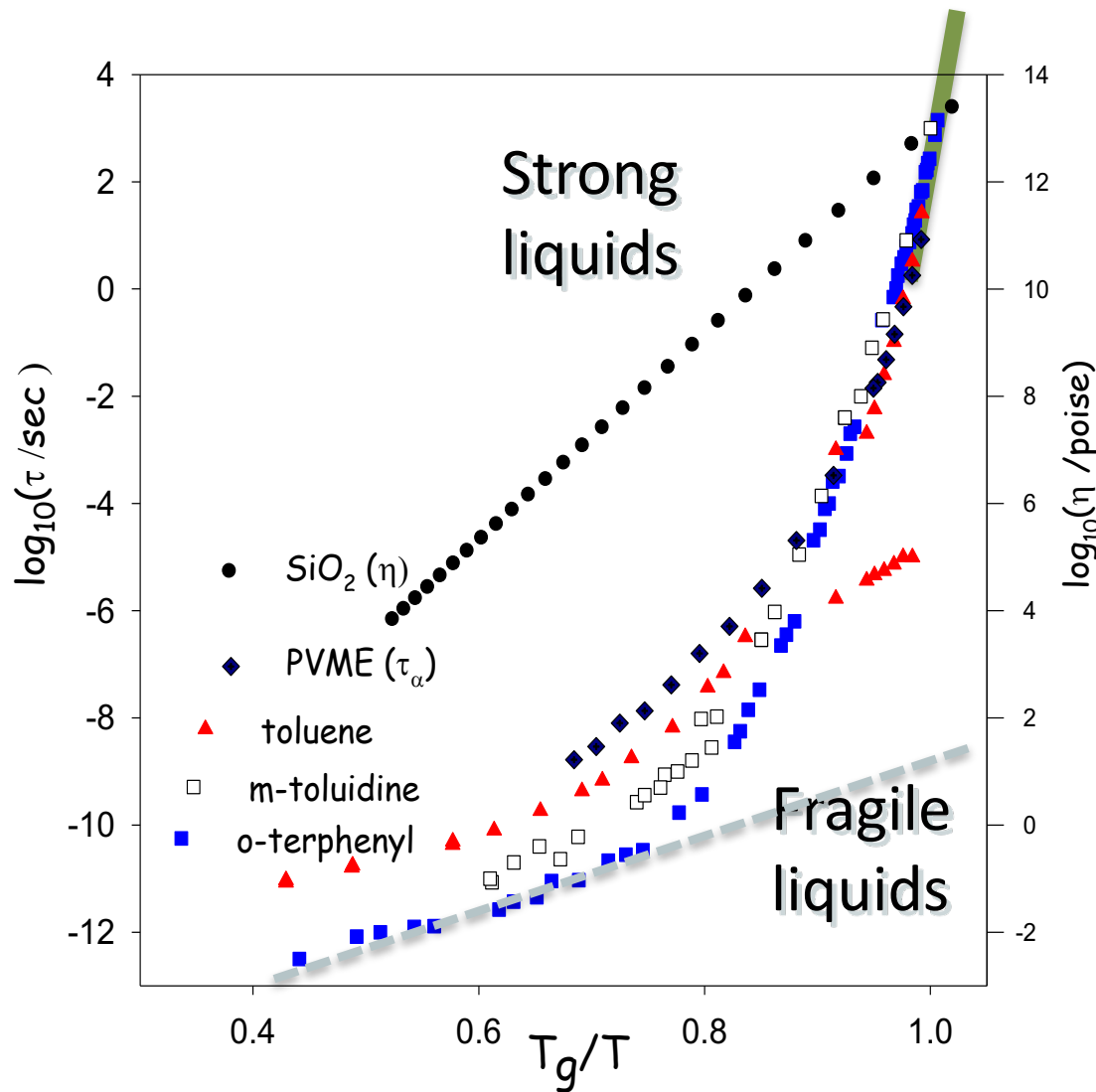
to

Space-time 4-point correlation function



Viscous slowing down and glass transition

1- Notion of fragility



$$m_P = \left. \frac{\partial \log_{10}(\tau)}{\partial T_g/T} \right|_P$$

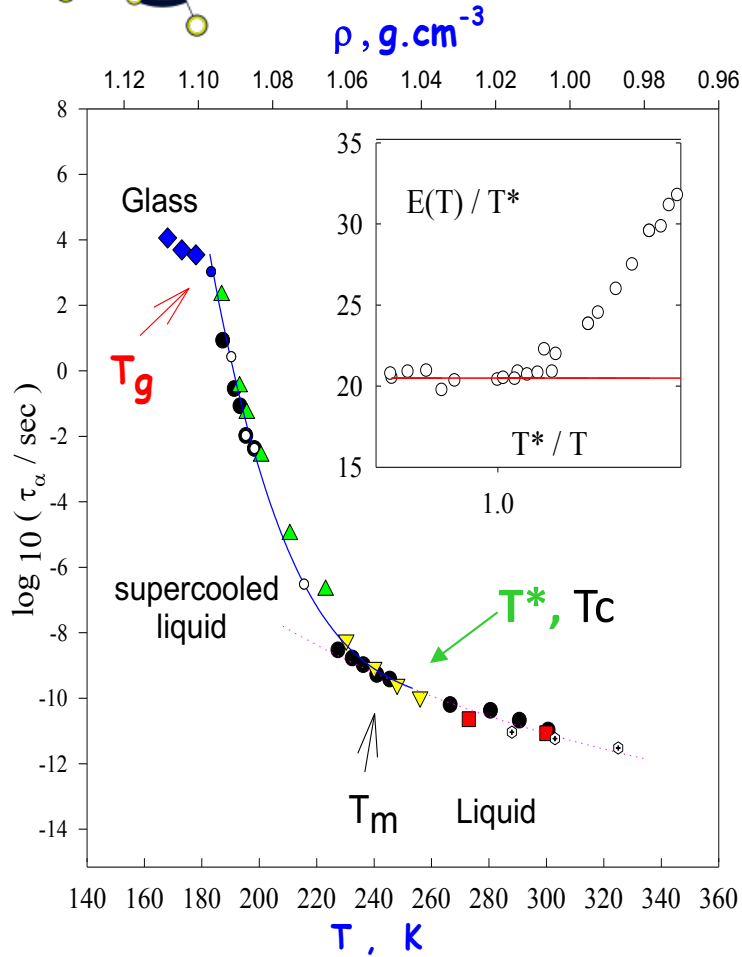
Fragility = the fact that the dynamical slowdown has stronger than Arrhenius temperature dependence

Correlations of m with various properties in the corresponding glass or liquid

Thermal expansivity $6-10 \cdot 10^{-4} \text{ K}^{-1}$

Angell's classification (1985) :
comparing systems different T-dependence of τ_α or η

How to scale out the density at atmospheric pressure ?



m-toluidine

$$\left(\frac{\partial \ln \tau}{\partial T}\right)_P = \underbrace{\left(\frac{\partial \ln \tau}{\partial T}\right)_\rho}_{\text{due to } T} + \underbrace{\left(\frac{\partial \ln \tau}{\partial \rho}\right)_T}_{\text{due to } \rho} \left(\frac{\partial \rho}{\partial T}\right)_P$$

$$= - \left(\frac{\partial \ln \tau}{\partial \ln \rho}\right)_T (-\alpha_\tau + \alpha_P)$$

isobaric expansivity

$$\alpha_P = -\rho^{-1} \left(\frac{\partial \rho}{\partial T}\right)_P$$

isochronic expansivity
(cst τ or η)

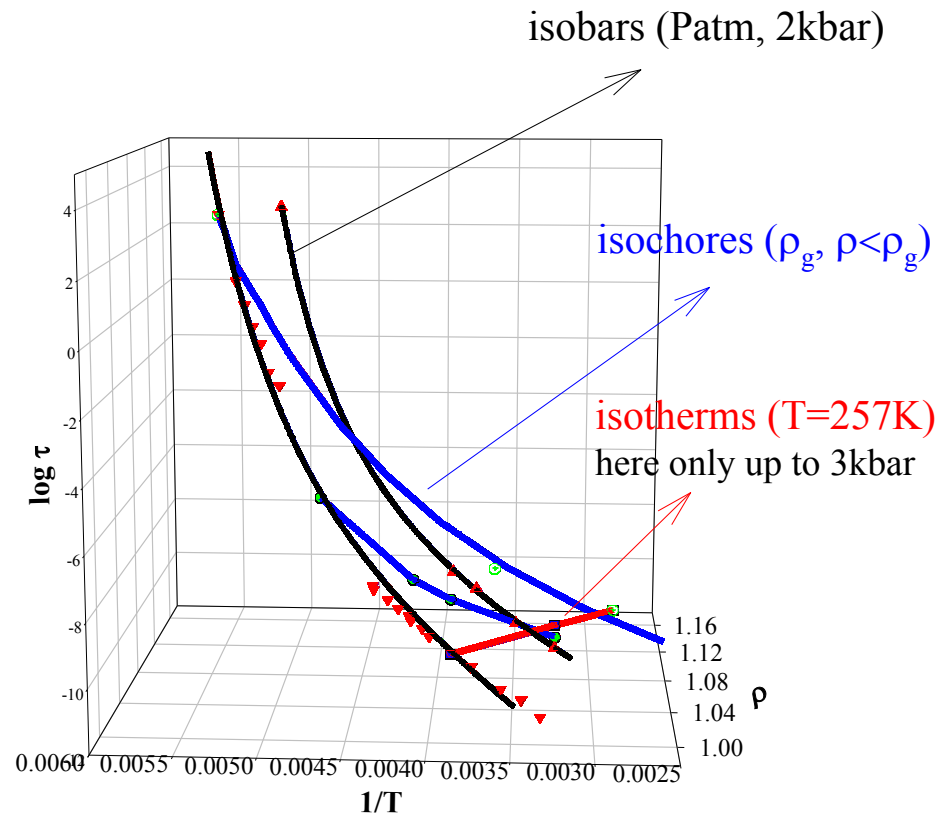
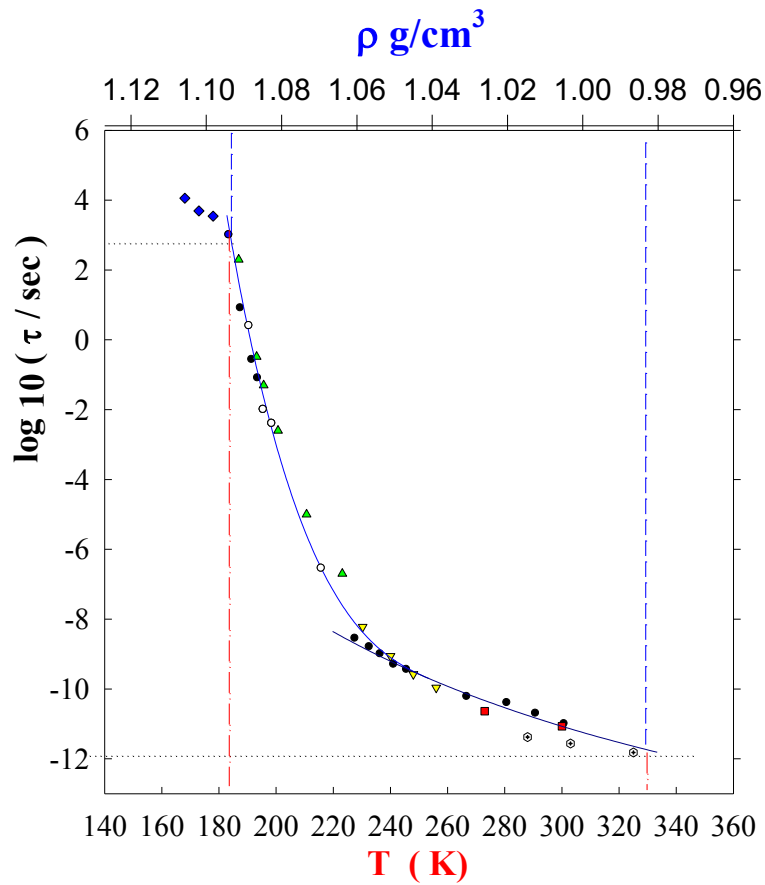
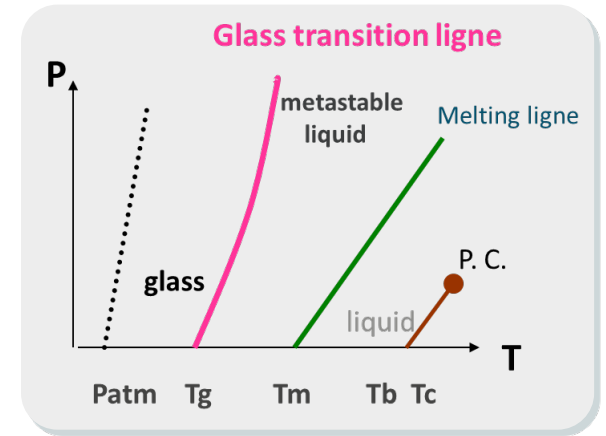
$$\alpha_\tau = -\rho^{-1} \left(\frac{\partial \rho}{\partial T}\right)_\tau$$

$\frac{ \alpha_\tau }{\alpha_P}$	the relative importance of density and temperature
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Suggesting high Pressure (or ρ) experiments

Molecular liquid m Toluidine

Patm high pressure



$$\left(\frac{\partial \ln \tau}{\partial T}\right)_P = \left(\frac{\partial \ln \tau}{\partial T}\right)_\rho + \left(\frac{\partial \ln \tau}{\partial \rho}\right)_T \left(\frac{\partial \rho}{\partial T}\right)_P$$

$$\left(\frac{\partial \ln \tau}{\partial T}\right)_P = \left(\frac{\partial \ln \tau}{\partial \rho}\right)_T (-\alpha_\tau + \alpha_P)$$

isobaric thermal expansivity coefficient

$$\alpha_P = -\rho^{-1} \left(\frac{\partial \rho}{\partial T}\right)_P$$

isochronic expansivity coefficient (at η or τ constant)

$$\alpha_\tau = -\rho^{-1} \left(\frac{\partial \rho}{\partial T}\right)_\tau$$

ratio $|\alpha_\tau|/\alpha_P$ = quantitative estimation of the relative contributions of the temperature and the density to the change of the relaxation time along a given isobar

	isochrone	$\alpha_P \cdot 10^4 \text{K}^{-1}$	$-\alpha_\tau \cdot 10^4 \text{K}^{-1}$	$ \alpha_\tau /\alpha_P$
polybutadiene	$\tau=10^{-9}\text{sec}$	6.5	10.6	1.6
glycerol	$\tau=10^{-9}\text{sec}$	4.2	25	6.2
m-toluidine	$\tau=10^{-9}\text{sec}$	7.3	11.5	1.6

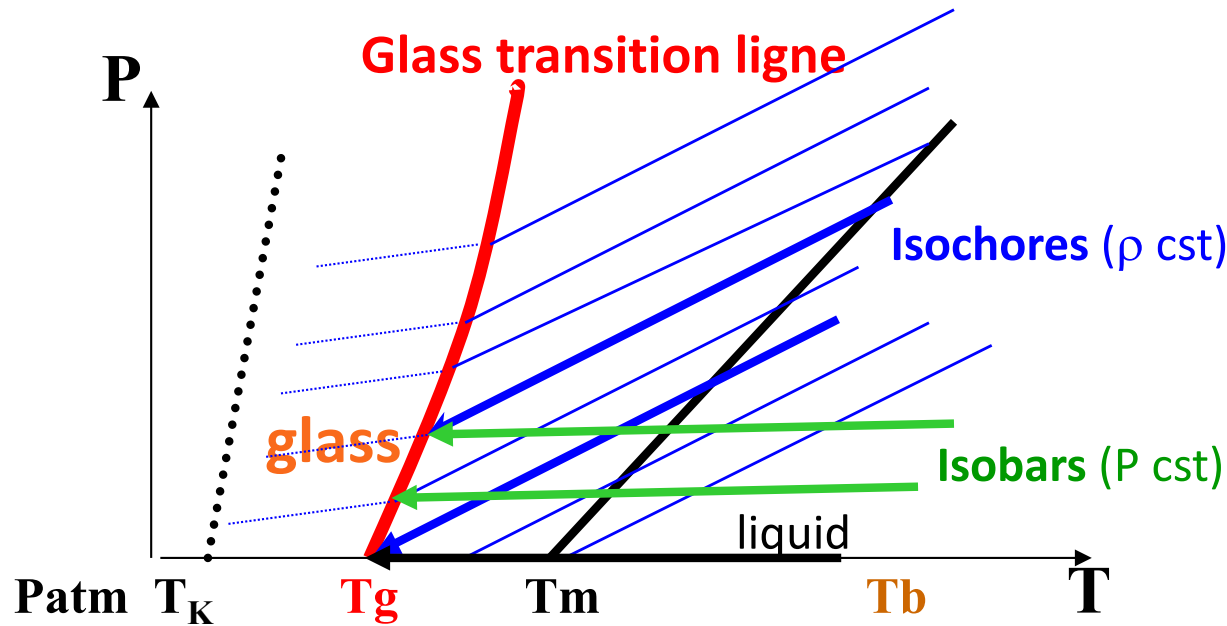
$$|\alpha_\tau|/\alpha_P > 1$$

T is the driving parameter of the viscous slowing down!

Test correlation to fragility by changing

- a) -the thermodynamic state or path
- b) -the relaxation time

A property is then considered as “**intrinsic**” when it does not depend on a) and b)



$$m_{\rho} = \left. \frac{\partial \log_{10}(\tau_{\alpha})}{\partial T_g(\rho)/T} \right|_{T_g(\rho)}$$

$$m_P = \left. \frac{\partial \log_{10}(\tau_{\alpha})}{\partial T_g(P)/T} \right|_{T_g(P)}$$

Isochoric and isobaric fragilities $m_{\rho, \tau}$ and $m_{P, \tau}$
related through

$$m_{P, \tau}(P) = m_{\rho, \tau}(1 + \alpha_P / |\alpha_{\tau}|) \Rightarrow m_{\rho, \tau} \leq m_{P, \tau}$$

$$\alpha_{\tau} = -\rho^{-1} \left(\frac{\partial \rho}{\partial T} \right)_{\tau}$$

$$\alpha_P = -\rho^{-1} \left(\frac{\partial \rho}{\partial T} \right)_P$$

Ex : o-TP $m_{P, g} = 82$, while $m_{\rho, g} = 45$



model-free assessment of the respective contributions of ρ and T in the viscous slowing down at P cst

$$\tau_{\alpha}(\mathbf{T}, \rho) = \tau_{\infty} \exp \left[\frac{E_a(\mathbf{T}, \rho)}{k_B T} \right]$$

two extremes physical pictures

dynamic congestion due increase of density and drainage of free volume

activated dynamics ($E_a > kT$)
thermally activated processes
in cst density energy landscape



$$\log (\tau_{\alpha}(\rho, T)) = F \left[\frac{E_{\infty}(\rho)}{T} \right]$$

ρ and T dependences of the α -relaxation time
described by a single scaling variable $X = E_{\infty}(\rho) / T$

$E_{\infty}(\rho)$ = effective activation energy
characteristic of high- T Arrhenius regime

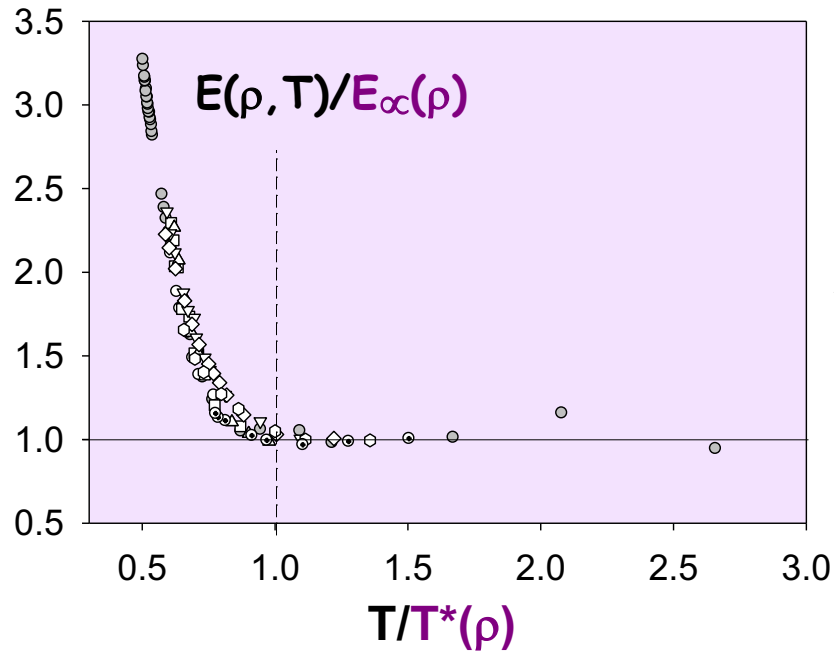
One single energy scale !

How does the scaling work ? How well established ? on the $E(\rho, T)$ or on $\log(\tau)$

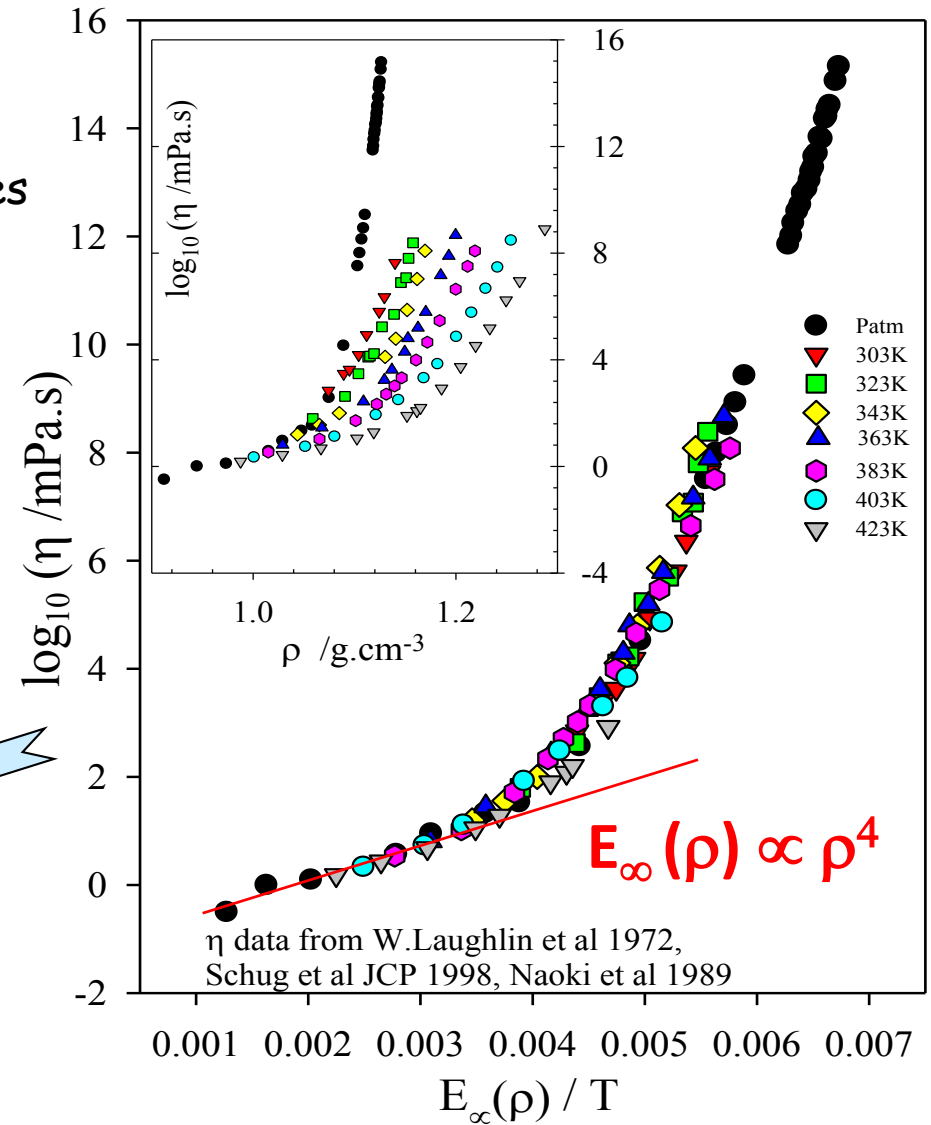
reminiscent of soft spheres
scaling for dynamic and thermodynamic properties

$$\Gamma \rightarrow \Gamma^{-1}$$

$$\rho T^{-1/4} \rightarrow \rho^4 / T$$



A crossover temperature $T^*(\rho) \propto E_{\infty}(\rho)$

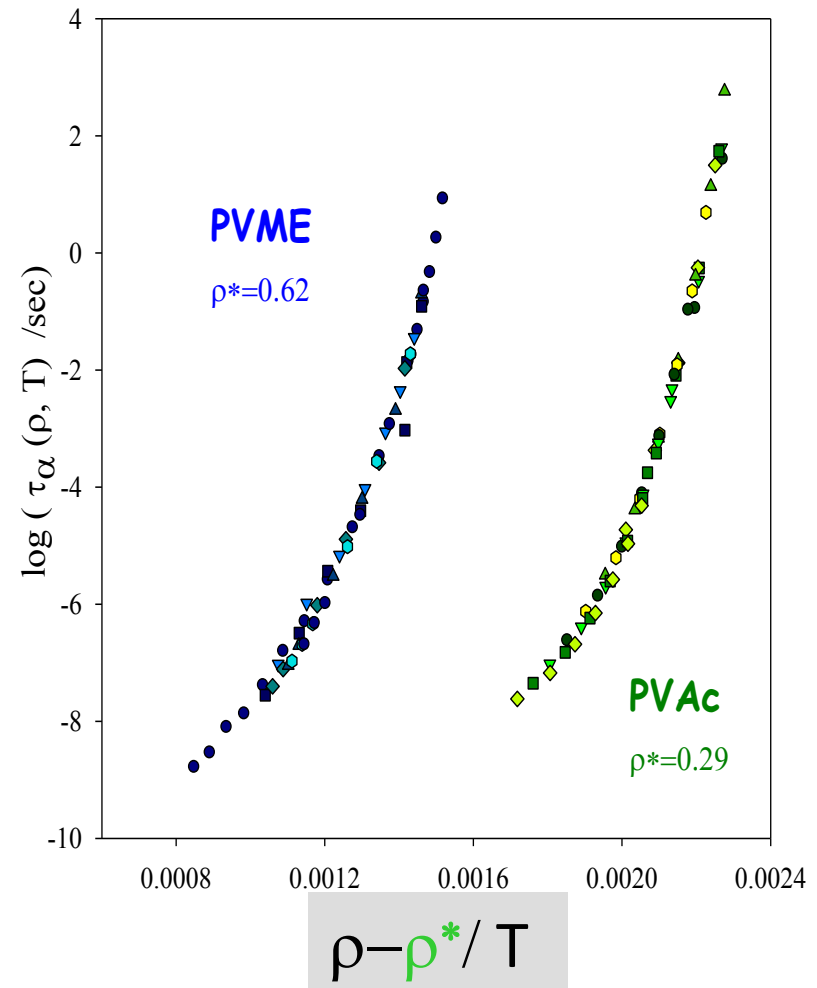
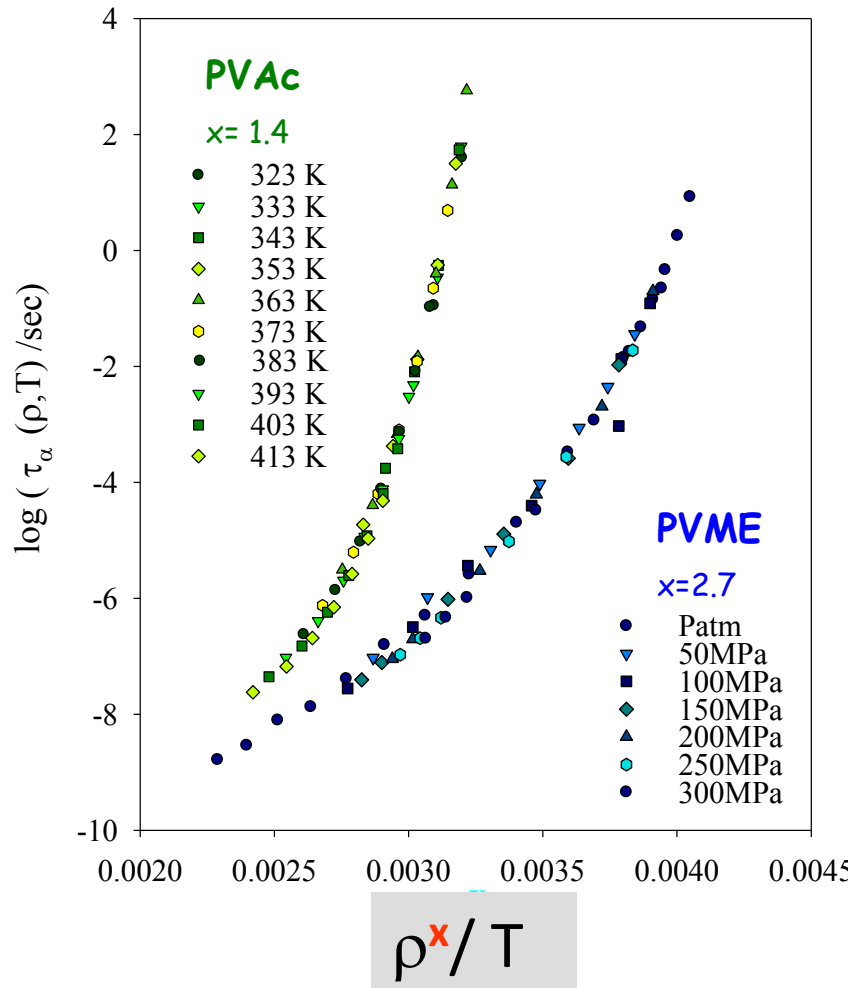


How well established ?

When **no high-T data** are available ,

extension to **polymers** with $e(\rho) \propto E_{\infty}(\rho)$

$$\log(\tau_{\alpha}(\rho, T)) = F\left(\frac{e(\rho)}{T}\right) \text{ where } e(\rho) = \rho^x \text{ or } e(\rho) = \rho - \rho^*$$





How useful?

SCALING plot at the G.T. temperature (100sec),
a modified Angell plot

at cst P (P_{atm}),
log(τ_α) vs T_g/T



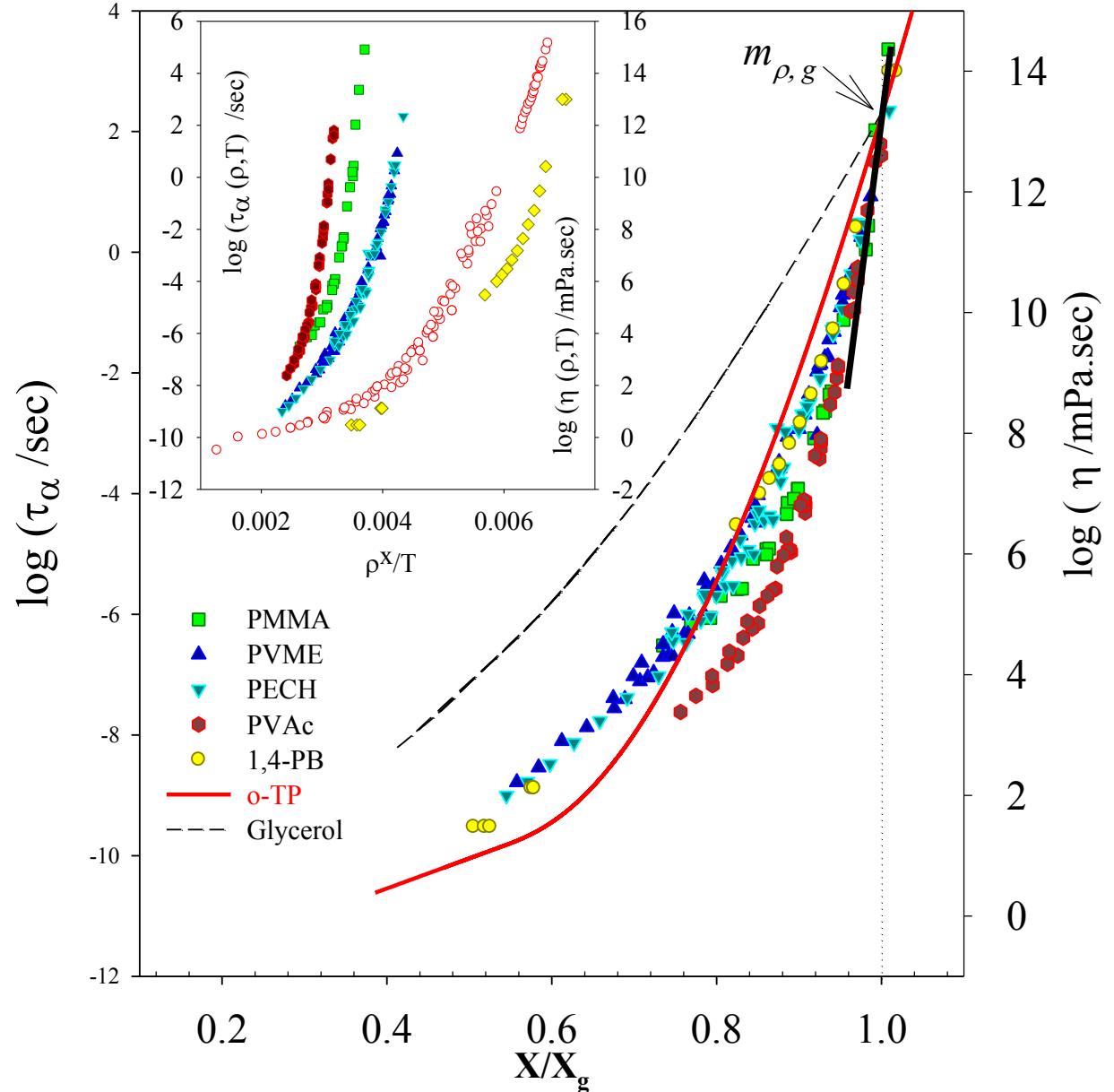
in a wide P-T range
log(τ_α) vs X/X_g

with X= e(ρ)/T
new scaling variable
X_g for τ=100sec

Fragility index ?

$$m = [d \log(\tau) / d(T/T_g)]_{T=T_g}$$

$$m_\rho = \left. \frac{\partial \ln[\tau_\alpha(\rho, T)]}{\partial [T_g(\rho)/T]} \right|_{T_g(\rho)}$$



III - local order and mesoscopic order underlying length scales high T activation energy ultra stable glasses

Looking for a simple* molecular system :

- - large number of partial contributions to $S(Q)$
simplifying the shape?
- ➔ - still able to form a (laboratory) glass and keep some disorder.
(resisting to the crystallization).
- ⊙ - low high Temperature activation energy
and high fragility?

* Simple = pseudo atomic ?

$C_{10}H_{18}$ Decalin



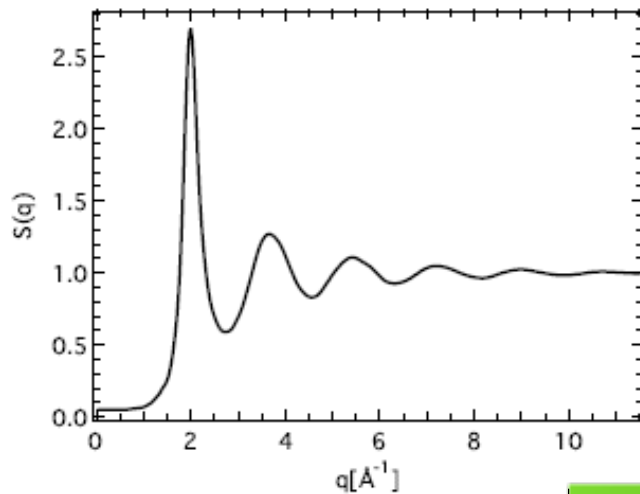
Structure of simple atomic liquids

(here assimilated to a sphere, no additional degrees of freedom)

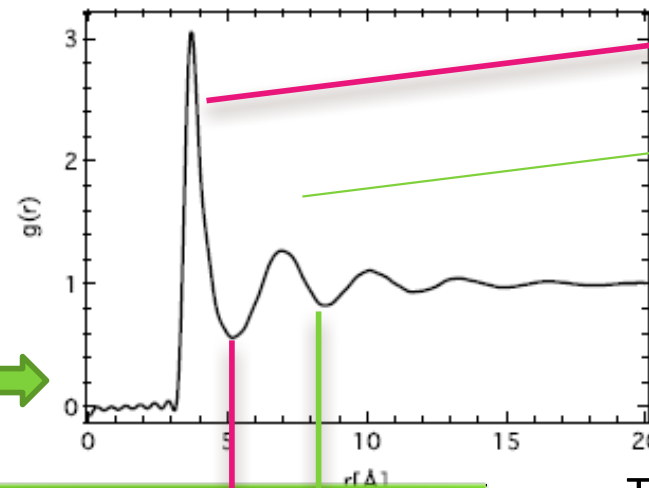
neutron measurement in the Q space (F.T.) \Rightarrow the radial pair correlation function in real space $g(r)$

$$S_M(Q) = f_i(Q) + \frac{4\pi}{Q} \rho_M \int (g_L(r) - 1) r \sin(Qr) dr$$

(a) Structure factor of Argon

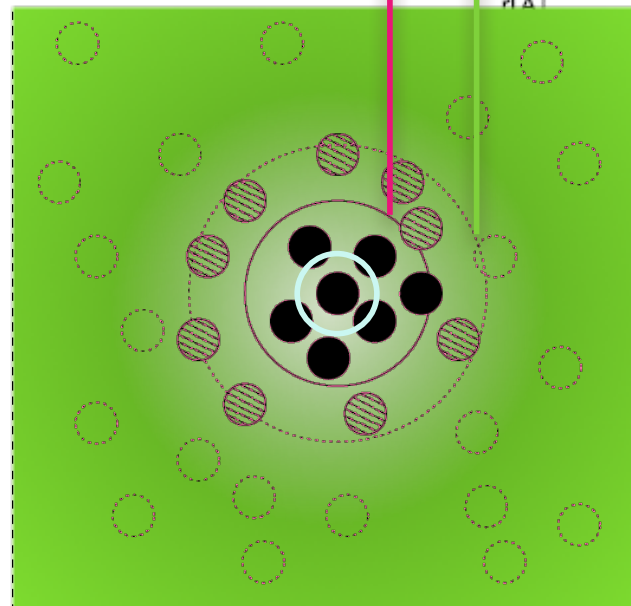


(b) Pair correlation function of Argon



First coordination shell

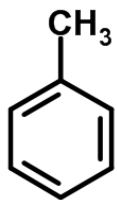
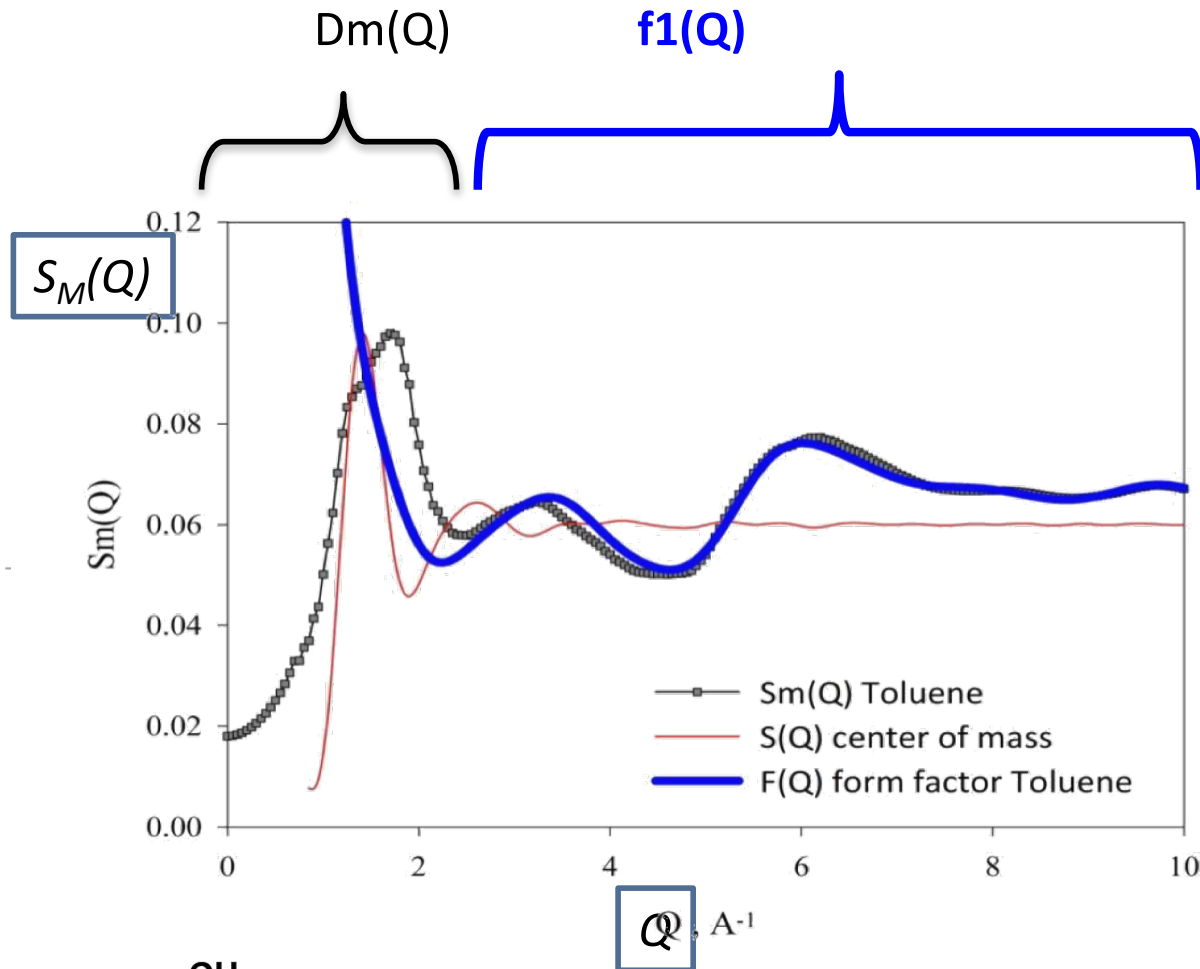
2nd coordination shell



The intermolecular potential governs everything, structure, thermodynamics, dynamics

Subtle local arrangement (bond-orientational order, locally preferred structures).

Decomposition for the structure factor (case of toluene C7D8)



$$S(Q) = Dm(Q) + f1(Q)$$

$$S_M(Q) = f1(Q) + \frac{4\pi}{Q} \rho_M \int (g_L(r) - 1) r \sin(Qr) dr$$

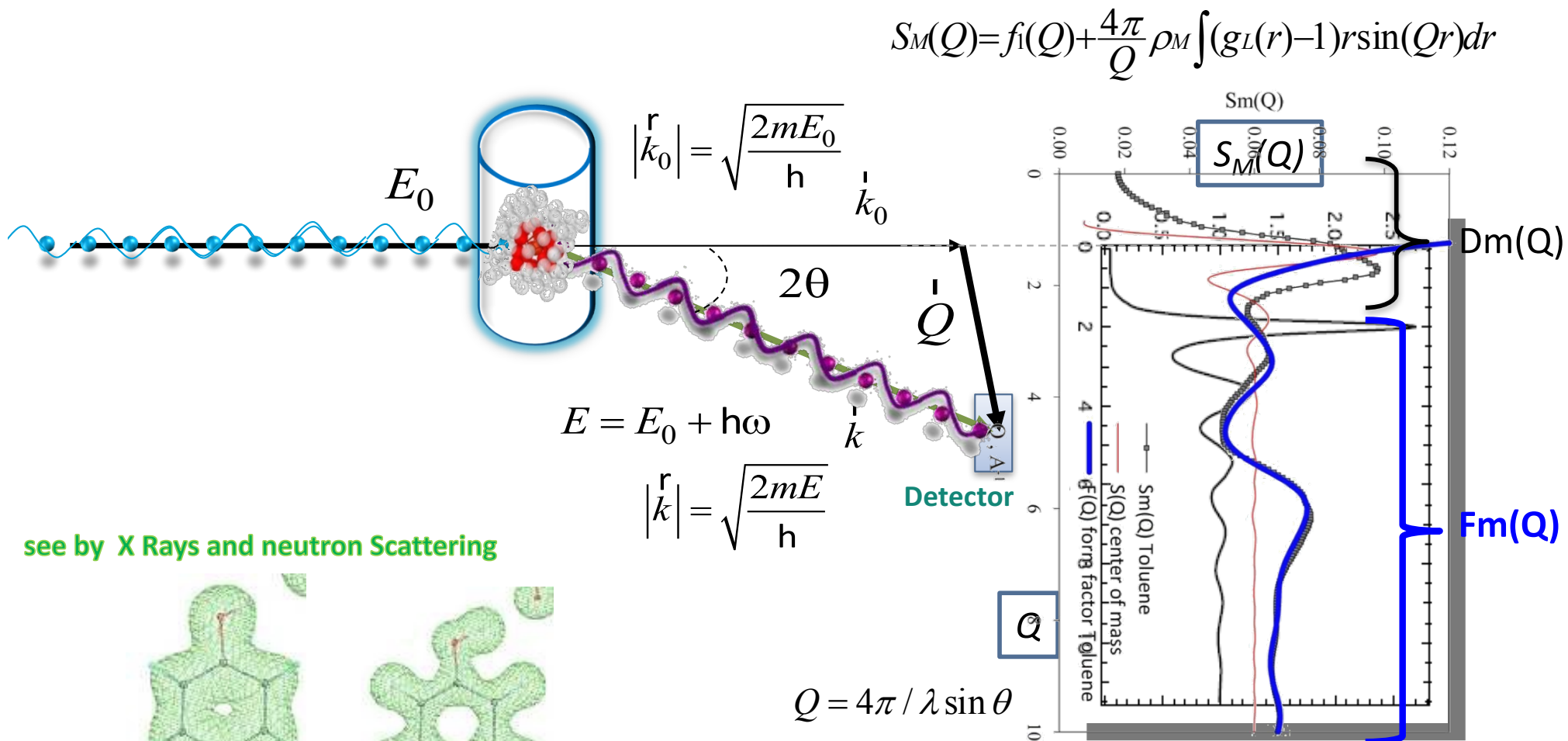
$$S(q) = \frac{1}{\bar{b}^2} \left[\frac{1}{N} \left\langle \left| \sum_i b_i e^{i\mathbf{q} \cdot \mathbf{r}_i} \right|^2 \right\rangle_t - (\bar{b}^2 - \bar{b}^2) \right]$$

$$\bar{b}^2 = \left| \sum_{\alpha} c_{\alpha} b_{\alpha} \right|^2$$

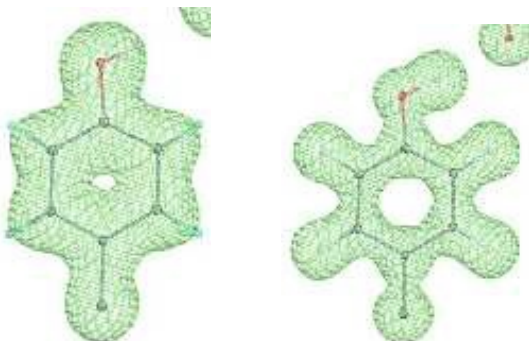
$$\bar{b}^2 = \sum_{\alpha} c_{\alpha} b_{\alpha}^2$$

N = number of atoms in a molecule
 b_i : scattering length of atom i ,
 r_i position vector,
 α, c_{α} atom type and its concentration

Schematic of a diffraction experiment



see by X Rays and neutron Scattering



X-ray (H)

Neutron(D)

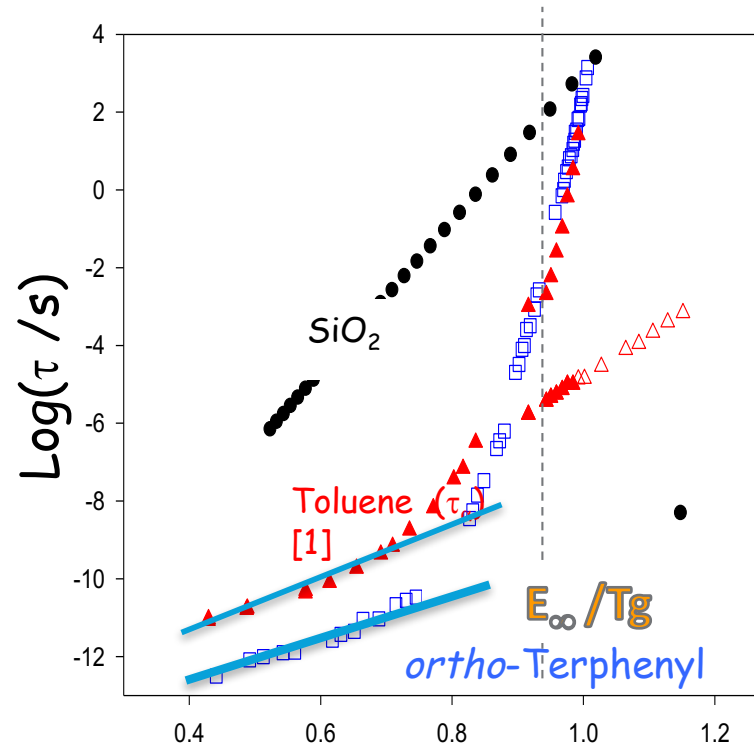
$$S(Q) = D_m(Q) + f_m(Q)$$

Aromatic systems : same intermolecular interaction

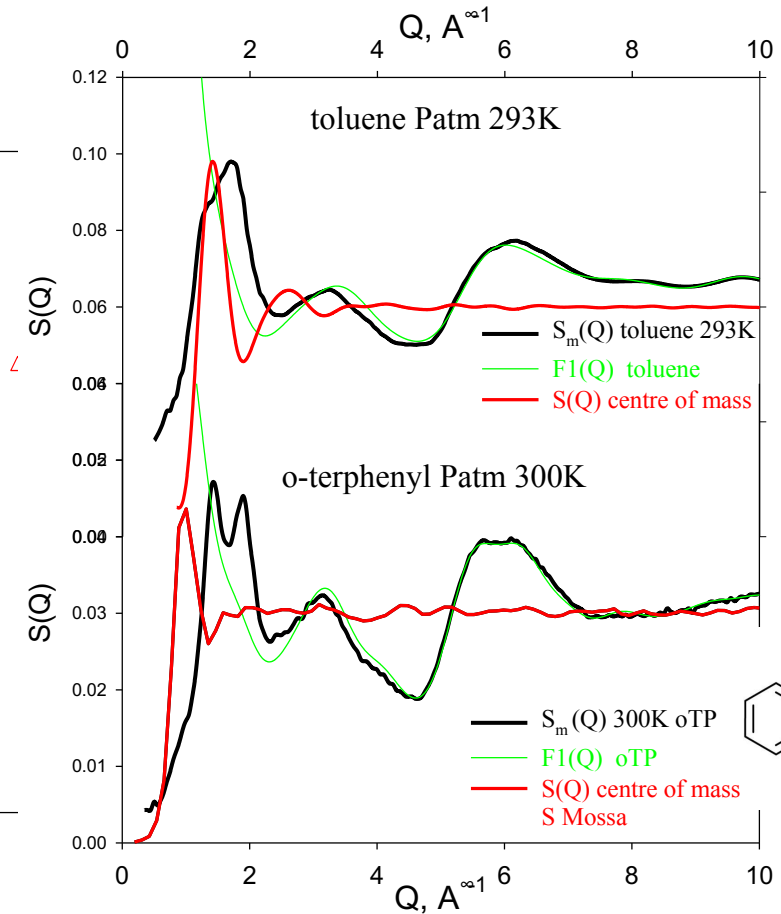
Importance of the High Temperature Activation Energy

Similar fragility index mp
close to T_g ($mp \sim 80-100$)

Tol : Rössler et al



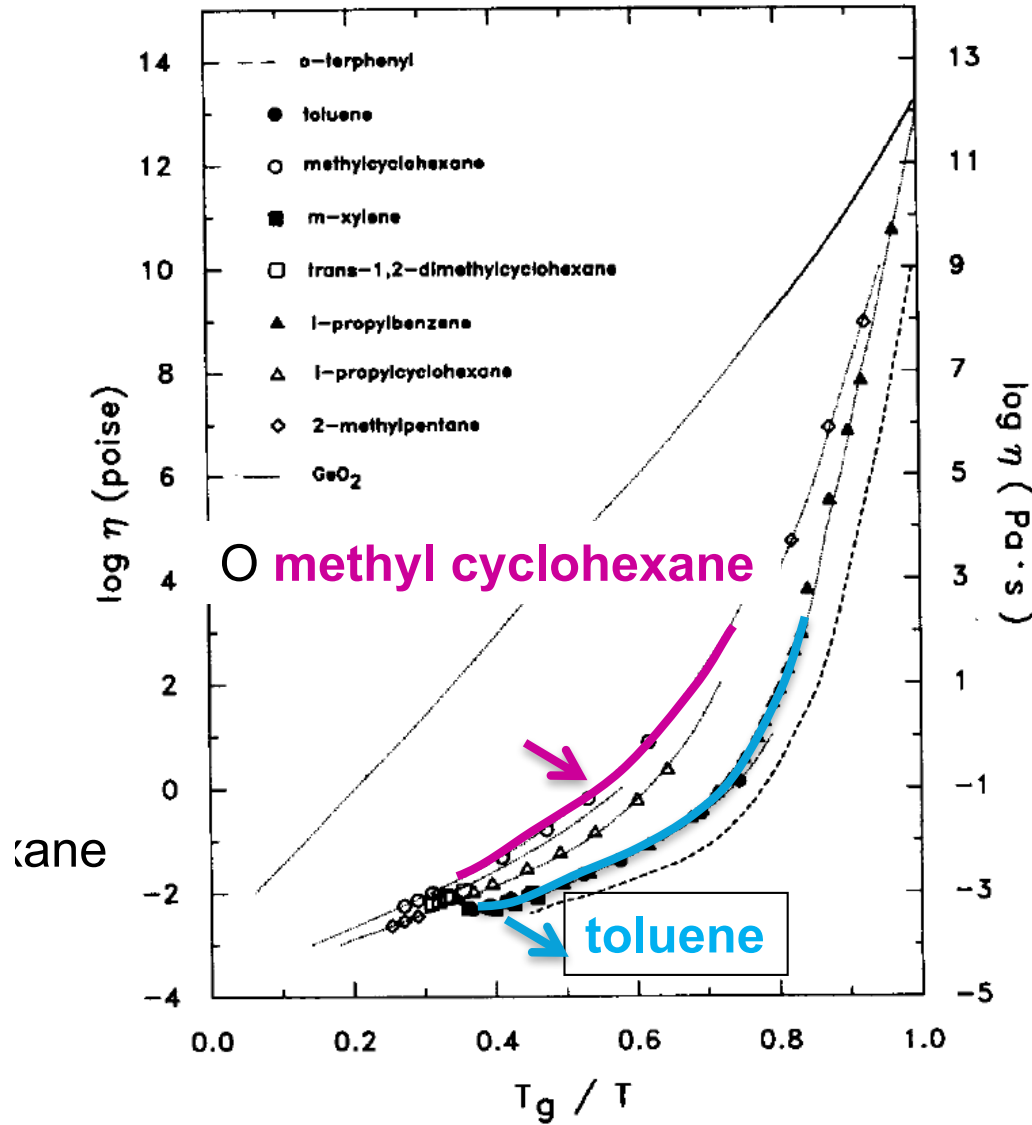
- ➔ Toluene $E_\infty = 1436 \sim 6.8 k_B T_m T_g / T$
- ➔ O-TP $E_\infty = 2431 \sim 7.4 k_B T_m$



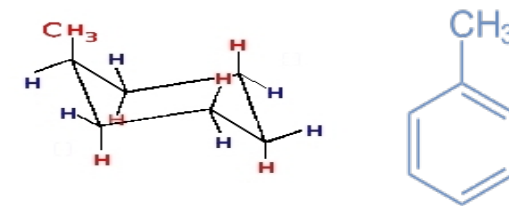
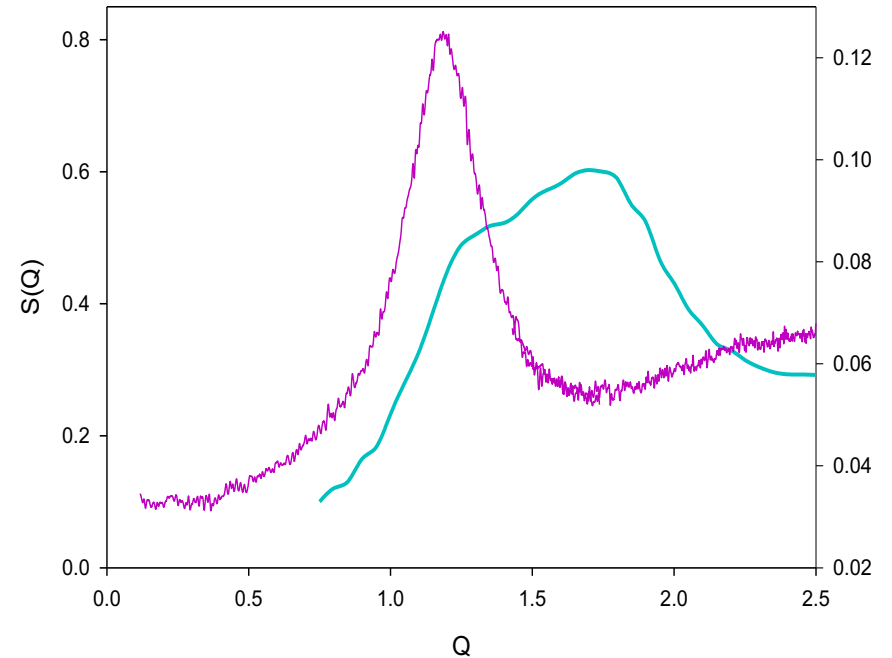
E_∞ changes from one molecule to another : combined effect of the strength of the intermolecular interactions, the shape of the molecule, and the short range correlations.



Effect of the shape : unsaturated and saturated analogues



Toluene ($T_m = 178\text{K}$) and methylcyclohexane ($T_m = 146.8\text{K}$)
 $P_{\text{atm}}, \sim 295\text{K}$



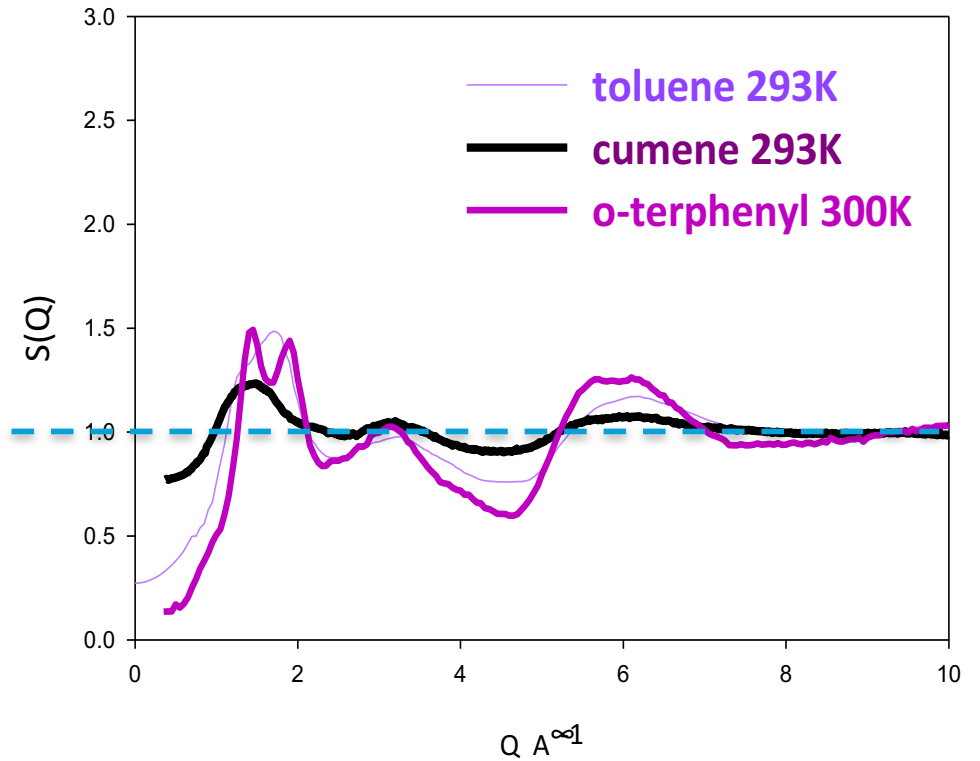
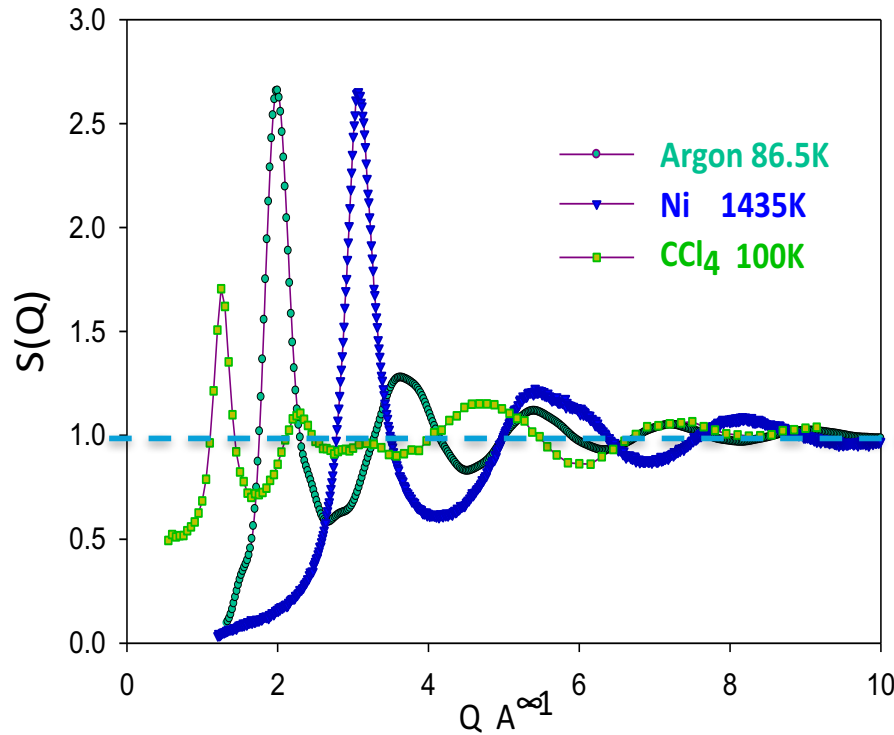
Molecular shape matters, not just size



Consequence of the molecular shape, ability to form a glass Introducing the comparison to metallic glasses

Atomic or spherical molecular liquids do not experience glass formation

Molecular Fragile Glassformers
Liquid, glass



Ar : Yarnell 1973
Ni Holland-Moritz 2008
CCl₄ Pusztai 2009

Very Simple Liquids
(spherical shape)
Only Crystal formation
+ water, benzene

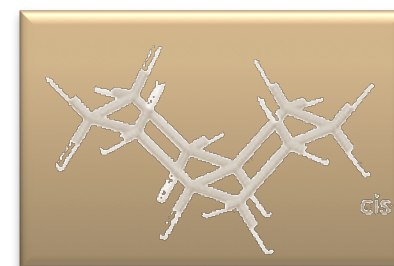
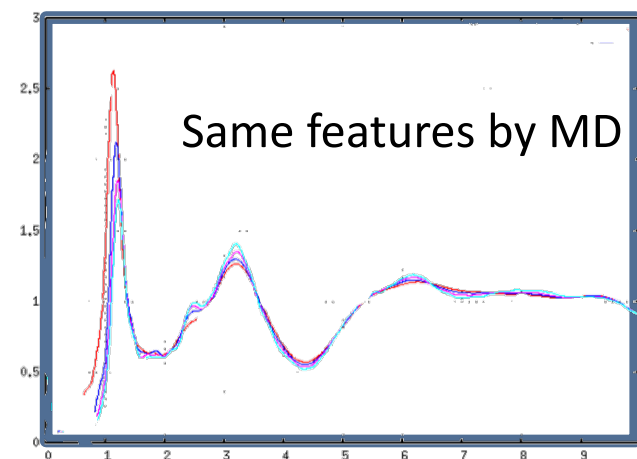


Analysis of the non-trivial behavior of the structure

Large decrease in intensity is related to the compaction of the system, better than simple spheres; whereas the S center-of-mass slightly increases as expected.

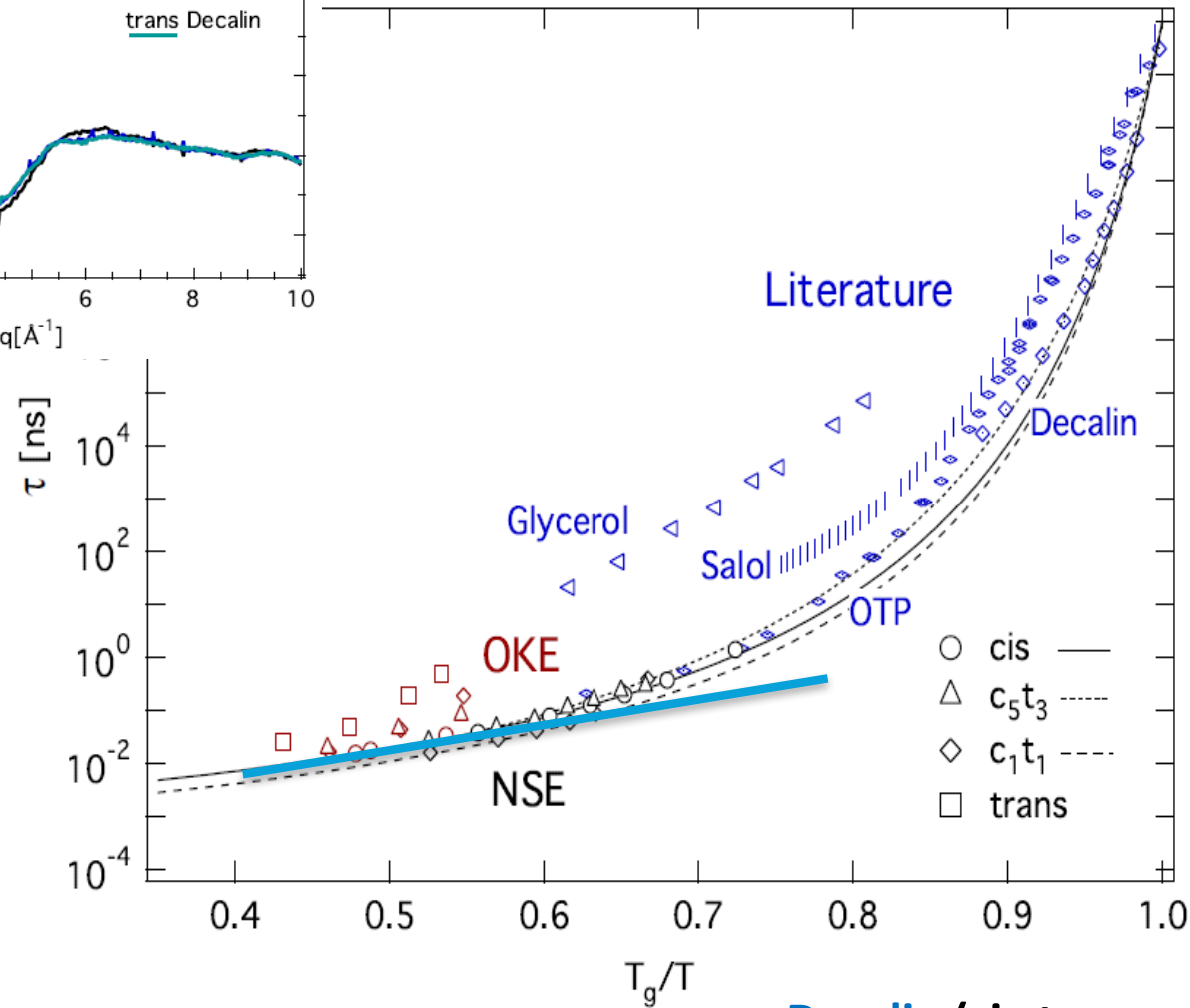
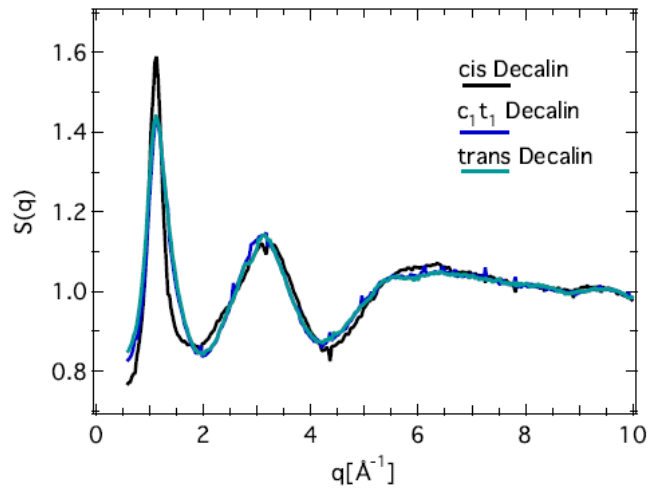


K_T



Changes at $\sim 2.4-2.5 \text{ \AA}^{-1}$ and decrease in intensity of main peak are due to strong changes observed in $S(Q)$ C-H partial, while growing of peak at $\sim 3.2 \text{ \AA}^{-1}$ is mainly due to the H-H partial. C-C partial does not seem to change much as the center of mass.

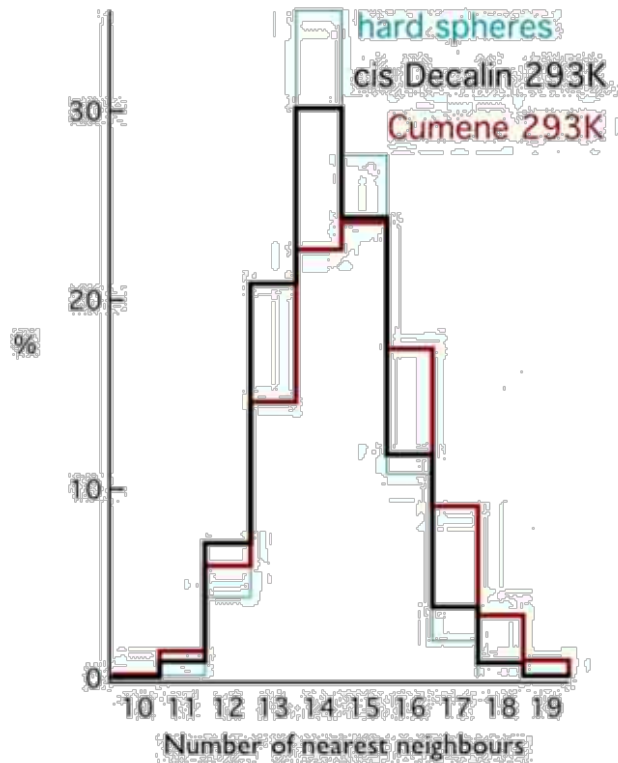
Moreover this liquid might be
 the most simple et fragile one $m_p \sim 149$



Decalin (cis-trans mixture), $m_p = 147$
From Duvvuri- Richert, JCP 2002

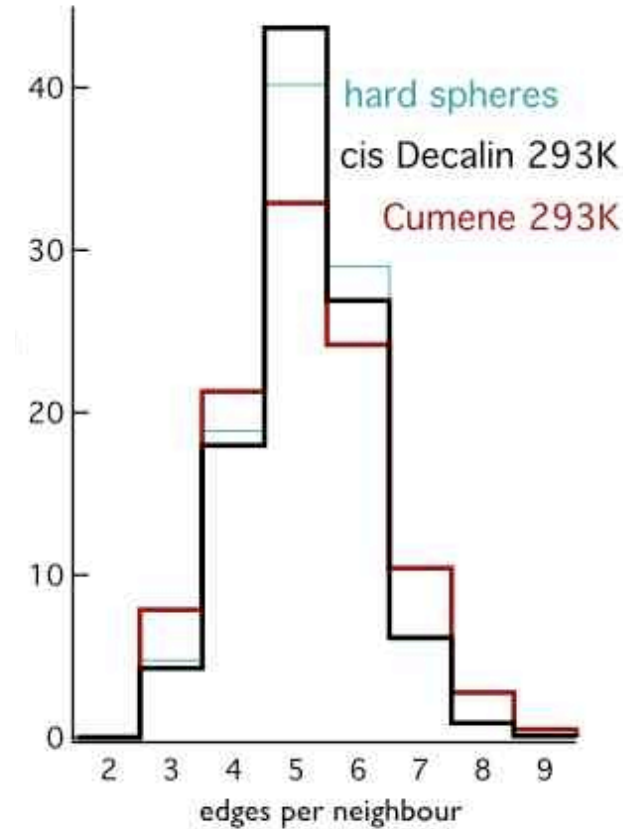
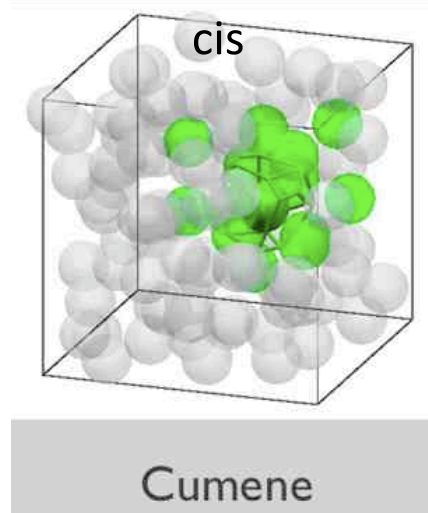
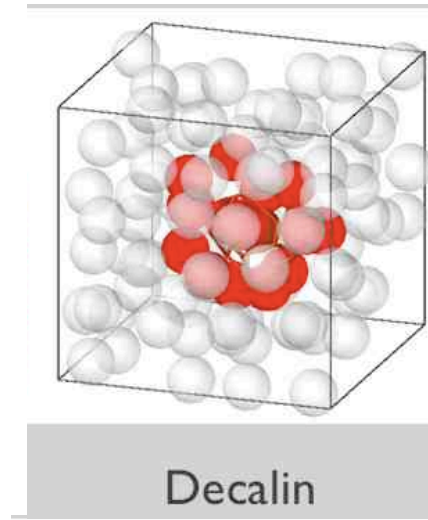
Comparison with **Hard spheres** and **a good molecular glassformer**.

The size and the shape of the local arrangement is different



Neighbour geometry:

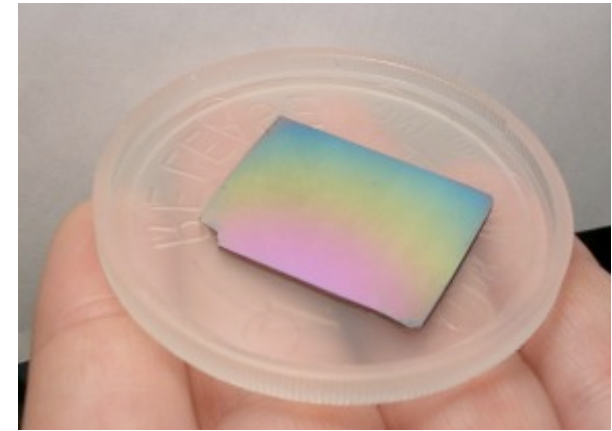
The neighbour geometry for Decalin is relatively close to the geometry of the model of randomly close packed hard spheres.



The ability to form a glass comes from the strong correlations between first neighbors and Interpenetrating aggregates in the high viscous regime ($T < T^*$).

Formation d'un verre stable

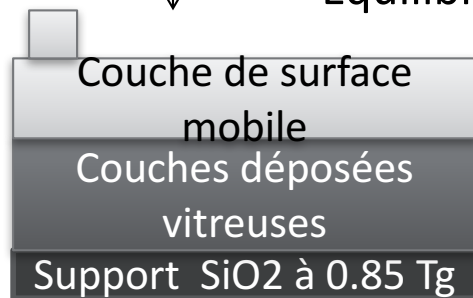
T du substrat = proche de 0.85 Tg
avec une vitesse de dépôt de 0.2 nm/s
jusqu'à une épaisseur de 10-15nm jquà 600nm



Creuset contenant quelques mg
de cristal



Équilibration en 10s



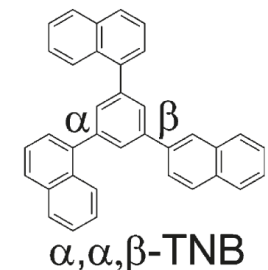
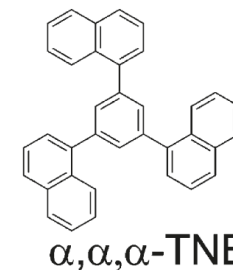
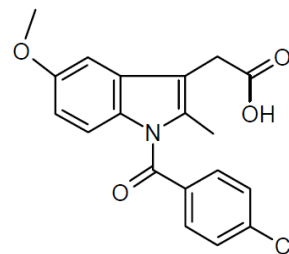
Enceinte à 10⁻⁸ torr

Empilement très efficace
Haute densité

Très différent des techniques de trempe
par dépôt de vapeur avec un substrat
à T << Tg

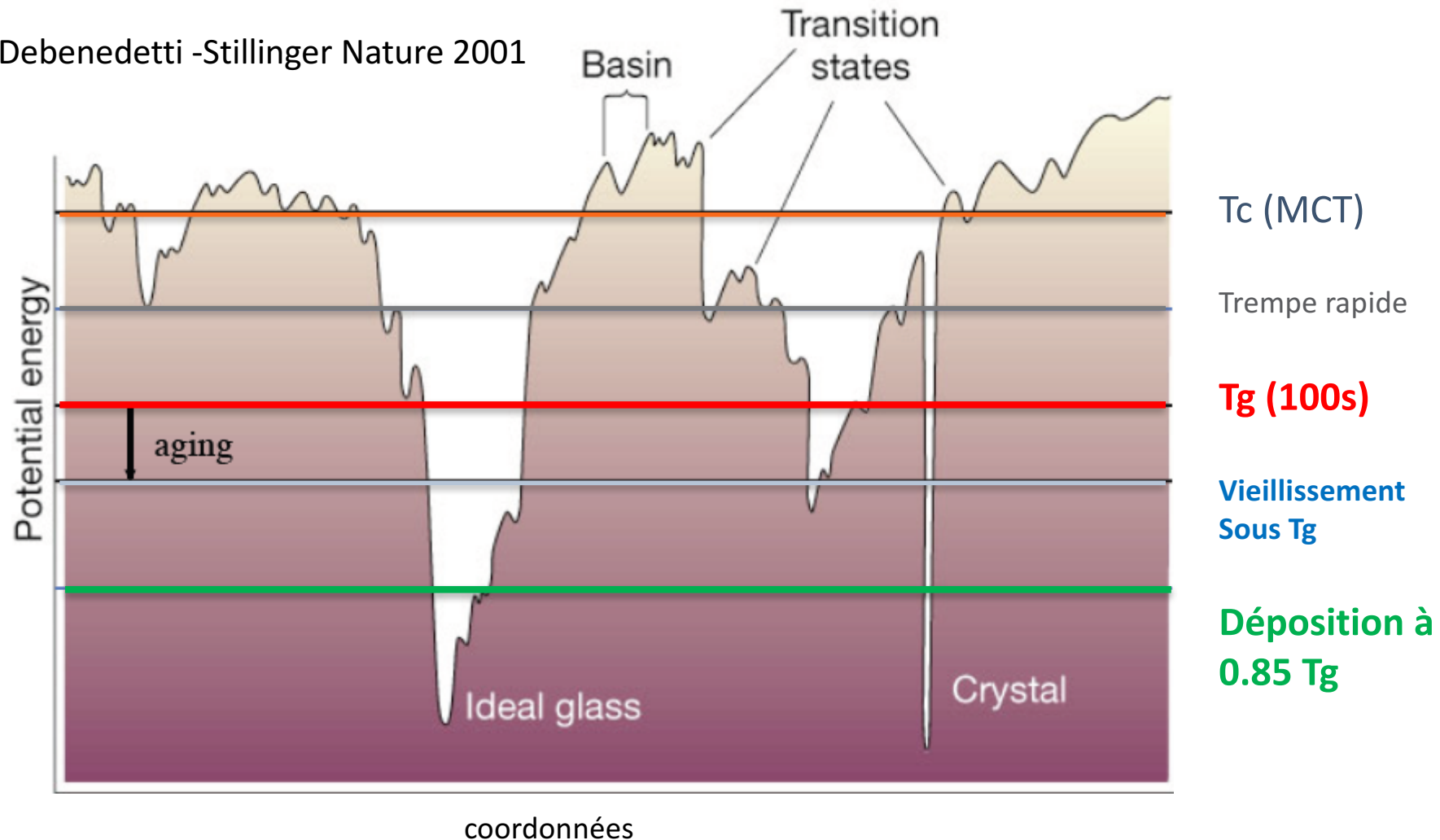
indomethacin (IMC, Tg~315 K)

trisnaphthylbenzene (TNB, Tg~348 K)



Schematic representation Potential energy landscape

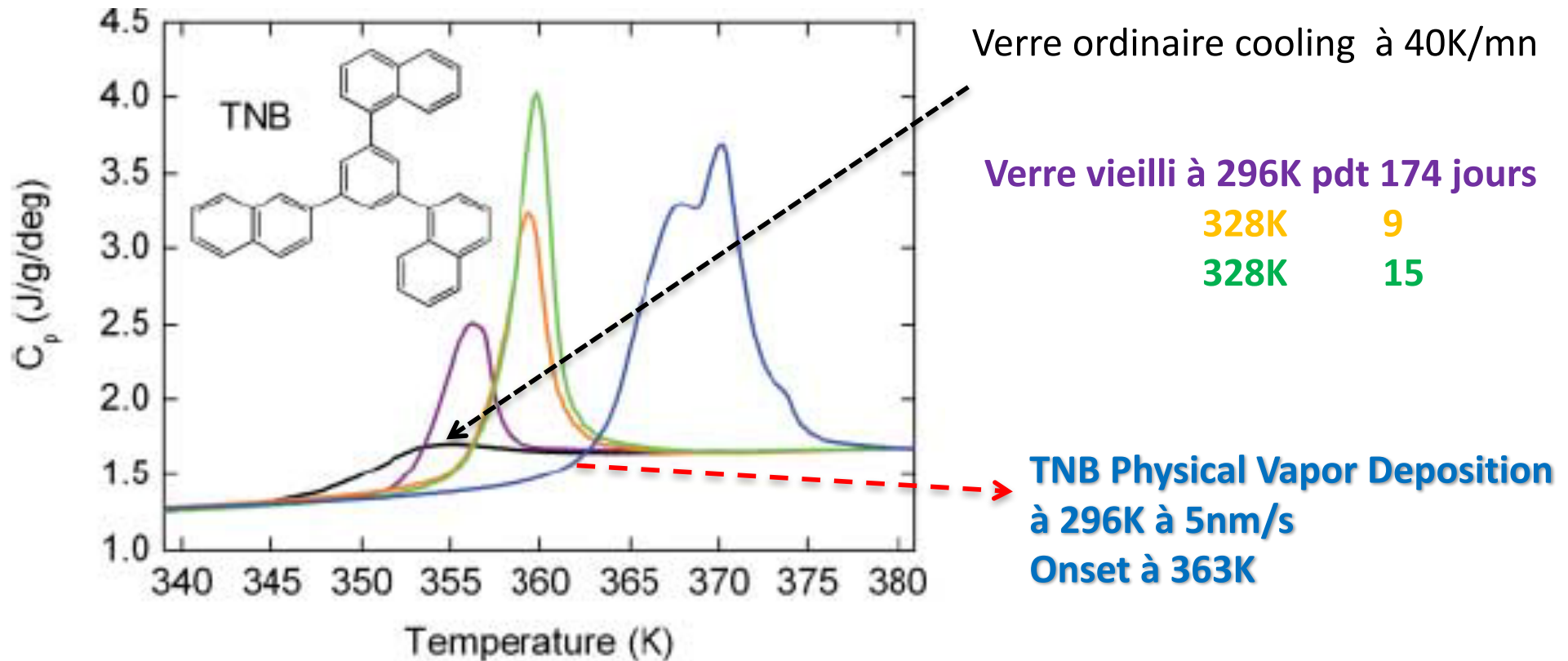
Debenedetti -Stillinger Nature 2001



**Offre la possibilité de découvrir les parties les plus basses de ce paysage et
approcher un possible verre idéal**

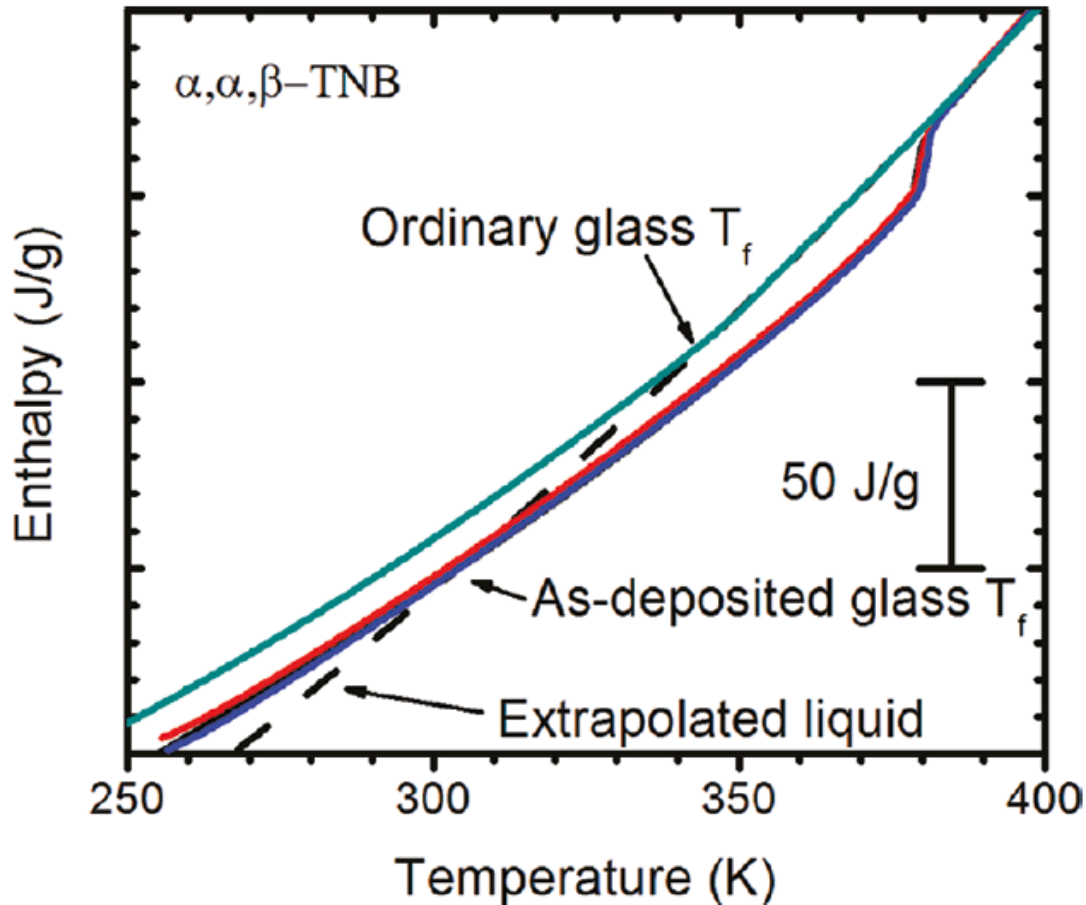
Sans compter les applications technologiques

Signature calorimétrique Cp par DSC TNB



363K est la T nécessaire pour activer la mobilité des molécules , indiquant que le verre PVD est beaucoup plus stable cinétiquement

Conséquence : calcul de la température fictive à partir de l'enthalpie



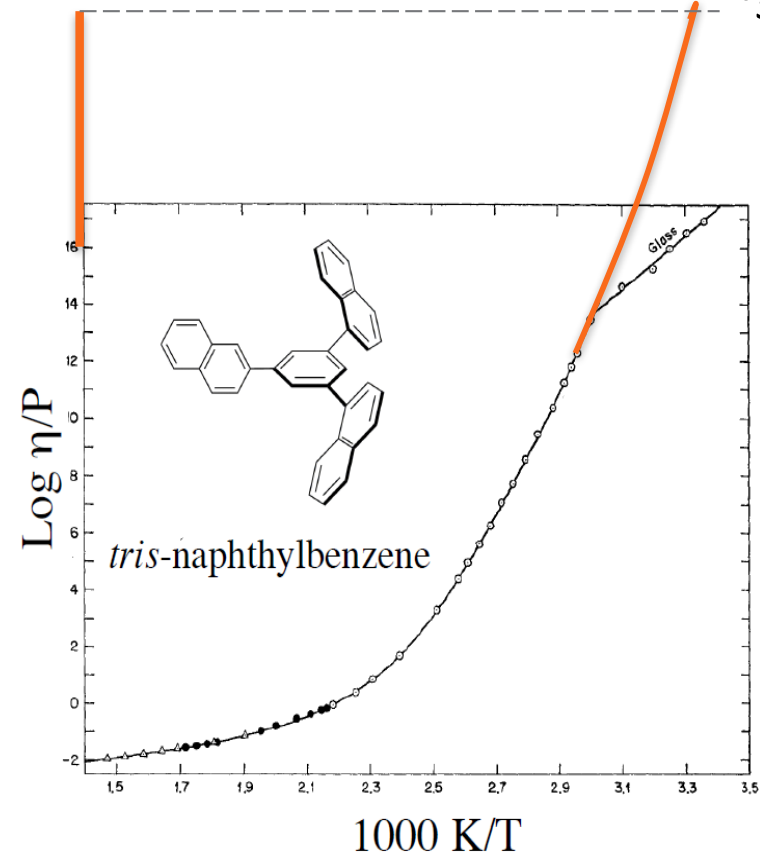
Dawson K, et al J. Phys. Chem. Lett. 2011, 2, 2683–2687

T_f au point de croisement du liquide extrapolé et du PVD verre

$T_f=304K$

En terme de viscosité η ou τ
par extrapolation

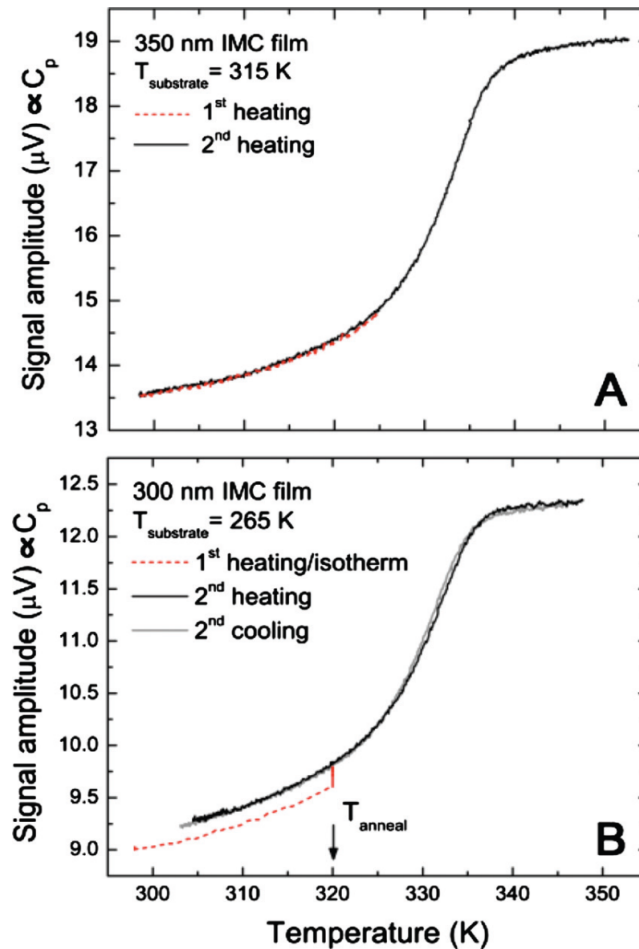
+ 8 décades 3200ans!! ~304K



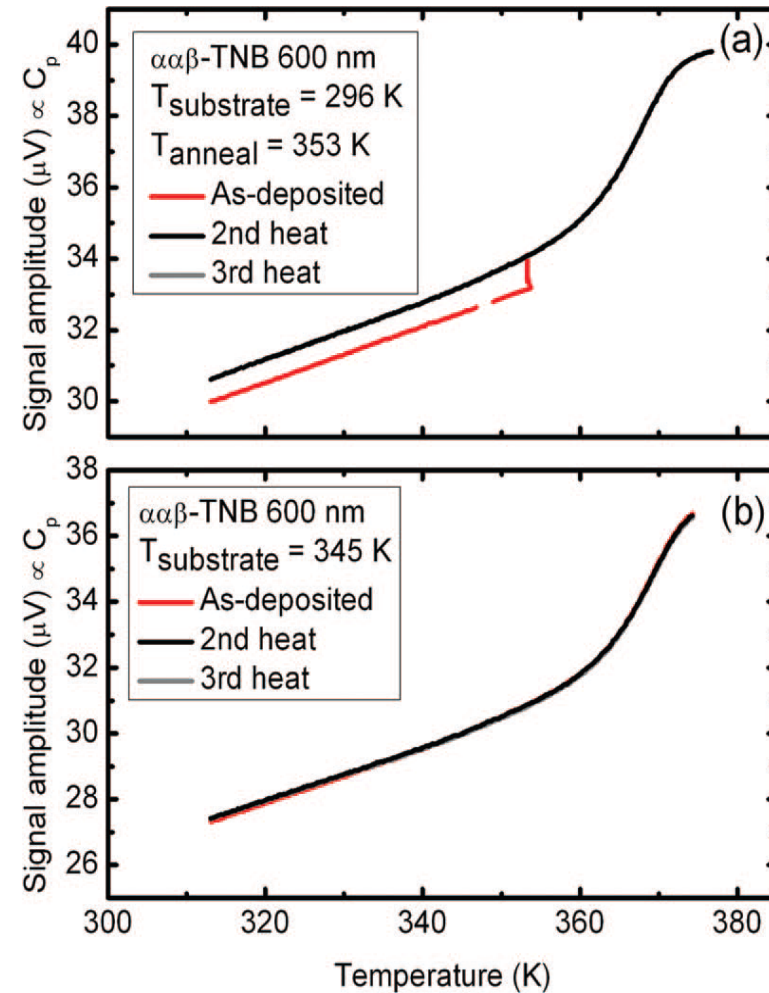


Autres propriétés du VERRE

Une différence de 4-5% sur la valeur de C_p à 300K dans le verre observé par AC nanocalorimétrie



Kenneth L. Kearns et al, J Chem Phys 2010

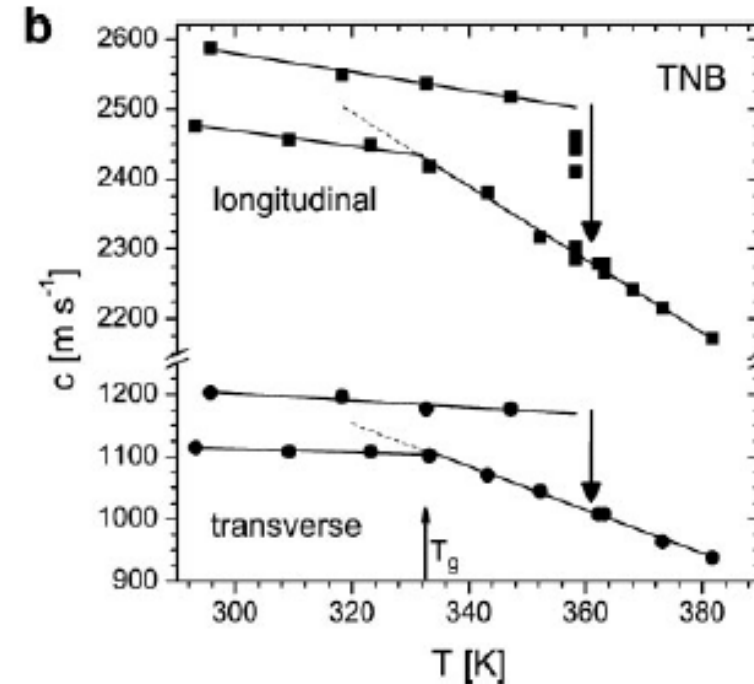
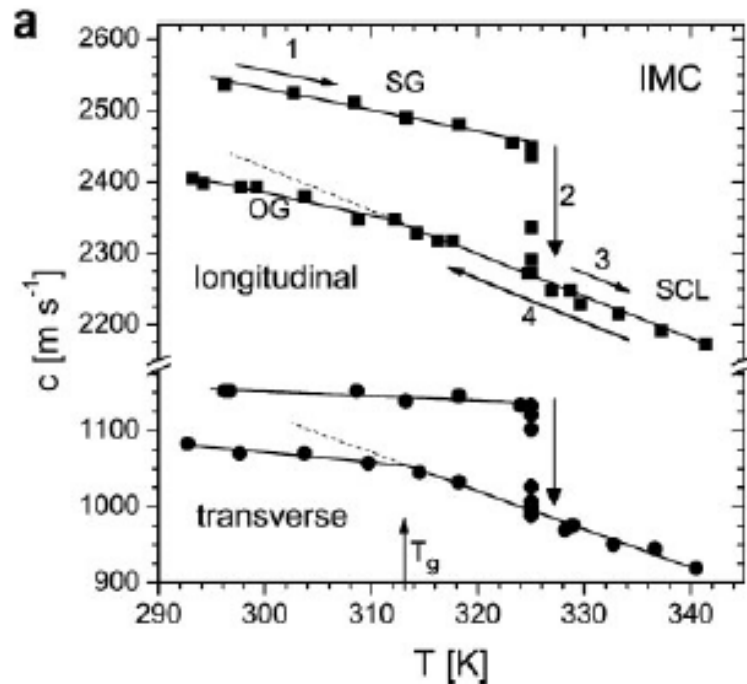


K. Whitaker et al, J Chem Phys 2012

Avec une densification de l'ordre de 1.3% est mesurée par ellipsométrie

High-Modulus Organic Glasses Prepared by Physical Vapor Deposition

Kenneth L. Kearns, et al Adv. Mater. 2009, 21, 1



Longitudinal bulk module M
Module de cisaillement G
Module de Young E

augmentation de 10% à 19%

Alors que par aging, ils ne varient que de .5-3%

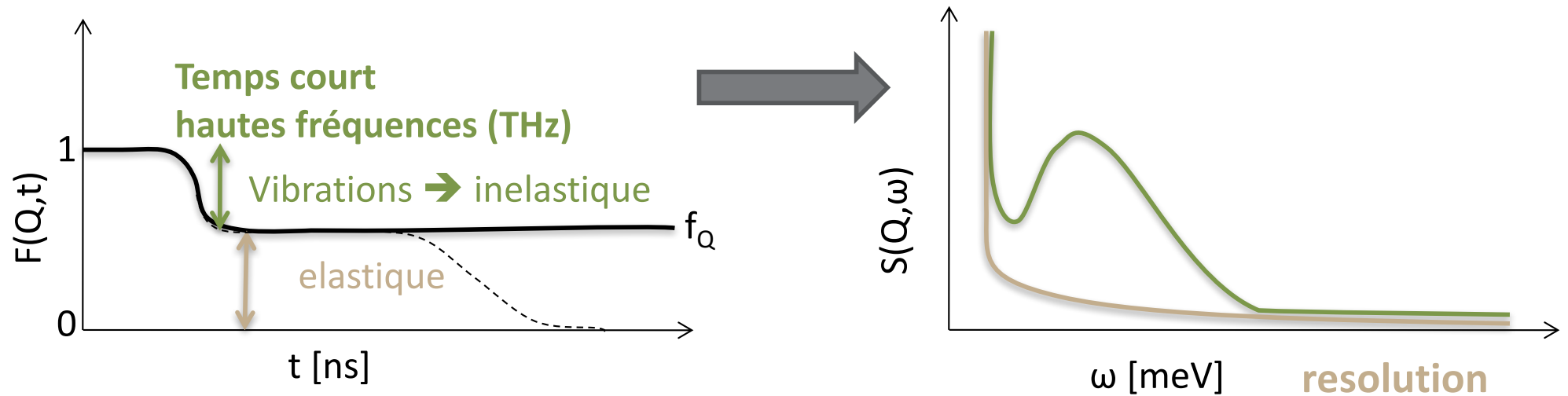


Fréquence de Debye des différents états

$$\omega_D = \left(18\rho\pi^2 / \left(1/c_L^3 + 2/c_T^3 \right) \right)^{1/3}$$

Autres propriétés du VERRE

Verre = photo instantanée du liquide gelé à T_g



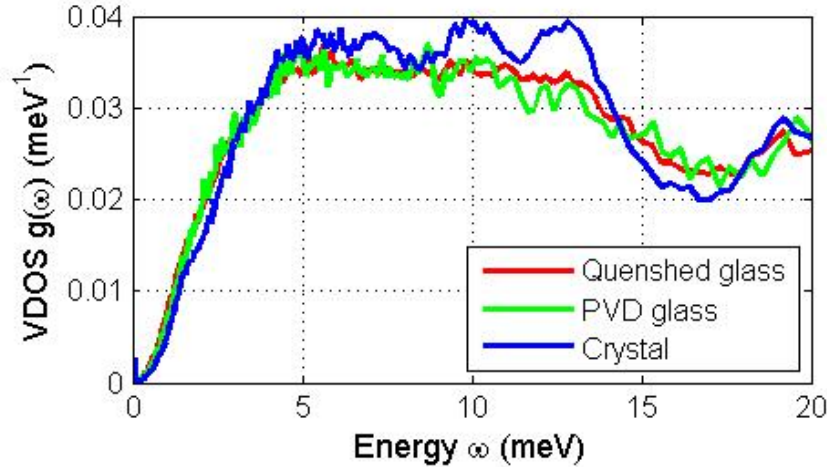
À partir des modes vibrationnelles dans le verre

→ description de la réponse du verre à de petites perturbations

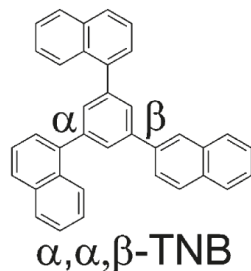
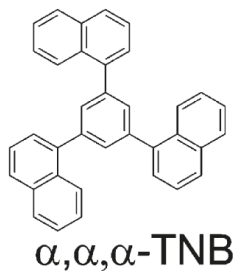
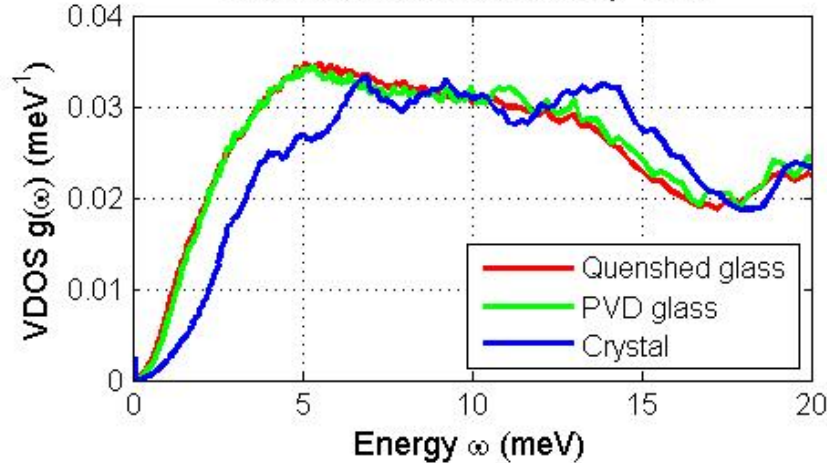
Pour les systèmes désordonnés, ces modes établissent une relation entre le désordre structural et la dynamique collective

Calcul de la VDOS à partir de $S(Q, \omega)$

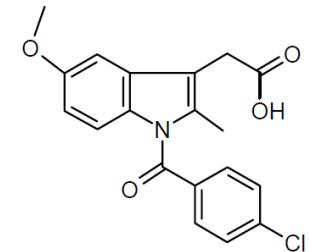
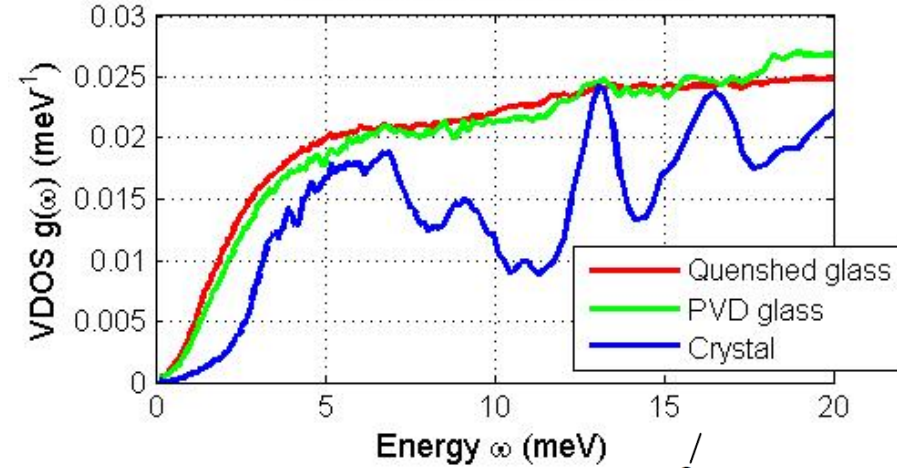
Normalized VDOS of $\alpha\alpha\alpha$ -TNB



Normalized VDOS of $\alpha\alpha\beta$ -TNB



Normalized VDOS of IMC



Des différences moins marquées
à basses fréquences, mais bien présentes

Systèmes moléculaires complexes avec
des modes intramoléculaires à basse freq.



Analyse de la VDOS

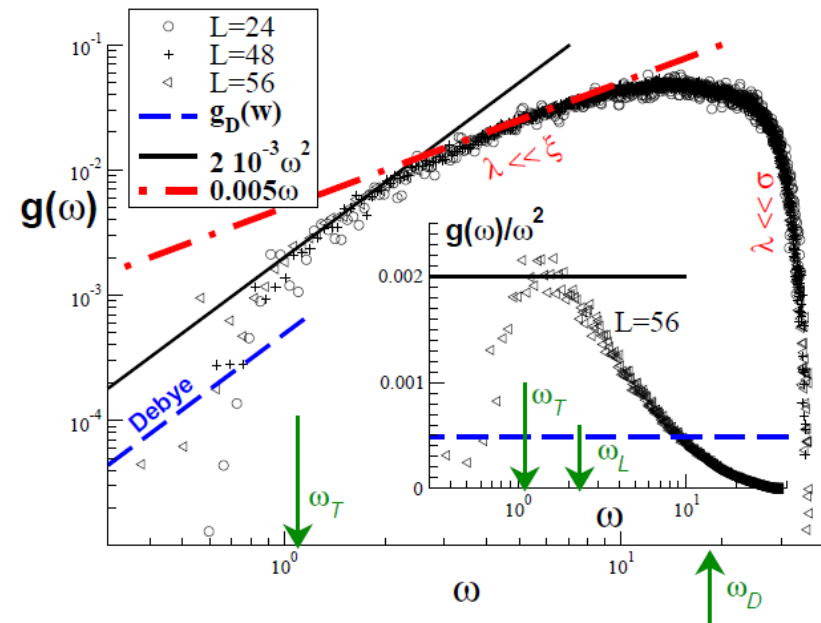
à partir des déplacements des particules d'un système soumis à une contrainte
 → Informations sur les hétérogénéités de l'amorphe,
 Toute déviation à une contrainte uniforme définit des déformations non affines

cross over entre un **comportement élastique continu (Debye)**
 et un **régime de déplacement non affine**

Léonforte et al PRE 2004, 2005

ω^2

ω^1

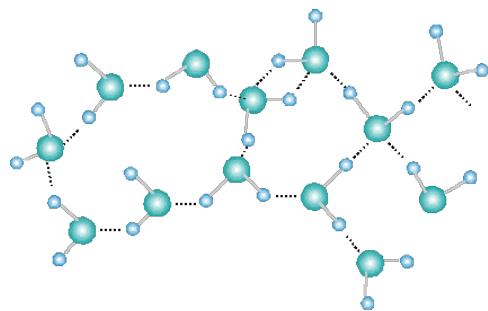
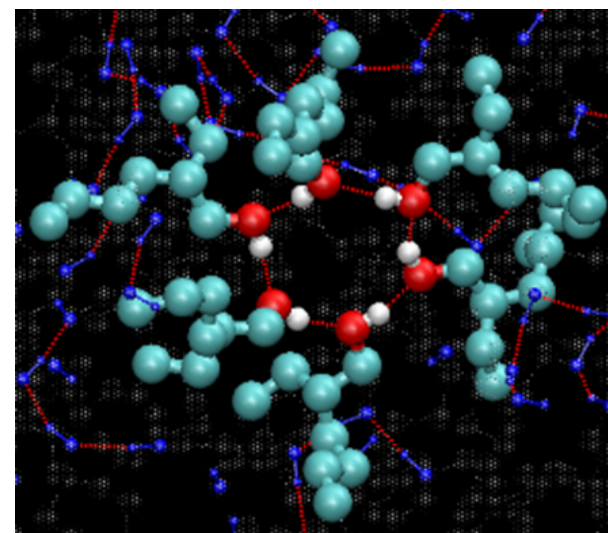
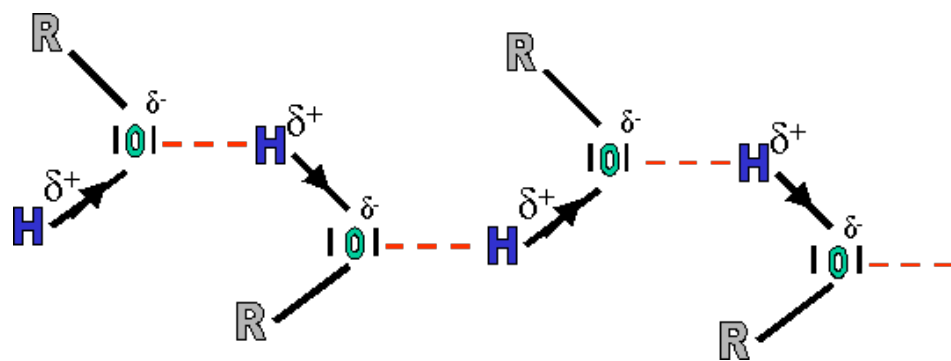


'hétérogénéités élastiques' de plus petites tailles dans les verres stables

Networks of molecular liquids ?

Role of Hydrogen Bonds

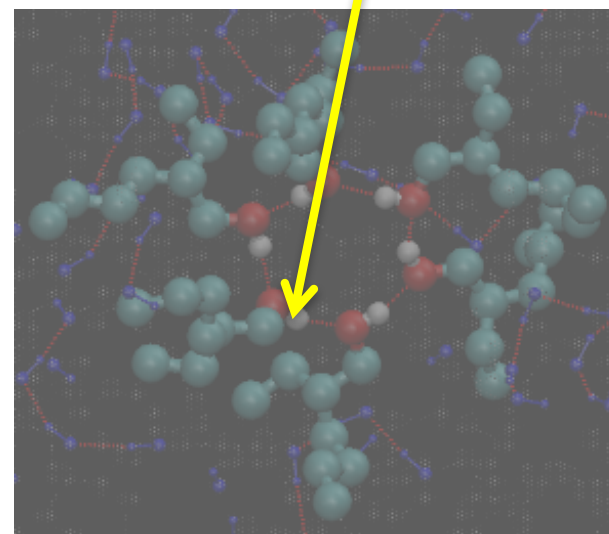
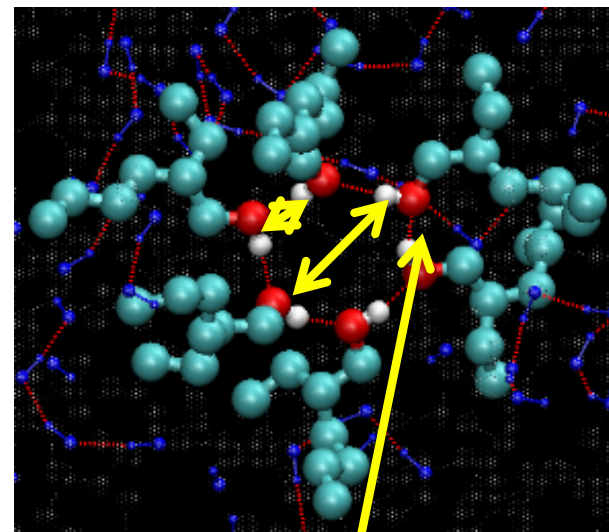
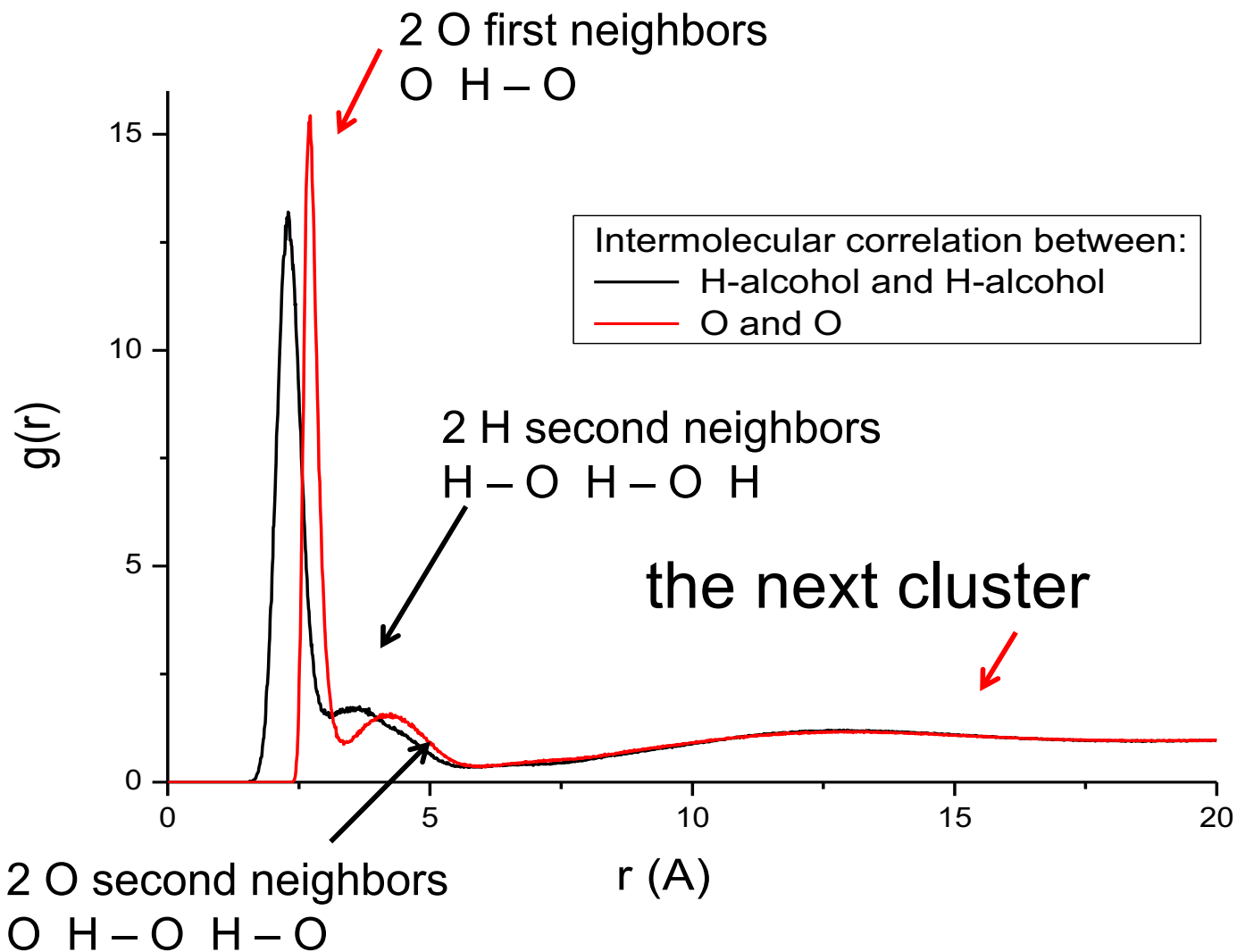
Auto-association of alcohols



water : 17 cristalline phases

1ethyl 2hexanol MD analysis 300K: Partial radial distribution function involving O -H atoms

**Evidence of various length scales,
showing cluster formation**





1ethyl 2hexanol static structure factor at Patm

Same remarkable feature: a well defined prepeak



Fully deuterated sample
Elastic neutron scattering

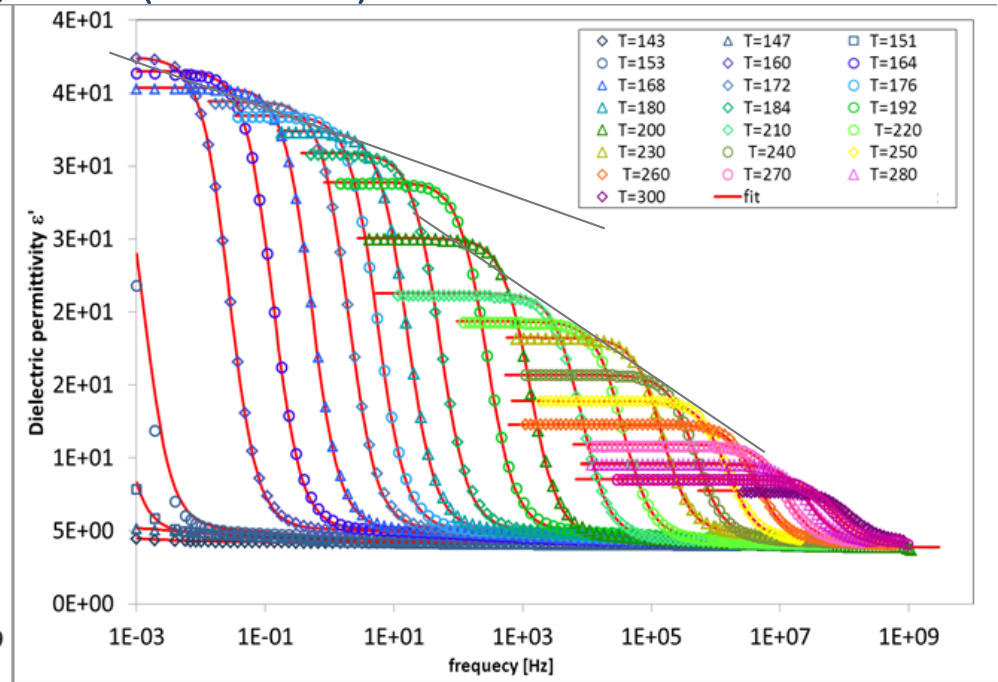
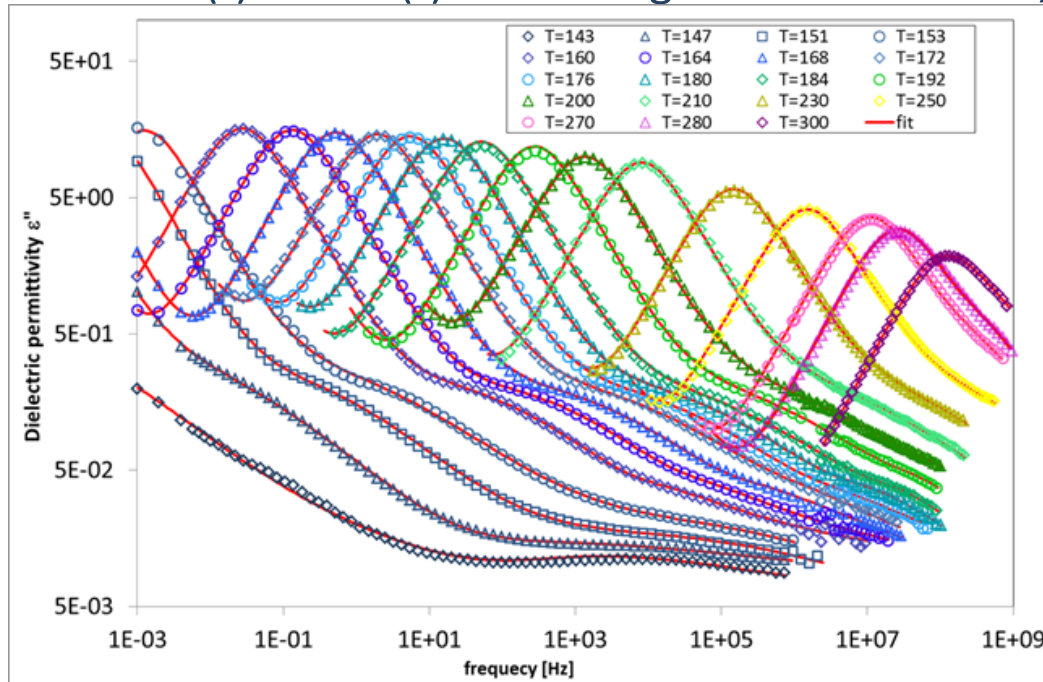
Fully hydrogenated
Xray scattering

Case of 1ethyl 2hexanol dielectric spectra

Prior work by Richert and Böhmer over the last years

The largest ratio τ_D / τ_α

$\epsilon''(f)$ and $\epsilon'(f)$ over a large T and ω domain, Patm (this work)



↓
mHz

↓
GHz

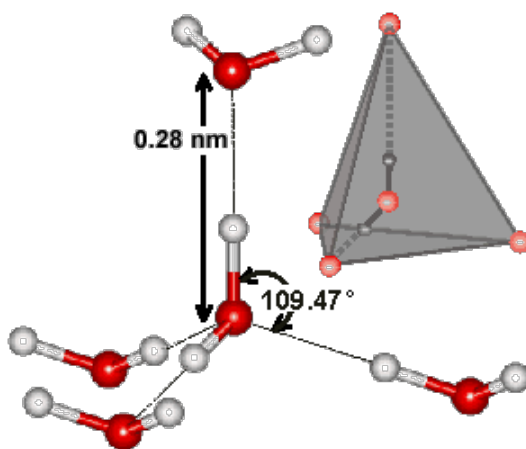
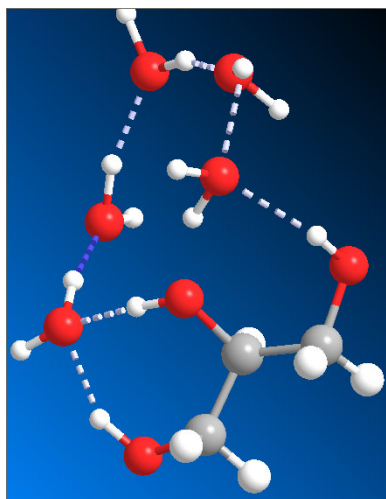
$$g_K = \frac{1}{\mu^2} \frac{9\epsilon_0 k_B T}{N} \frac{(\epsilon_s - \epsilon_\infty)(2\epsilon_s + \epsilon_\infty)}{\epsilon_s (\epsilon_\infty + 2)^2}$$

45

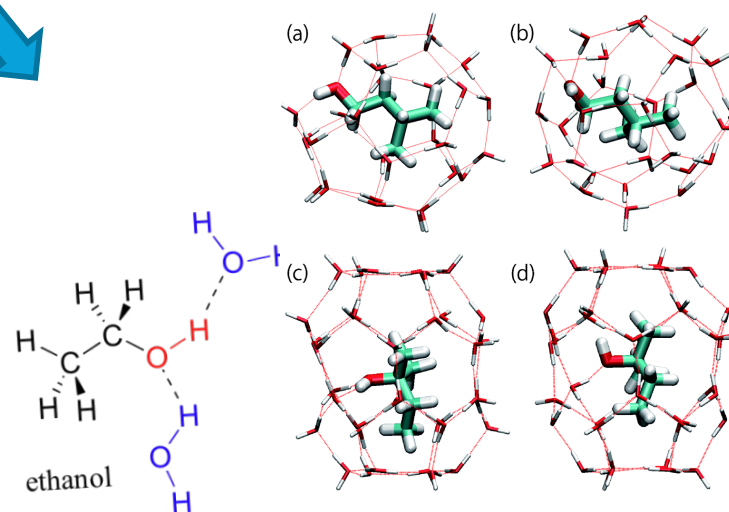
Role of Hydrogen Bonds in solutions

Water + alcohol, aminoacids...

lower hydrophobicity
better to delay nucleation
and retard growth,
working by
disrupting the water
hydrogen bond network,



higher hydrophobicity
strengthens local water
structure
control of hydrate formation



Clathrates Yasuoka, 2015

Water + glycerol phase diagram case 1 : strongly hydrophilic alcohol

Inaba (2007)

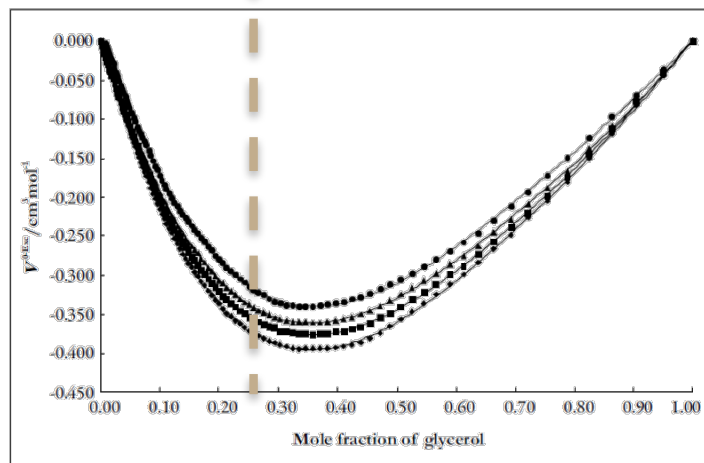
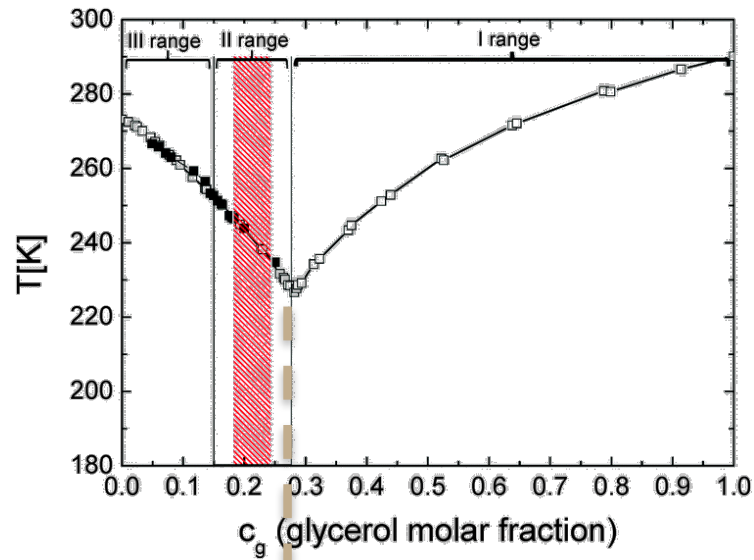
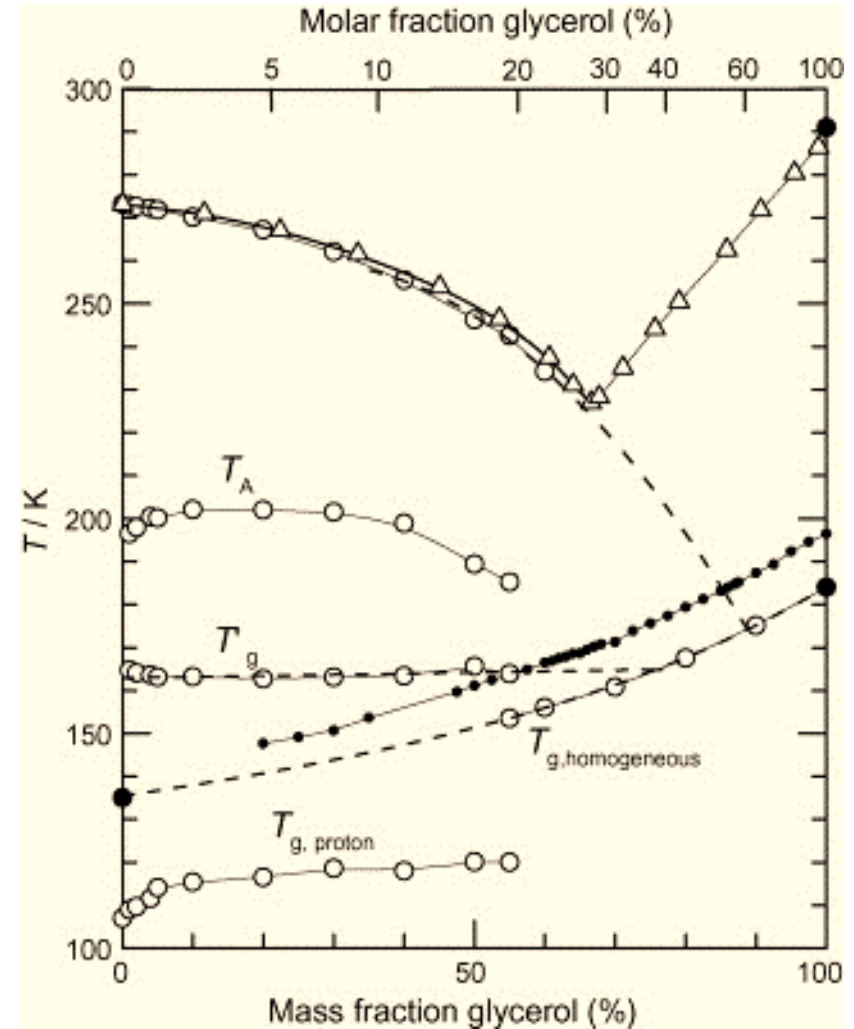
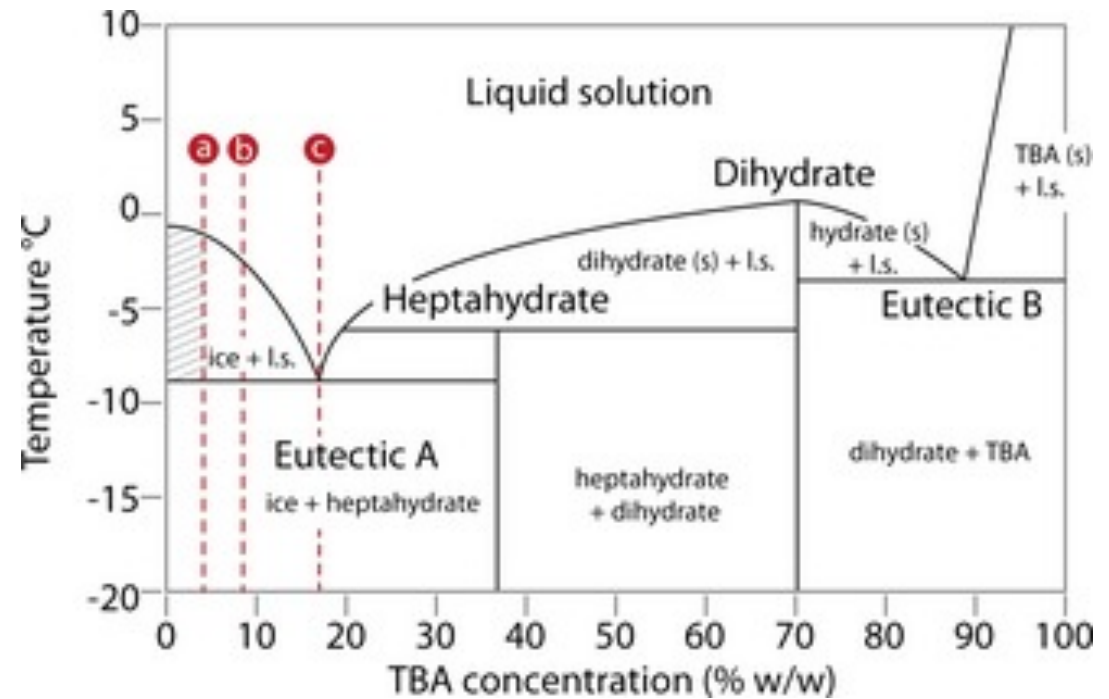
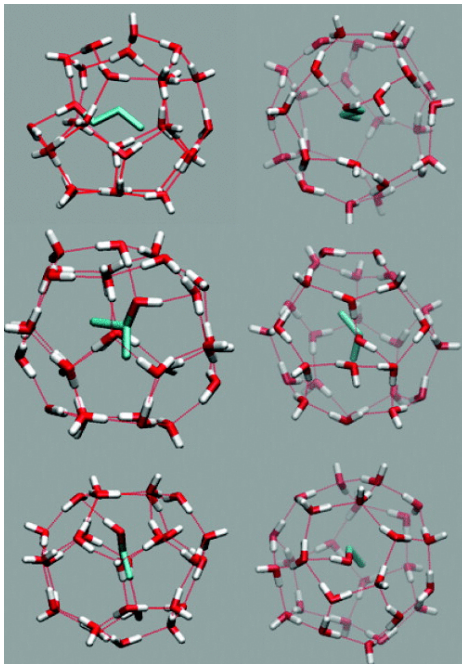
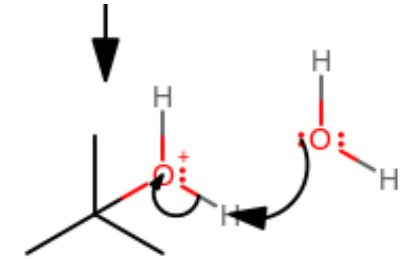


Figure 1. Excess molar volumes of glycerol + water mixtures at several temperatures. (◆): 288.15 K; (■): 293.15 K; (▲): 298.15 K; (●): 303.15 K.



Tanaka (2012) mais dans Nat mat

Water + monoalcohol phase diagram case 1 : partly hydrophobic alcohol



Clathrates

GT is a combination of cooperative effects and local, molecular ones.

The local structure is distinct from that of the crystal

the molecular details play a fundamental role in its dynamics (de Gennes, Narrowing, rigid vs flexible)

The high T behavior, the ability to form a glass etc.... results from the combination of

- the strength of the intermolecular interactions, attraction, dispersion
- the shape of the molecule,
- the short-range (spatial) correlations associated with the local arrangement of the molecules in the liquid.

The simplest system that forms a glass is a **molecule**, with a reduced number of partial contributions to $S(Q)$, a simplified the shape keeping intrinsic and self generated disorder.

It presents a low activation energy at high Temperature and very high fragility.

No increase in static correlation length while dynamical ones varie by a factor of 3 to 4.

