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#### **Refining techniques and developments**

Franck Pigeonneau Mines Paris | PSL Univ. - Centre of Materials Forming, CNRS UMR 7635, Sophia Antipolis, France



2. Mass transfer around a bubble

#### 3. Fining process

- 3.1 Sulphate fining
- 3.2 Water fining
- 3.3 Helium fining
- 3.4 Centrifugal and low pressure fining

#### 4. Synthesis





Glass bath is an open reactor system.

Existence of a residence time distribution E(t).

Arbitrary reactor



 $E(t) \sim$  probability density function:

$$\int_{0}^{\infty} E(t)dt = 1, \qquad (1)$$
$$\langle t \rangle = \int_{0}^{\infty} tE(t)dt. \qquad (2)$$





- Glass bath is an open reactor system.
- Existence of a residence time distribution E(t).





 $\delta(t-t_s). \qquad (3)$ 

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$$\int_{0}^{\infty} E(t)dt = 1, \qquad (1) \qquad E(t) =$$
$$\langle t \rangle = \int_{0}^{\infty} t E(t)dt. \qquad (2)$$





- Glass bath is an open reactor system.
- Existence of a residence time distribution E(t).





#### Continuous Stirred Tank Reactor



 $E(t) \sim$  probability density function:

$$\int_{0}^{\infty} E(t)dt = 1, \qquad (1)$$
$$\langle t \rangle = \int_{0}^{\infty} tE(t)dt. \qquad (2)$$

$$E(t) = \delta(t - t_s).$$
 (3)  $E(t) = \frac{e^{-t/t_s}}{t_s}.$  (4)





- Glass bath is an open reactor system.
- Existence of a residence time distribution E(t).



 $t_s$  is useful to normalise t and E(t)

$$\bar{t} = \frac{t}{t_s}, \ \bar{E}(\bar{t}) = t_s E(t).$$
 (6)





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Figure 1: Residence time distribution of PFR, CSTR and industrial glass furnace.  $\ensuremath{\texttt{PSL}\xspace}$  | Cemef





Figure 2: Bubble population in borosilicate liquid from post-mortem analysis<sup>1</sup>.

<sup>&</sup>lt;sup>1</sup>L. Pereira et al.: Experimental study of bubble formation in a glass-forming liquid doped with cerium oxide, in: J. Am. Ceram. Soc. 103 (2020), pp. 2453–2462.



Figure 3:  $\mu_0 = N(t)$  vs. t for T=1320, 1370 and 1420 °C<sup>2</sup>. Solid circles from Bastick<sup>3</sup>.  $N(t) = N_0 e^{-\alpha t}, N_0 \approx 10^8 \text{ m}^{-3}.$  (7)

<sup>2</sup>F. Pigeonneau/L. Pereira/A. Laplace: Dynamics of rising bubble population undergoing mass transfer and coalescence in highly viscous liquid, in: Chem. Eng. J. 455.2 (2023), p. 140920.
<sup>3</sup>R. E. Bastick: Laboratory experiments on the refining of glass, in: Symposium sur l'affinage du verre, Paris 1956, pp. 127–138.

#### In close reactor, degree of conversion:

$$X_{CR}(t) = 1 - \frac{N(t)}{N_0} = 1 - e^{-\alpha t}.$$
 (8)

7

▶ In open reactor, the degree of conversion is given by<sup>4</sup>:

$$X_{OR} = \int_0^\infty X_{CR}(t) E(t) dt = 1 - G(\alpha),$$
(9)  
$$G(\alpha) = \int_0^\infty E(t) e^{-\alpha t} dt, \text{ Laplace transformation.}$$
(10)

<sup>&</sup>lt;sup>4</sup>J. Villermaux: Génie de la réaction chimique, Paris 1993.

Flat glass: less than 1 bubble of  $200 \,\mu\text{m}/20 \,\text{m}^2 \implies 10 \text{ bubbles/m}^3$ :

$$G(\alpha) = 1 - X_{OR}(\alpha) = 10^{-7}.$$
 (11)

▶ Container glass: less than 1 bubble/bottle  $\rightarrow$  10<sup>4</sup> bubbles/m<sup>3</sup>:

$$G(\alpha) = 1 - X_{OR}(\alpha) = 10^{-4}.$$
 (12)

• According to Bastick<sup>5</sup>, for T=1350 °C:

$$\alpha \approx 5 \cdot 10^{-4} \text{ s}^{-1}. \tag{13}$$

<sup>&</sup>lt;sup>5</sup>Bastick: Laboratory experiments on the refining of glass (see n. 3).

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Figure 4:  $G(\alpha) = 1 - X(\alpha)$  vs.  $\alpha t_s$  for PFR, CSTR and industrial glass furnace.



According to Pereira et al.<sup>6</sup>:

$$\alpha \sim \frac{\rho g a^2}{\eta(T) H}$$

 $<sup>^{6}\</sup>mbox{Pereira et al.:}$  Experimental study of bubble formation in a glass-forming liquid doped with cerium oxide (see n. 1).



<sup>&</sup>lt;sup>6</sup>G. E. Kunkle/W. M. Welton/R. L. Schwenninger: Melting and vacuum refining of glass or the like and composition of sheet, US Patent 4,738,938, PPG Industries, Inc., 1988.









<sup>&</sup>lt;sup>6</sup>J. Ferguson: Centrifugal glass-melting furnace, US Patent 2,006,947, 1930.



<sup>&</sup>lt;sup>6</sup>P. Jeanvoine et al.: Procédé et dispositif de fusion et d'affinage de matières vitrifiables, Eur. Patent 0 970 021 B1, Saint-Gobain Glass France, 2005.

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Figure 5:  $Eo - Re diagram^7$ .

<sup>&</sup>lt;sup>7</sup>R. Clift/J. R. Grace/M. E. Weber: Bubbles, Drops, and Particles, New York 1978.



$$Eo = Bo = \frac{\rho g D^2}{\gamma} < 40$$
$$Re = \frac{\rho U D}{\eta} < 1$$
$$Mo = \frac{g \eta^4}{\rho \gamma^2} \approx 3 \times 10^2$$

Figure 5:  $Eo - Re diagram^7$ .

<sup>&</sup>lt;sup>7</sup>Clift/Grace/Weber: Bubbles, Drops, and Particles (see n. 7).



Figure 6: Drag law in glass forming liquids.

<sup>&</sup>lt;sup>8</sup>R. B. Jucha et al.: Bubble rise in glassmelts, in: J. Am. Ceram. Soc. 65 (1982), pp. 289–292.
<sup>9</sup>E. J. Hornyak/M. C. Weinberg: Velocity of a freely rising gas bubble in a soda-lime silicate glass melt, in: J. Am. Ceram. Soc. 67 (1984), pp. C244–C246.



(b)



Figure 7: Film drainage at a free surface of high viscous liquids<sup>10</sup>: (a) expe. setup, (b)  $\overline{h}$  vs.  $\overline{t}$ .

<sup>&</sup>lt;sup>10</sup>H. Kočárková/F. Rouyer/F. Pigeonneau: Film drainage of viscous liquid on top of bare bubble: Influence of the Bond number, in: Phys. Fluids 25 (2013), p. 022105.



Figure 8: Stream functions inside and outside of a rising bubble.

According to Hadamard<sup>a</sup> & Rybczynski<sup>b</sup>:

$$\mathbf{F} = -4\pi\eta a (\mathbf{V} - \mathbf{U}),$$
 (14)  
 $\mathbf{V} = -rac{
ho \mathbf{g} a^2}{3\eta}.$  (15)

<sup>a</sup>J. Hadamard: Mouvement permanent lent d'une sphére liquide et visqueuse dans un liquide visqueux, in: C. R. Acad. Sci. Paris 152 (1911), pp. 1735–1738.

<sup>b</sup>W. Rybczynski: Uber die fortschreitende bewegun einer flussingen kugel in einem zaben medium, in: Bull. de l'Acad. des Sci. de Cracovie, série A 1 (1911), pp. 40–46.



$$\frac{dn_i}{dt} = \int_S \mathcal{D}_i \nabla C_i \cdot \mathbf{n} dS = 4\pi a^2 \mathbf{J}_i, \quad (16)$$
$$\left[ P_0 + \rho(H - z) + \frac{2\gamma}{a} \right] \frac{4\pi a^3}{3} = \left( \sum_{i=1}^N n_i \right) \mathcal{R} \mathcal{T}, \quad (17)$$
$$\frac{d\mathbf{x}}{dt} = \mathbf{u} - \frac{\rho \mathbf{g} a^2}{3\eta}. \quad (18)$$





























▶ Without convection and in steady-state regime:

$$\frac{d}{dr}\left(r^{2}\frac{dC_{i}}{dr}\right) = 0,$$

$$C_{i}(r) = C_{i}^{\infty} - \frac{(C_{i}^{\infty} - C_{i}^{s})a}{r},$$

$$J_{i} = \frac{\mathcal{D}_{i}(C_{i}^{\infty} - C_{i}^{s})}{a}.$$
(19)
(20)
(21)





In the case of a rising bubble, the problem is normalised:

$$\bar{\mathbf{x}} = \frac{\mathbf{x}}{2a}, \ \bar{C}_i = \frac{C_i - C_i^{\infty}}{C_i^s - C_i^{\infty}}.$$
(22)

In steady-state and in spherical coordinate system:

$$\begin{bmatrix} \bar{u}_{r} \frac{\partial \bar{C}_{i}}{\partial \bar{r}} + \frac{\bar{u}_{\theta}}{\bar{r}} \frac{\partial \bar{C}_{i}}{\partial \theta} \end{bmatrix} = \frac{1}{\operatorname{Pe}_{i} \bar{r}^{2}} \begin{bmatrix} \frac{\partial}{\partial \bar{r}} \left( \bar{r}^{2} \frac{\partial \bar{C}_{i}}{\partial \bar{r}} \right) + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \bar{C}_{i}}{\partial \theta} \right) \end{bmatrix}, \quad (23)$$
$$\operatorname{Pe}_{i} = \frac{V_{T} 2a}{\mathcal{D}_{i}}. \quad (24)$$







Figure 10: Re & Pe vs. *a* (m).







(b) Stretched coordinate<sup>a</sup>:  $r = [1 + \delta(\text{Pe}_i)\zeta]/2$ 



<sup>a</sup>E. J. Hinch: Perturbation Methods, 1991.

Figure 11: Chemical boundary layer around a rising bubble.

$$\left(\frac{\bar{u}_r}{\delta}\frac{\partial\bar{C}_i}{\partial\zeta} + \frac{\bar{u}_\theta}{1+\delta\zeta}\frac{\partial\bar{C}_i}{\partial\theta}\right) = \frac{1}{2\operatorname{Pe}_i\delta^2(1+\delta\zeta)^2}\left\{\frac{\partial}{\partial\zeta}\left[(1+\delta\zeta)^2\frac{\partial\bar{C}_i}{\partial\zeta}\right] + \frac{4\delta^2}{\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial\bar{C}_i}{\partial\theta}\right)\right\} (25)$$

$$\frac{\bar{u}_r}{\delta} = -\frac{\cos\theta}{2} \left[ \frac{2\zeta}{1+\hat{\eta}} + \frac{3\hat{\eta}-2}{2(1+\hat{\eta})} \delta\zeta^2 + \mathcal{O}(\delta^2) \right],\tag{26}$$

$$\frac{\bar{u}_{\theta}}{1+\delta\zeta} = -\frac{\sin\theta}{4} \left[ -\frac{2}{1+\hat{\eta}} - \frac{4+10\hat{\eta}}{1+\hat{\eta}}\delta\zeta + \frac{6+19\hat{\eta}}{1+\hat{\eta}}\delta^2\zeta^2 + \mathcal{O}(\delta^3) \right].$$
(27)

From the *principle of least degeneracy*<sup>11</sup>: For solid particle or immobile interface  $\Rightarrow$ For bubble  $\Rightarrow \hat{\eta} \rightarrow 0$ :  $\hat{\eta} \to \infty$ :  $\delta \propto 1/\sqrt{\mathrm{Pe}_i}$ . (29)

$$\delta \propto 1/\sqrt[3]{\mathrm{Pe}_i}.$$
 (28)

<sup>&</sup>lt;sup>11</sup>M. Van Dyke: Perturbation methods in fluid mechanics, Stanford, California 1975.

The molar flux becomes:

$$\frac{dn_i}{dt} = 4\pi a^2 J_i, \tag{30}$$

$$J_{i} = \frac{(C_{i}^{s} - C_{i}^{\infty})\mathcal{D}_{i}}{2a} \operatorname{Sh}_{i}, \text{ without motion } \operatorname{Sh}_{i} = 2.$$
(31)  
$$\operatorname{Sh}_{i} = \frac{2a}{\pi\delta} \int_{S} \frac{\partial \bar{C}}{\partial n} dS, \text{ Sherwood number.}$$
(32)

Levich's solution<sup>12</sup>:

$$Sh_i = 0.991\sqrt[3]{Pe_i}.$$
 (33)  $Sh_i = 0.651\sqrt{Pe_i}.$  (34)

<sup>&</sup>lt;sup>12</sup>V. G. Levich: Physicochemical hydrodynamics, Englewood Cliffs, N.J. 1962.



Figure 12: Sh vs. Pe for solid particle and bubble<sup>13</sup>.

<sup>&</sup>lt;sup>13</sup>Clift/Grace/Weber: Bubbles, Drops, and Particles (see n. 7).

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Figure 13: *a* vs. *t* of O<sub>2</sub> bubble at T=1400 °C,  $C_{\rm Fe}=2.8 \times 10^{-2}$  wt%,  $\mathcal{R}_{\rm Fe}=0.575$ .



Solve the advection/diffusion/reaction equation:

$$\frac{DC_{O_2}}{Dt} = \mathcal{D}_{O_2} \nabla^2 C_{O_2} + \dot{r}_{O_2}.$$
(35)

Assumptions:

- The flow around the bubble is in the Stokes regime.
- Interface between the bubble and glass is fully mobile.
- Oxidation-reduction reaction of iron oxide is in chemical equilibrium<sup>14</sup>.
- Diffusion of iron is assumed very low.

$$\dot{r}_{\rm O_2} = -\frac{C_{\rm Fe}K_{\rm Fe}}{16C_{\rm O_2}^{3/4}(K_{\rm Fe} + C_{\rm O_2}^{1/4})^2} \frac{DC_{\rm O_2}}{Dt}.$$
(36)

<sup>&</sup>lt;sup>14</sup>R. G. C. Beerkens/H. de Waal: Mechanism of oxygen diffusion in glassmelts containing variable-valence ions, in: J. Am. Ceram. Soc. 73 (1990), pp. 1857–1861; F. Pigeonneau: Mass transfer of a rising bubble in molten glass with instantaneous oxidation-reduction reaction, in: Chem. Eng. Sci. 64.13 (2009), pp. 3120–3129.



Figure 14: Sh vs. Pe for  $O_2^{15}$ .

<sup>&</sup>lt;sup>15</sup>F. Pigeonneau/L. Pereira/A. Laplace: Mass transfer around a rising bubble in a glass-forming liquid involving oxidation-reduction reaction: Numerical computation of the Sherwood number, in: Chem. Eng. Sci. 232 (2021), p. 116382.

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Figure 15: *a* vs. *t* of O<sub>2</sub> bubble at T=1400 °C,  $C_{\rm Fe}=2.8 \times 10^{-2}$  wt%,  $\mathcal{R}_{\rm Fe}=0.575$ .



2. Mass transfer around a bubble

#### 3. Fining process

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#### 4. Synthesis





$$\frac{dn_i}{dt} = 2\pi a \mathcal{D}_i \operatorname{Sh}_i \left( C_i^{\infty} - \mathcal{L}_i P_i^{\beta_i} \right), \qquad (37)$$

$$\left[P_0 + \rho(H - z) + \frac{2\gamma}{a}\right] \frac{4\pi a^3}{3} = \left(\sum_{i=1}^N n_i\right) \mathcal{R}T,$$

$$\frac{d\mathbf{x}}{dt} = \mathbf{u} - \frac{\rho \mathbf{g} a^2}{3\eta}.$$
(38)





$$\frac{dn_i}{dt} = 2\pi a \mathcal{D}_i \operatorname{Sh}_i \left( \frac{C_i^{\infty} - \mathcal{L}_i P_i^{\beta_i}}{dt} \right),$$
(37)

$$\left[P_{0} + \rho(H - z) + \frac{2\gamma}{a}\right] \frac{4\pi a^{3}}{3} = \left(\sum_{i=1}^{N} n_{i}\right) \mathcal{R}T, \qquad (38)$$
$$\frac{d\mathbf{x}}{dt} = \mathbf{u} - \frac{\rho \mathbf{g} a^{2}}{3\eta}. \qquad (39)$$





Multivalent elements oxidation-reduction state described by

$$M^{(m+k)+}(l) + \frac{k}{2}O^{2-}(l) \longleftrightarrow M^{m+}(l) + \frac{k}{4}O_2(g),$$
 (40)

$$\mathrm{SO}_4^{2+}(I) \longleftrightarrow \mathrm{SO}_2(g) + \frac{1}{2}\mathrm{O}_2(g) + \mathrm{O}^{2-}(I).$$
 (41)

Equilibrium with a gaseous atmosphere is given by

$$O_2(g) \longleftrightarrow O_2(l).$$
 (42)









Figure 16: (a) Ion concentrations and (b) redox state  $C_{\rm Fe^{2+}}/C_{\rm Fe}$  without and with thermodynamic equilibrium of gaseous phase.

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#### 3.1 Sulphate fining





Figure 17: Bubble histogram vs. 2a [SGR, Paris, N. McDonald].

Model coupling mass transfer between bubble and the liquid<sup>16</sup>:

$$\frac{dC_{A_{i}}}{dt} = \sum_{r=1}^{R} \nu_{ri} \frac{d\zeta_{r}}{dt}, \ i \in [1, N_{I}]. \ (43) \qquad \qquad S_{b,G_{j}} = -4\pi \sum_{k=1}^{N_{cl}} a_{k}^{2} k_{G_{j},k} \left(C_{G_{j}} - \mathcal{L}_{G_{j}} P_{G_{j},k}^{\alpha_{G_{j}}}\right) N_{b,k}, \ (46) \\ \frac{dC_{G_{j}}}{dt} = \sum_{r=1}^{R} \beta_{rj} \frac{d\zeta_{r}}{dt} + S_{b,G_{j}}, \ j \in [1, N_{fg}], \ (44) \qquad \qquad \sum_{k=1}^{R} M_{rk} \frac{d\zeta_{k}}{dt} = \frac{d \ln K_{r}}{dT} \frac{dT}{dt} - \sum_{j=1}^{N_{cl}} \frac{\beta_{rj}}{C_{G_{j}}} S_{b,G_{j}}, \ r \in [1; R], \ (47) \\ \frac{dC_{G_{j}}}{dt} = S_{b,G_{j}}, \ j \in [N_{fg} + 1; N_{g}]. \ (45) \qquad \qquad M_{rk} = \sum_{i=1}^{N_{l}} \frac{\nu_{ri}\nu_{ki}}{C_{A_{i}}} + \sum_{j=1}^{N_{fg}} \frac{\beta_{rj}\beta_{kj}}{C_{G_{j}}}. \ (48)$$

<sup>&</sup>lt;sup>16</sup>J. Kloužek et al.: The redox distribution at the interface of glass melts with different oxidation state, in: Ceram. Silik. 44 (2000), pp. 91–95; F. Pigeonneau: Coupled modelling of redox reactions and glass melt fining processes, in: Glass Technol.: Eur. J. Glass Sci. Technol. A 48.2 (2007), pp. 66–72.



Figure 18: Amount of H<sub>2</sub>O vs.  $x_{H_2O}$  with a solubility equal to  $0.68e^{-613/T}$  at  $P=10^5$  Pa according Beerkens<sup>17</sup>.

<sup>&</sup>lt;sup>17</sup>R. G. C. Beerkens: Analysis of advanced and fast fining processes for glass melts, in: Advances in Fusion and Processing of Glass III, New York 2004, pp. 3–24.



Figure 19: a vs. t with  $a(0)=10^2 \mu m$  and  $T=1400 \circ C$ .





Patent proposed by Lazet<sup>18</sup>.

Class	Ini. compo. ini.	$N_{b}~({ m m}^{-3})$	<i>a</i> 0 (mm)
1	CO <sub>2</sub>	$2.25 imes10^4$	0.1
2	$CO_2$	$2.25 imes10^4$	0.25
3	$CO_2$	$2.25 imes10^4$	0.75
4	$CO_2$	$2.25 imes10^4$	1
5	$H_2O$	10 <sup>5</sup>	10

Table 1: Composition, volume concentration and initial radii of bubble classes.

<sup>&</sup>lt;sup>18</sup>F. J. Lazet: Preparing alkali metal silicate glass with bubbles, US Patent 3,960,532 A, Philadelphia Quartz Company, 1975.



Figure 20: *a* vs. *t* for the 5 classes of bubbles at T=1200 °C.





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Figure 21: Molar concentration of dissolved gas in the liquid vs. t at T=1200 °C. PSL® | Genef



# **3. Fining process** 3.3 Helium fining

- Patent proposed by Kobayashi et al.<sup>19</sup>.
- ► He bubbling ➤ diffusion of He in the liquid ➤ migration of He in bubbles produced by the melting.
- ▶ Test with two populations of bubbles, one from melting with  $10^8 \text{ m}^{-3}$  and one due to the bubbling of He with  $10^7 \text{ m}^{-3}$  at  $T=1400 \circ \text{C}$ .
- Initial radius equal to 100 μm.

<sup>&</sup>lt;sup>19</sup>H. Kobayashi/S. E. Jaynes/R. G. C. Beerkens: Process of fining glassmelts using helium bubbles, US Patent 2006/0174655A1, Praxair, Inc., 2003.

3.3 Helium fining



Figure 22: *a* vs. *t* for the 2 classes of bubbles at T=1400 °C.





3.3 Helium fining

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Figure 23: Molar concentration of He dissolved in the liquid vs. t at T=1400 °C.



3.4 Centrifugal and low pressure fining



Figure 24: Centrifugal finer according to Spinosa<sup>20</sup>.

<sup>&</sup>lt;sup>20</sup>E. D. Spinosa: Modular refining methods, in: Am. Ceram. Soc. Bull. 83.10 (2004), pp. 25–27.

3.4 Centrifugal and low pressure fining

Tonarová et al.<sup>21</sup> provided an optimisation of the centrifugal fining.

$$P = P_0 + \rho g \left( H - \frac{\omega^2 R^2}{4g} - z \right) + \rho \frac{\omega^2 r^2}{2}.$$
(49)

<sup>&</sup>lt;sup>21</sup>V. Tonarová/L. Němec/J. Kloužek: The optimal parameters of bubble centrifuging in glass melts, in: J. Non-Cryst. Solids 357 (2011), pp. 3785–3790.

3.4 Centrifugal and low pressure fining



Figure 25: Fining time vs.  $\omega$  according to Tonarová et al.<sup>22</sup>.

<sup>&</sup>lt;sup>22</sup>Tonarová/Němec/Kloužek: The optimal parameters of bubble centrifuging in glass melts (see n. 21).

3.4 Centrifugal and low pressure fining



Figure 26: *a* vs. *t* for various  $\omega$  for 3 atmospheric pressure.





3.4 Centrifugal and low pressure fining



Figure 27: Fining time vs.  $\omega$  for 3 atmospheric pressure.





## 4. Synthesis

Method	Structure	Installation	Operation	Development
Helium	Small	Familiar to industry	Complicated	Laboratory
Sonic	Small	Familiar to industry	Simple	Laboratory
Centrifugal	Large	Unfamiliar to industry	Extremely challenging	Pilot scale
Vacuum	Large	Unfamiliar to industry	Challenging	Commercial

Table 2: Comparison of refining methods<sup>23</sup>.

<sup>&</sup>lt;sup>23</sup>Spinosa: Modular refining methods (see n. 20).

### 4. Synthesis





Figure 28: Segmented melter RAMAR of Owens-Illinois<sup>a</sup>.

<sup>&</sup>lt;sup>a</sup>F. G. Pellett et al.: Method for rapid melting and refining glass, US Patent 3,819,350, 1974.

#### Thank you all of you for your attention!

Students:

H. Kočárková, M. Perrodin, M. Guémas, D. Boloré, L. Pereira.

Colleagues:

- SGR, Paris: M.-H. Chopinet, E. Gouillart, D. Martin, N. McDonald;
- Lab. Navier (Univ. Paris-Est/Marne la Vallée): F. Rouyer;
- Lab. Génie Chim. (Toulouse): P. Chamelot, O. Masbernat;
- ▶ I.M.F.T. (Toulouse): E. Climent;
- LadHyX (Ecole Polytechnique, Paris Saclay): A. Sellier;
- CEA Marcoule: A. Laplace;
- Lab. Inorg. Mater., Univ. of Prague (Czech Republic): J. Kloužek.

#### Contact:

franck.pigeonneau@minesparis.psl.eu



