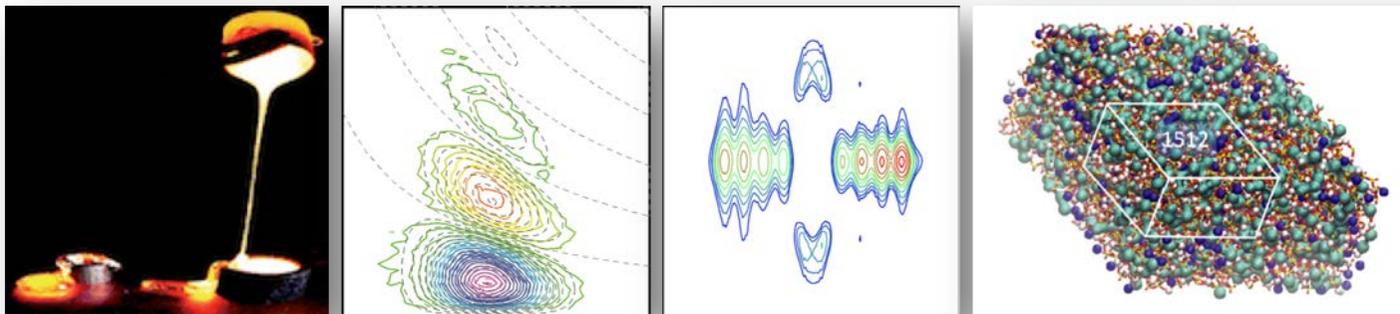


Solid-State Nuclear Magnetic Resonance of glasses: from basics to more advanced concepts

Dominique Massiot, Pierre Florian

CEMHTI-CNRS UPR3079 Orléans

<http://www.cemhti.cnrs-orleans.fr/?nom=massiot>



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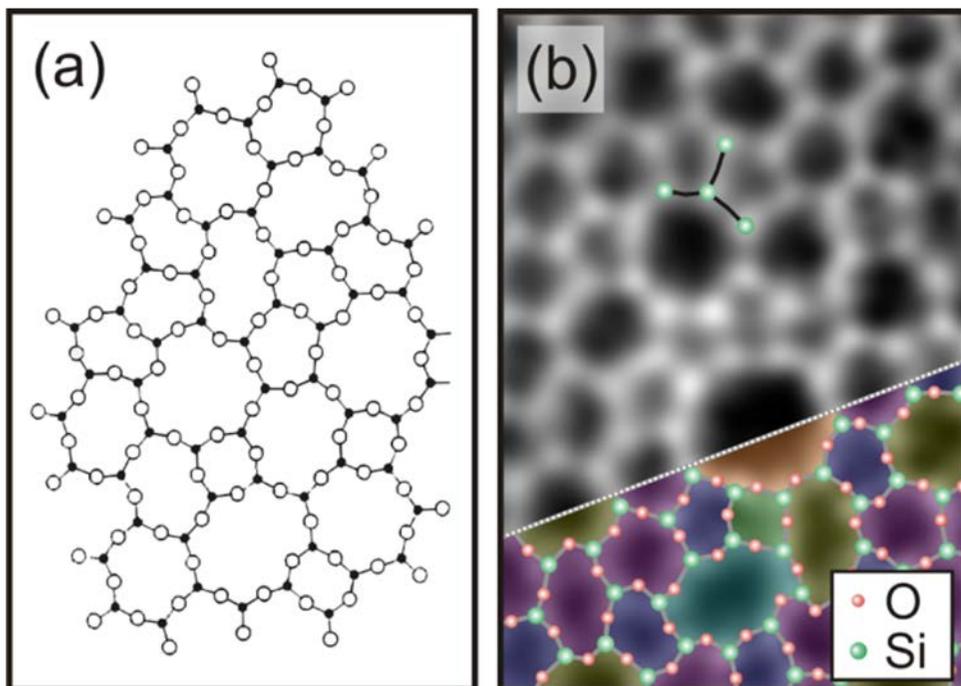
Amorphous materials: Order within disorder

Phillip S. Salmon

Nature Material **1**, 87 - 88 (2002)

doi:10.1038/nmat737

Chemical order / Geometrical disorder



THE JOURNAL OF
PHYSICAL CHEMISTRY C

Article

pubs.acs.org/JPC

Atomic Arrangement in Two-Dimensional Silica: From Crystalline to Vitreous Structures

Leonid Lichtenstein, Markus Heyde,* and Hans-Joachim Freund

NANO LETTERS

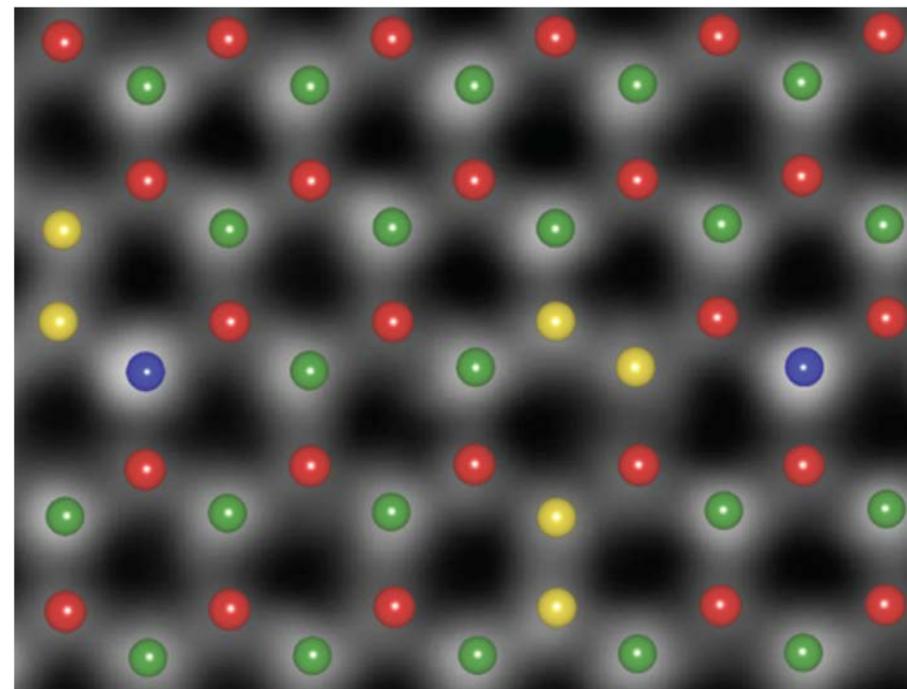
Letter

pubs.acs.org/NanoLett

Direct Imaging of a Two-Dimensional Silica Glass on Graphene

Pinshane Y. Huang,[†] Simon Kurasch,[‡] Anchal Srivastava,^{§,¶} Viera Skakalova,^{§,||} Jani Kotakoski,^{||,⊥} Arkady V. Krasheninnikov,^{⊥,¶} Robert Hovden,[†] Qingyun Mao,[†] Jannik C. Meyer,^{‡,||} Jurgen Smet,[§] David A. Muller,^{*,†,□} and Ute Kaiser^{*,‡}

Geometrical order / Chemical disorder



Vol 464 | 25 March 2010 | doi:10.1038/nature08879

nature

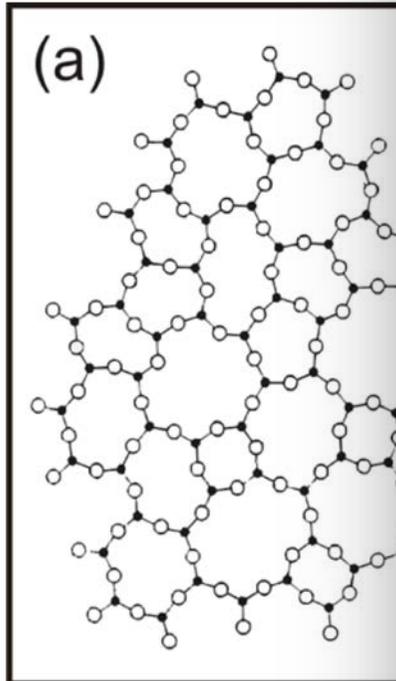
nature

LETTERS

Atom-by-atom structural and chemical analysis by annular dark-field electron microscopy

Ondrej L. Krivanek¹, Matthew F. Chisholm², Valeria Nicolosi³, Timothy J. Pennycook^{2,4}, George J. Corbin¹, Niklas Dellby¹, Matthew F. Murfitt¹, Christopher S. Own¹, Zoltan S. Szilagyi¹, Mark P. Oxley^{2,4}, Sokrates T. Pantelides^{2,4} & Stephen J. Pennycook^{2,4}

Chemical order



THE JOURNAL OF
PHYSICAL CHEMISTRY C

Atomic Arrangement in Two Vitreous Structures

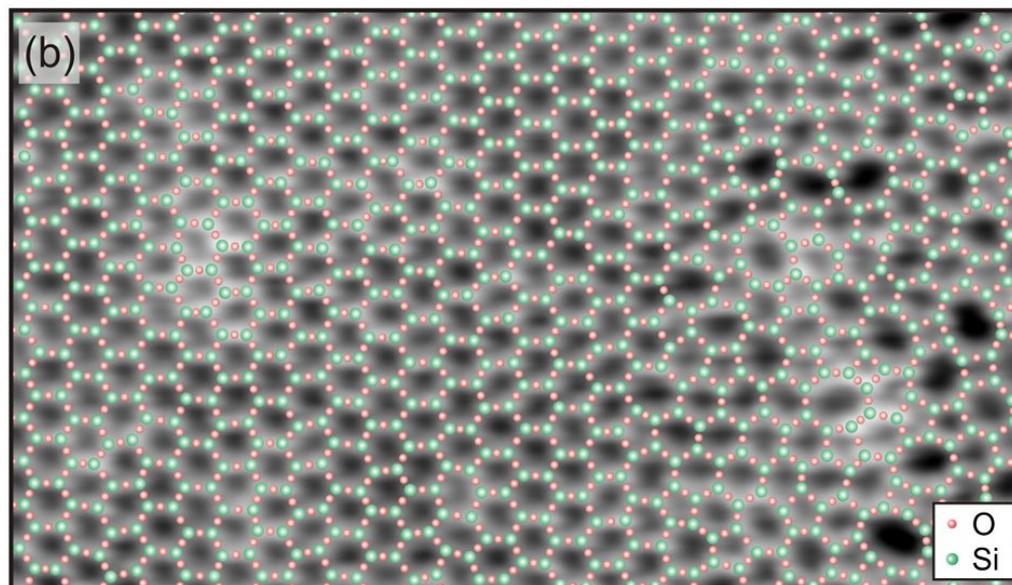
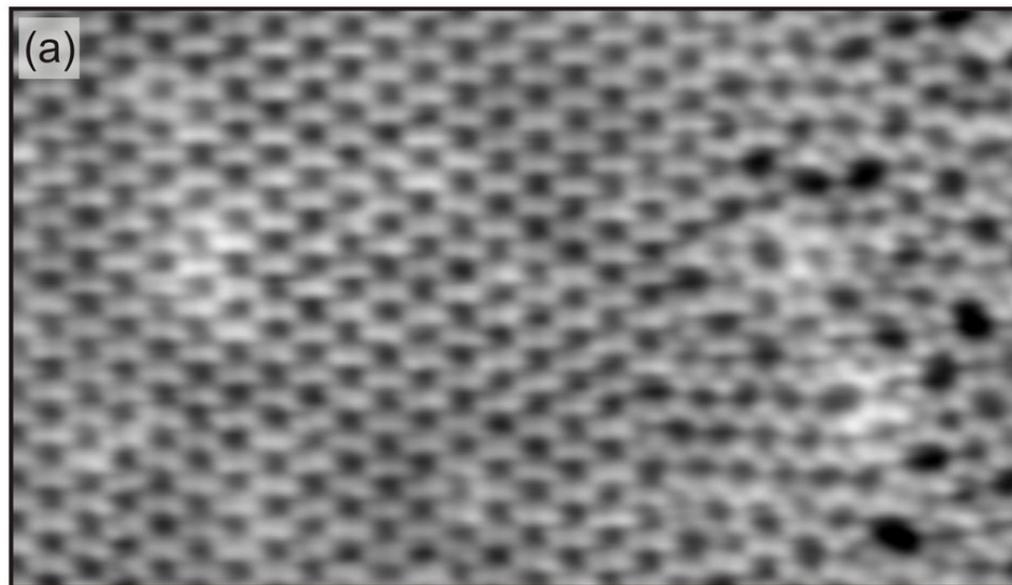
Leonid Lichtenstein, Markus Heyde,* and

NANO LETTERS

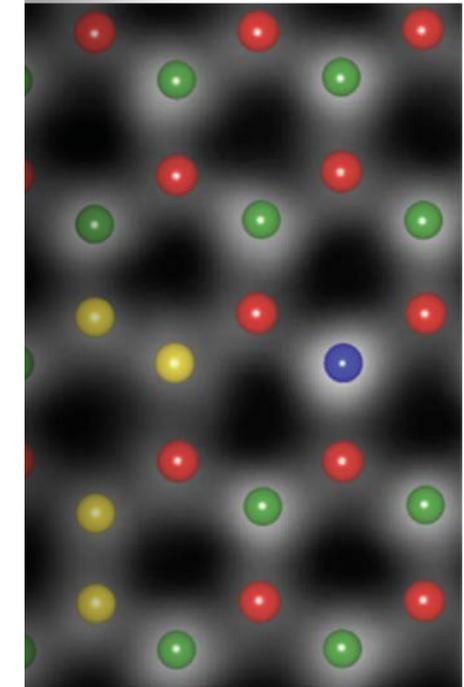
Direct Imaging of a Two-D

Pinshane Y. Huang,[†] Simon Kurasch,[‡]
Arkady V. Krasheninnikov,^{1,§} Robert He
David A. Muller,^{*;‡,□} and Ute Kaiser^{*;‡}

© Dominique Massiot



Chemical disorder

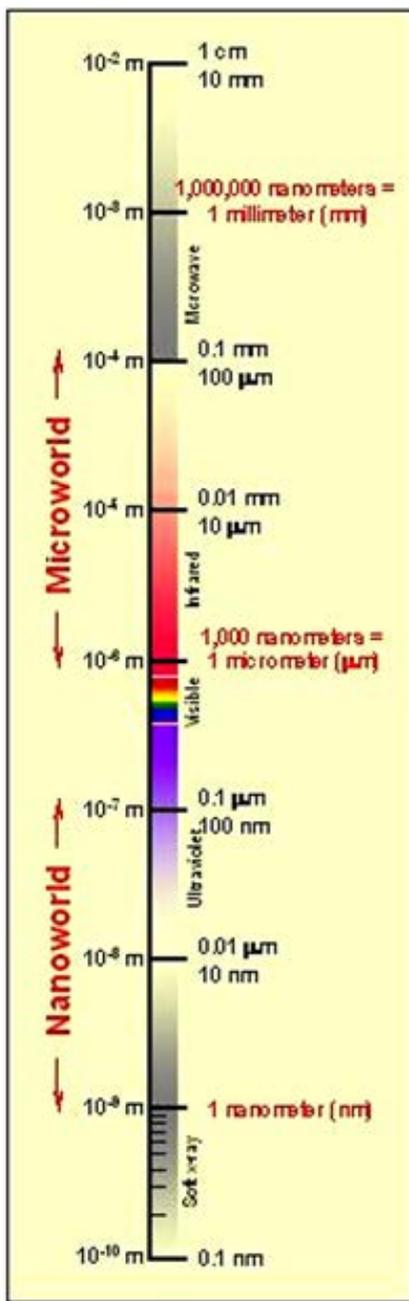


nature

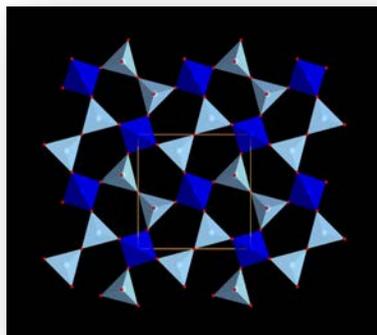
LETTERS

and chemical analysis by microscopy

si³, Timothy J. Pennycook^{2,4}, George J. Corbin¹,
Zoltan S. Szilagy¹, Mark P. Oxley^{2,4},



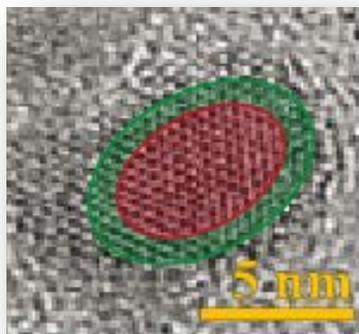
Diffraction X-Rays Neutrons e⁻



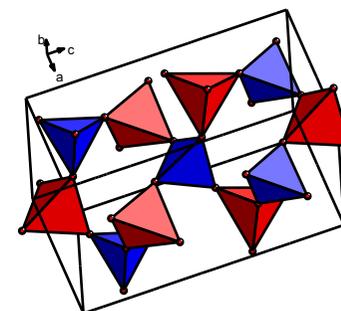
Order



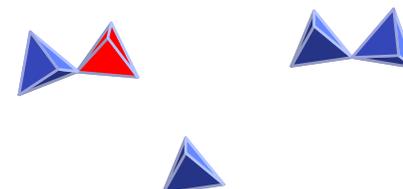
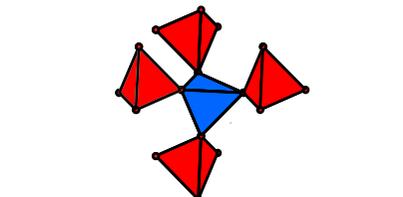
Disorder



Local Order

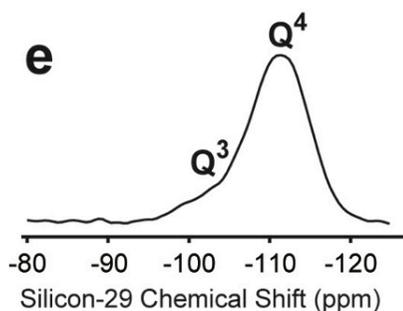
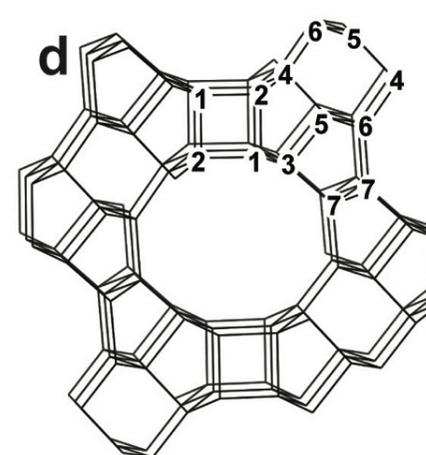
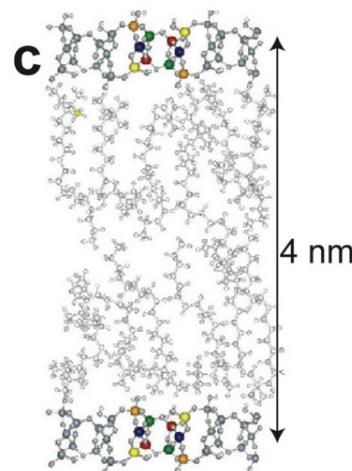
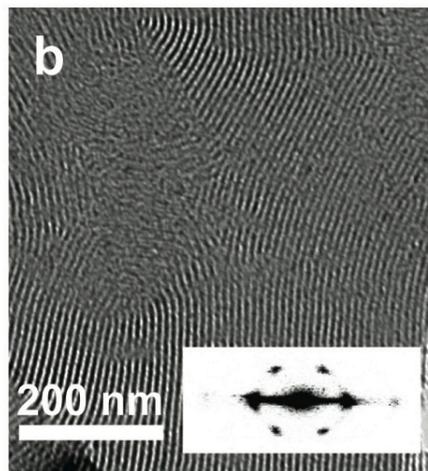
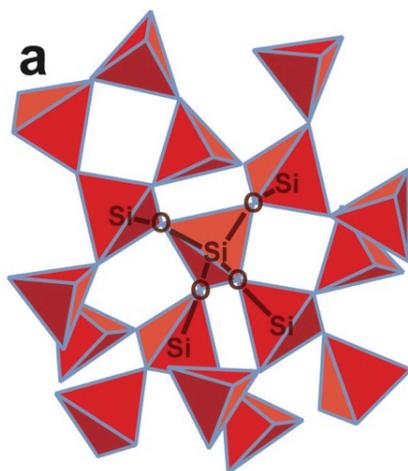


nm

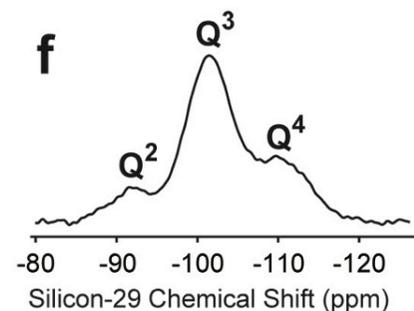


Å

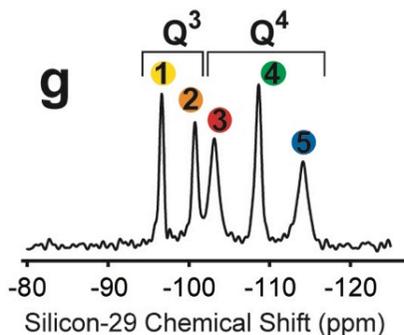
NMR XAS Raman / IR



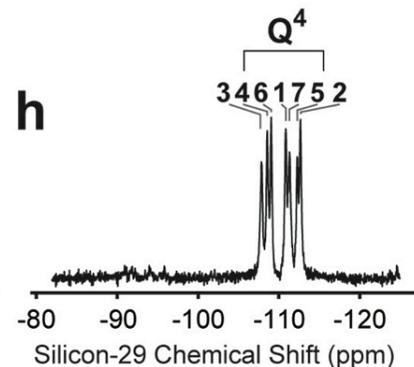
Glass



Mesoporous Silica



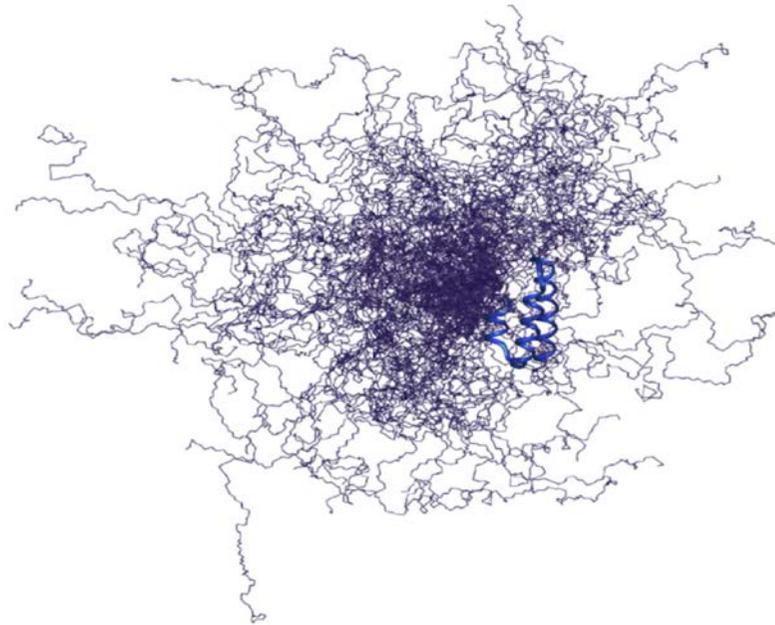
SiO_2 -surfactant
Mesophase



Zeolite

*Resolution is gained by averaging out anisotropic signatures
Magic Angle Spinning MAS*

Intrinsically Disordered Protein IDPs



Macromolecules

secondary and tertiary

Structure & dynamics

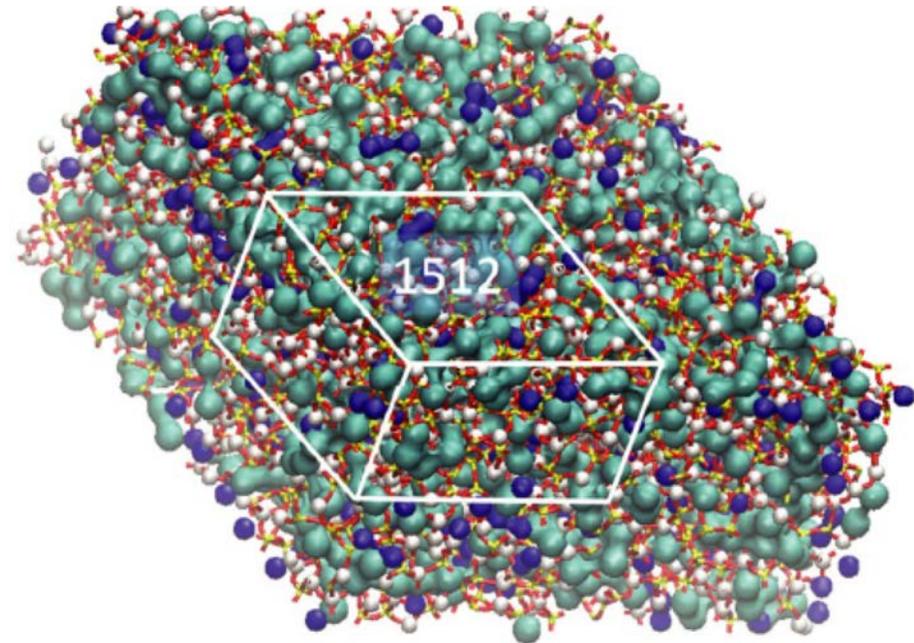
Constraints from identified spin pairs

Through-bond and through space

^1H / ^{13}C / ^{15}N

Glass

3D interconnected network



Multicomponent Glass SiO₂-Al₂O₃ MO

Extended 3D *rigid* network of tetrahedral

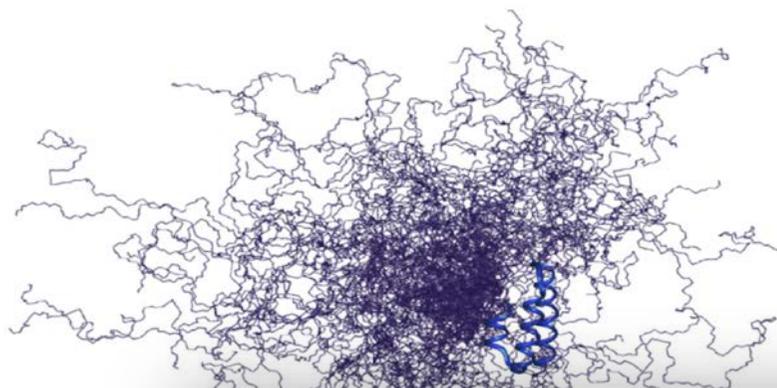
Connectivity

Very similar information from

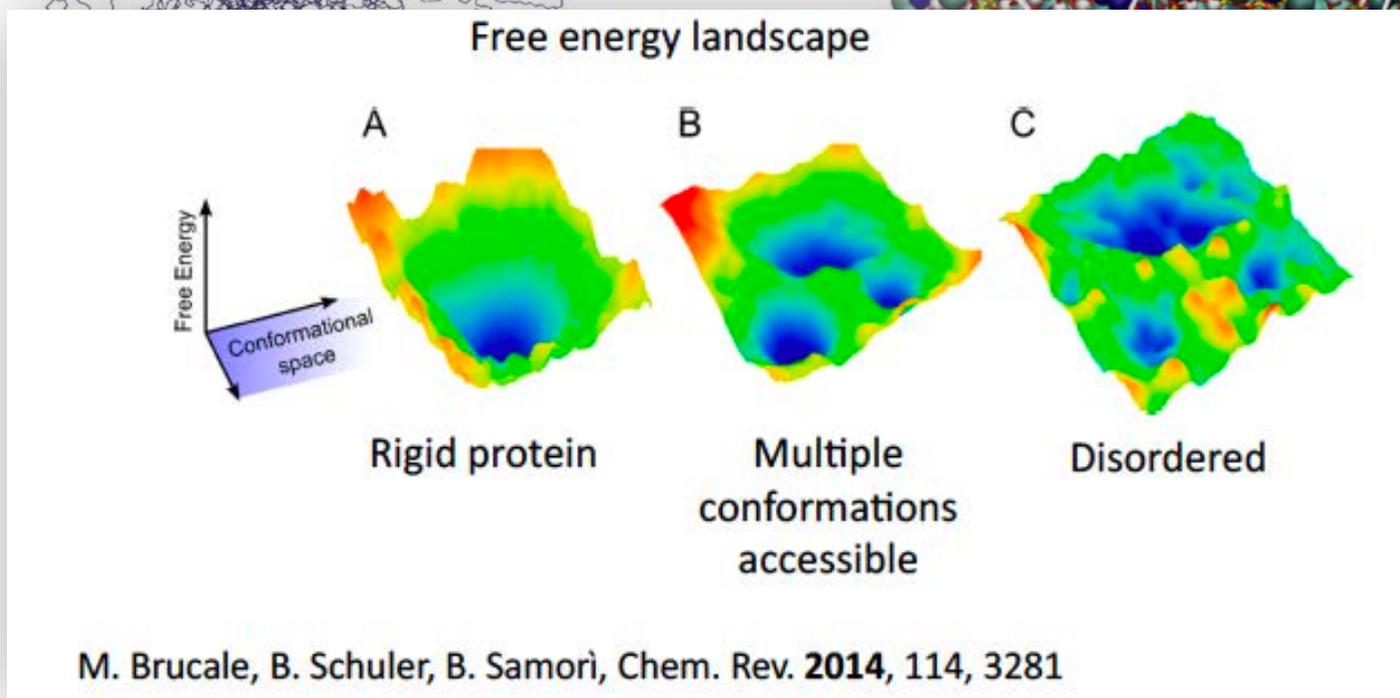
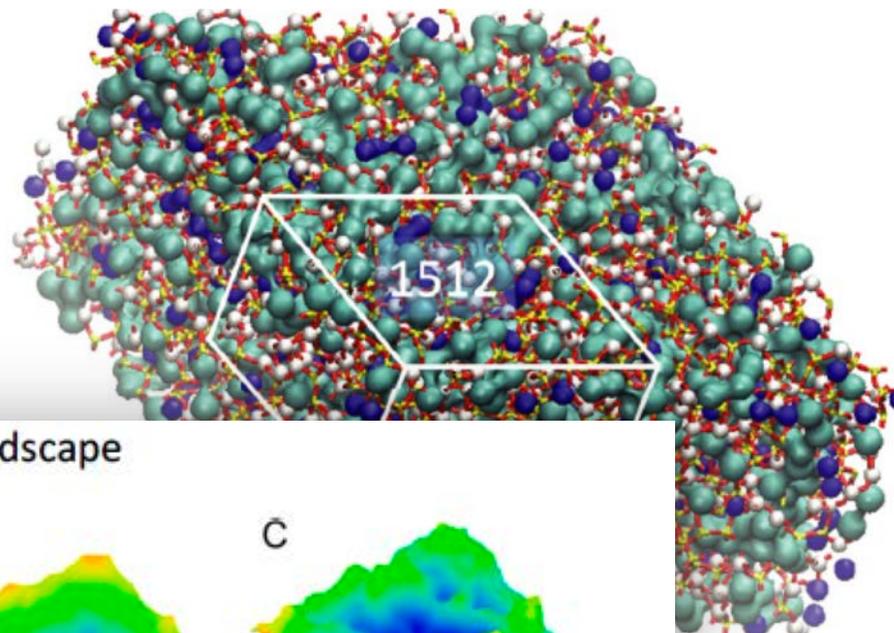
Dipolar and J-coupling

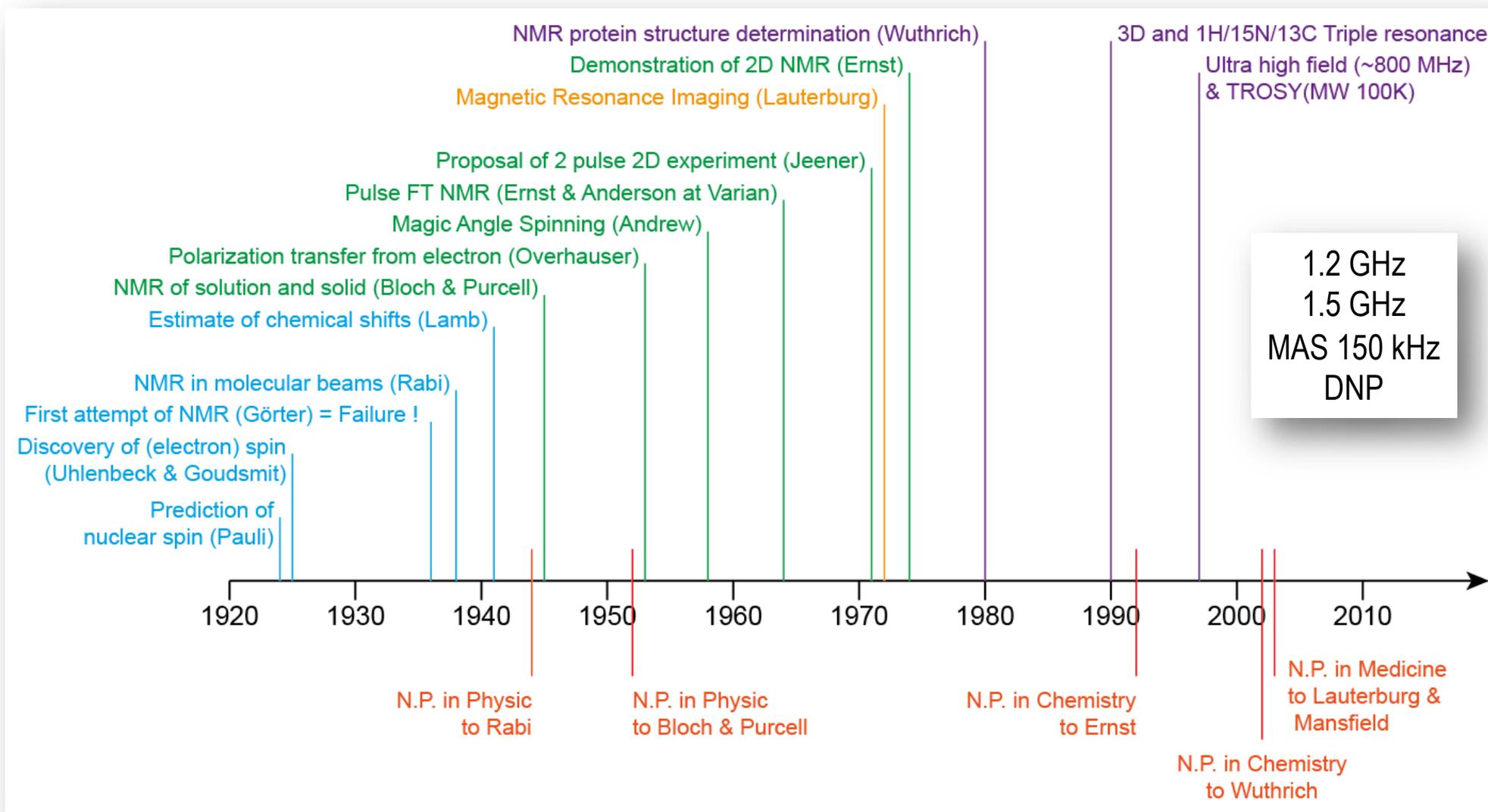
^{29}Si / ^{23}Na / ^{27}Al / ^{43}Ca / ^{17}O

Intrinsincally Disordered Protein IDPs



Glass 3D interconnected network





© Pierre Florian

The Nuclear Magnetic Resonance

The Nobel Prize in Physics 1944

Isidor Isaac Rabi

by Professor E. Hulthén, Stockholm

Let us now for a moment touch upon Rabi's achievements in this field. Returning to the essential point of the problem, let us put the question: How does the atom react to the magnetic field? According to a theorem stated by the English mathematician Larmor, this influence may be ascribed to a relatively slow precession movement on the part of the electron and the atomic nucleus around the field direction - a gyromagnetic effect most closely recalling the gyroscopic movement performed by a top when it spins around the vertical line. If the strength of the magnetic field is known, the magnetic factor of the electron and of the atomic nucleus can also be estimated by this means, provided that we can observe and measure these precessional frequencies. Rabi solved the problem in a manner as simple as it was brilliant. Within the magnetic field was inserted a loop of wire, attached to an oscillating circuit the frequency of which could be varied in the same manner as we tune in our radio receiving set to a given wavelength. Now, when the atomic beam passes through the magnetic field, the atoms are only influenced on condition that they precess in time with the electric current in the oscillating circuit. This influence might perhaps be described graphically: the nucleus performs a vault (salto) - the technical term for which is a "quantum jump" - thereby landing in another positional direction to the field. But this means that the atom has lost all chance of reaching the detector and of being registered by it. The effect of these quantum jumps is observable by the fact that the detector registers a marked resonance minimum, the frequency position of the registration being determined with the extraordinary precision achievable with the radio frequency gauge. By this method Rabi has literally established radio relations with the most subtle particles of matter, with the world of the electron and of the atomic nucleus.

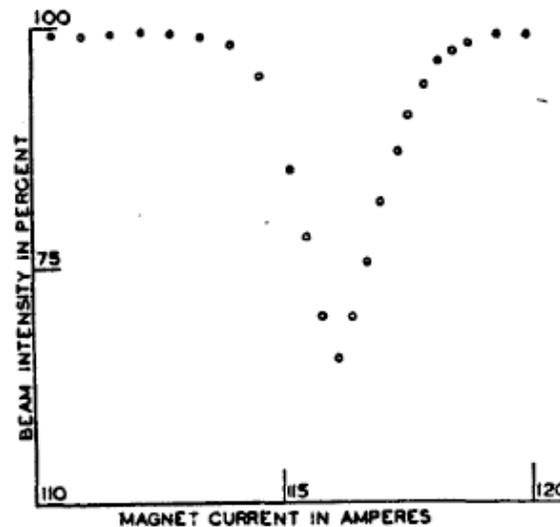
The Nuclear Magnetic Resonance

The Nobel Prize in Physics 1944

Isidor Isaac Rabi

by Professor E. Hulthén, Stockholm

Let us now for a moment touch upon the problem, let us put the question in the form of a theorem stated by the English physicist Rabi. precession movement on the gyromagnetic effect most closely around the vertical line. If the angular momentum and of the atomic nucleus can be compared to a top, these precessional frequencies are proportional to the magnetic field. When a magnetic field was inserted and varied in the same manner as the frequency of the atomic beam passes through the magnetic field, the nucleus precesses in time with the electric field. graphically: the nucleus performs a quantum jump, landing in another position.

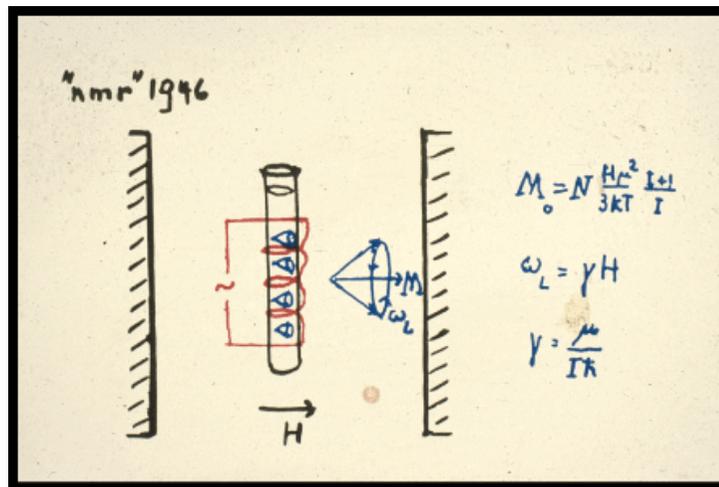


Resonance of LiCl from Rabi's 1938 paper.

According to the essential point of the problem? According to a theorem stated by the English physicist Rabi. described to a relatively slow precession movement on the gyromagnetic effect most closely around the vertical line - a If the angular momentum and of the atomic nucleus can be compared to a top when it spins, these precessional frequencies are proportional to the magnetic field. When a magnetic field was inserted and varied in the same manner as the frequency of the atomic beam passes through the magnetic field, the nucleus precesses in time with the electric field. graphically: the nucleus performs a quantum jump - thereby losing all chance of

reaching the detector and of being registered by it. The effect of these quantum jumps is observable by the fact that the detector registers a marked resonance minimum, the frequency position of the registration being determined with the extraordinary precision achievable with the radio frequency gauge. By this method Rabi has literally established radio relations with the most subtle particles of matter, with the world of the electron and of the atomic nucleus.

E.M. Purcell et F. Bloch, Nobel of Physics 1952 for the discovery of NMR



$$M_0 = N \frac{\mu^2 \hbar^2 I(I+1)}{3kT}$$

$$\omega_L = \gamma H$$

$$\gamma = \frac{\mu}{\hbar}$$

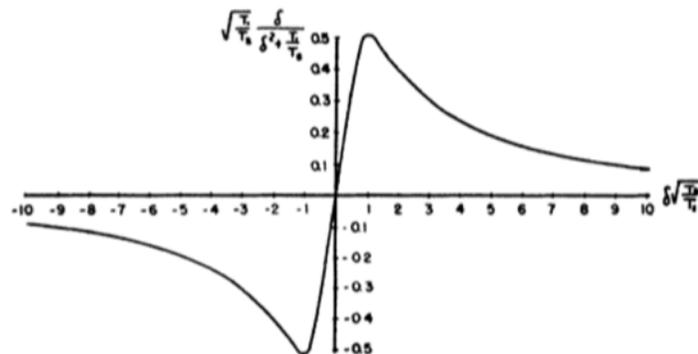


FIG. 2. Schematic representation of the voltage amplitude in the case of slow passage. T_1 and T_2 are the longitudinal and transversal relaxation times, respectively, and the scale is chosen such as to make the plot independent of their values. The significance of abscissa and ordinate is otherwise the same as in Fig. 1.

Nuclear Induction

F. BLOCH, W. W. HANSEN, AND MARTIN PACKARD
Stanford University, Stanford University, California
 January 29, 1946

THE nuclear magnetic moments of a substance in a constant magnetic field would be expected to give rise to a small paramagnetic polarization, provided thermal equilibrium be established, or at least approached. By superposing on the constant field (z direction) an oscillating magnetic field in the x direction, the polarization, originally parallel to the constant field, will be forced to precess about that field with a latitude which decreases as the frequency of the oscillating field approaches the Larmor frequency. For frequencies near this magnetic resonance frequency one can, therefore, expect an oscillating induced voltage in a pick-up coil with axis parallel to the y direction. Simple calculation shows that with reasonable apparatus dimensions the signal power from the pick-up coil will be substantially larger than the thermal noise power in a practicable frequency band.

We have established this new effect using water at room temperature and observing the signal induced in a coil by the rotation of the proton moments. In some of the experiments paramagnetic catalysts were used to accelerate the establishment of thermal equilibrium.

H 2																	He
Li 6,7	Be 9											B 10,11	C	N 14	O 17	F	Ne 21
Na 23	Mg 25											Al 27	Si	P	S 33	Cl 35,37	Ar
K 39,41	Ca 43	Sc 45	Ti 47,49	V 50,51	Cr 53	Mn 55	Fe	Co 59	Ni 61	Cu 63,65	Zn 67	Ga 69,71	Ge 73	As 75	Se	Br 79,81	Kr 83
Rb 85,87	Sr 87	Y	Zr 91	Nb 93	Mo 97	Tc	Ru 99,101	Rh	Pd 105	Ag	Cd	In 113-5	Sn	Sb 121-3	Te	I 127	Xe 129-31
Cs 133	Ba 135-7	La 138-9	Hf 177-9	Ta 181	W	Re 185-7	Os 187-9	Ir 191-3	Pt	Au 197	Hg 201	Tl	Pb	Bi	Po	At	Rn
Fr	Ra	Ac															

Ce	Pr 141	Nd 143-5	Pm	Sm 147-9	Eu 151-3	Gd 155-7	Tb 159	Dy 161-3	Ho 165	Er 167	Tm	Yb 173	Lu 175-6
Th	Pa	U 235	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr

Element n'ayant que des isotopes de spin $I=1/2$

Element ayant des isotopes de spin quadripolaire ($I>1/2$)

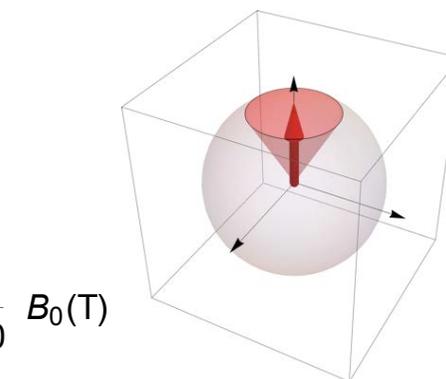
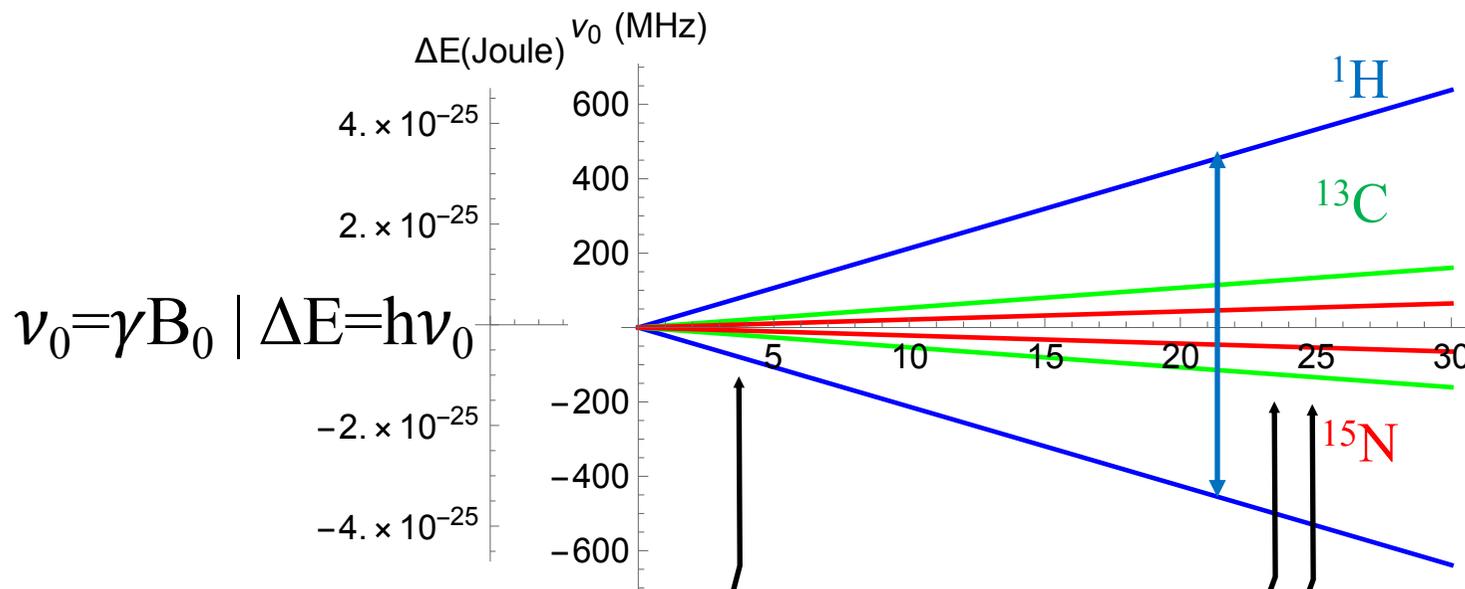
❖ Observability

- ❖ abundance
- ❖ Gyromagnetic ratio
- ❖ Quadrupolar momentum

Numerous possibly sensitive nuclei but fewer easily observed

The most usually observed are «light» nuclei

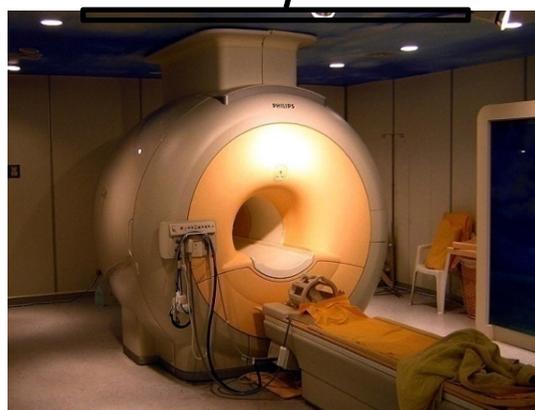
- ❖ $I=1/2$: ^1H , ^{29}Si , ^{31}P , ^{13}C ,
- ❖ $I=3/2$: ^{23}Na , ^{11}B
- ❖ $I=5/2$: ^{27}Al , ^{17}O (isotopic enrichment)



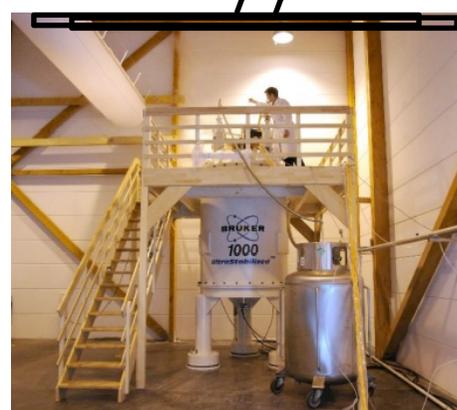
NMR signal
 $\Delta N \propto \exp(-\Delta E/kT)$



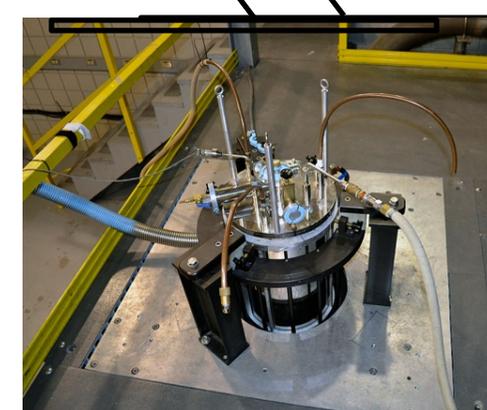
Earth Field
 $\sim 50\mu\text{T}$ (0.5G)
 magnetometer



3T [128MHz]
 Whole body (Philips)

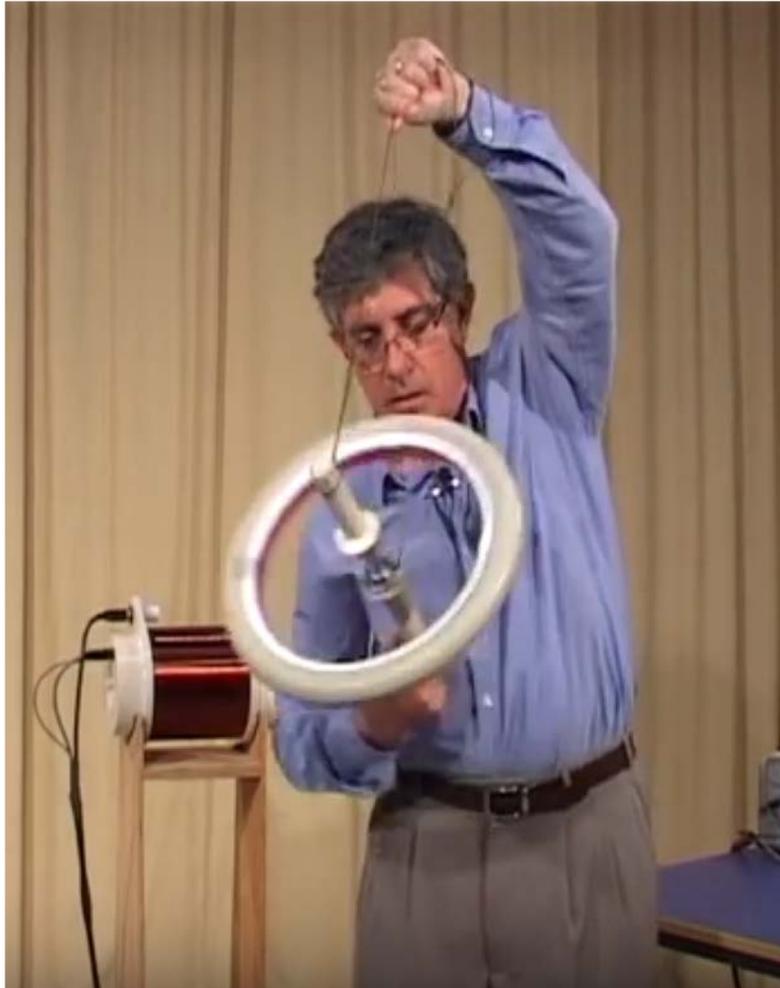


23.4T
 1GHz (Lyon)
 1.2GHz Bruker 2019



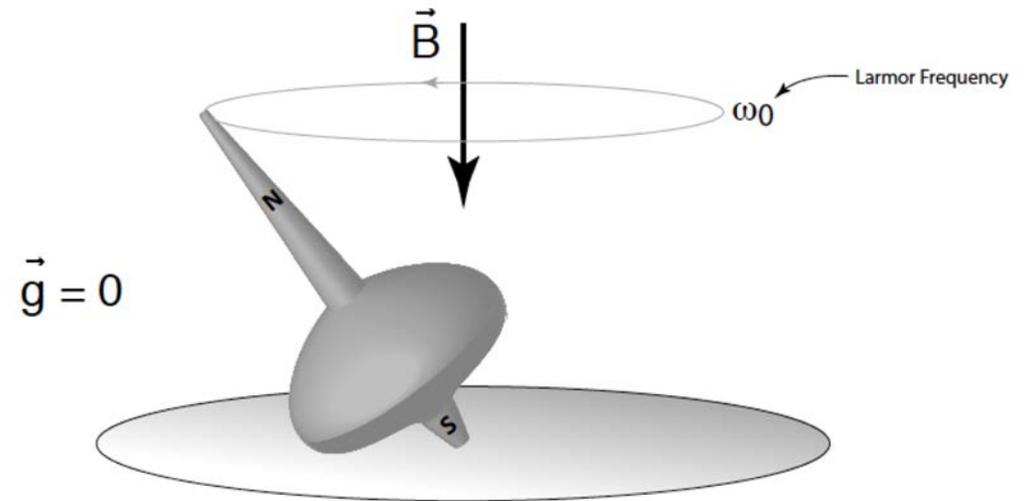
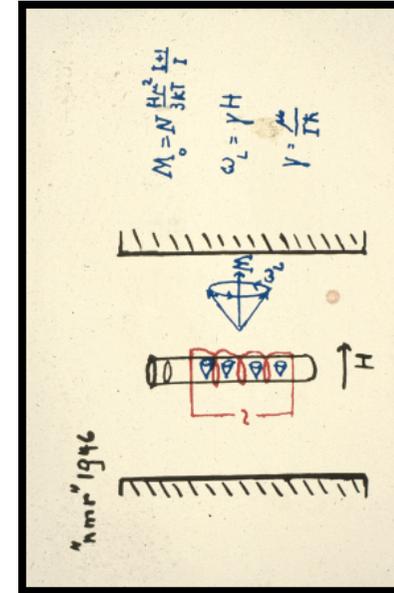
35 T
 1.5 GHz (Thallahassee)
 40T inhomogeneous

Top in a Gravitational field

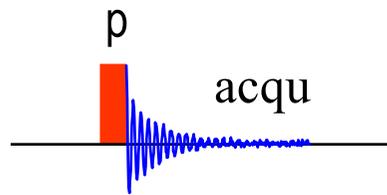
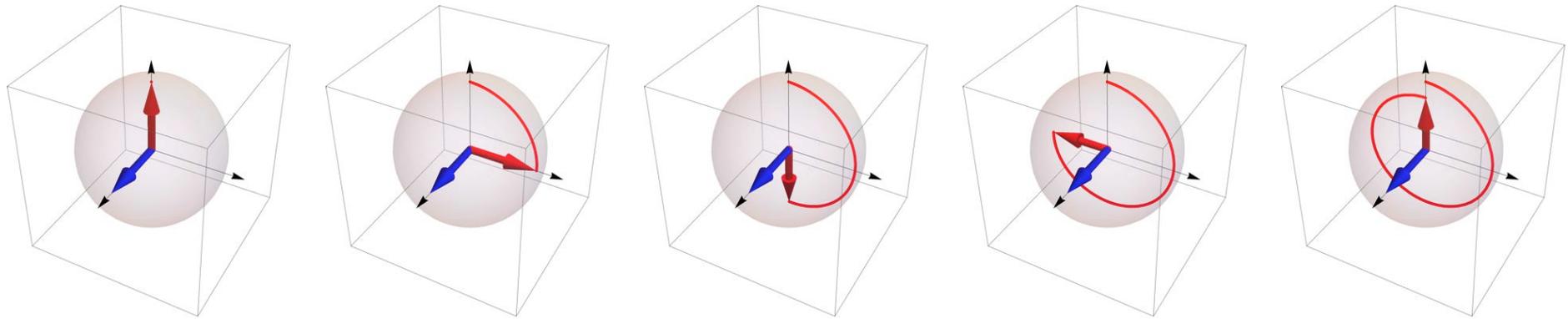


Paul Callaghan applying a pulse in the gravitational field
<https://www.youtube.com/watch?v=7aRKAXD4dAg>

"Spin" in Magnetic Field



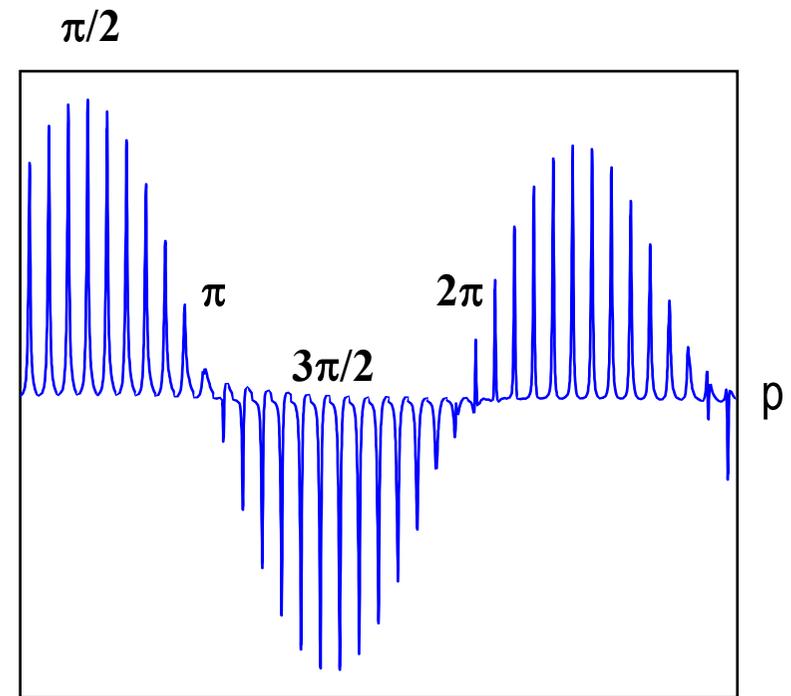
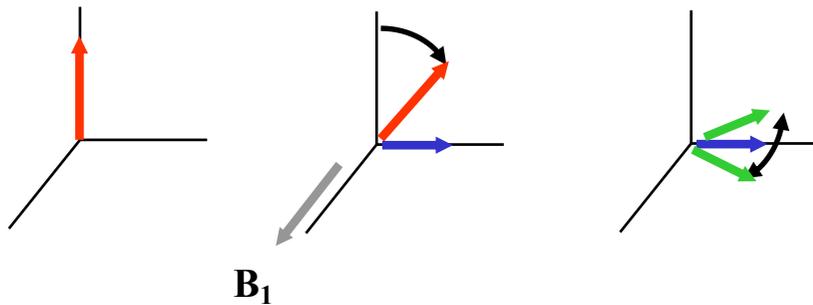
Rotating Frame



Equilibrium

Pulse

Evolution (t_2)



Precession in the B_1 field [$\nu_1 = \gamma B_1$] – remaining far from thermal equilibrium

Bloch Equations

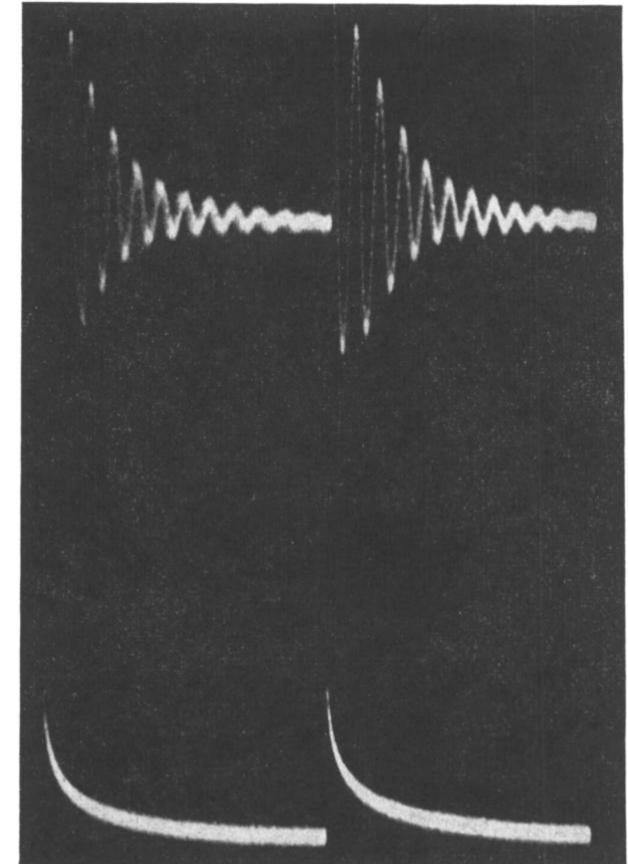
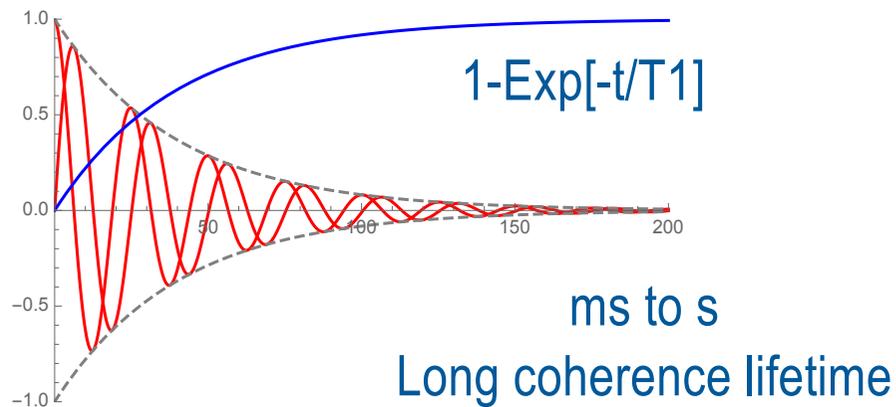
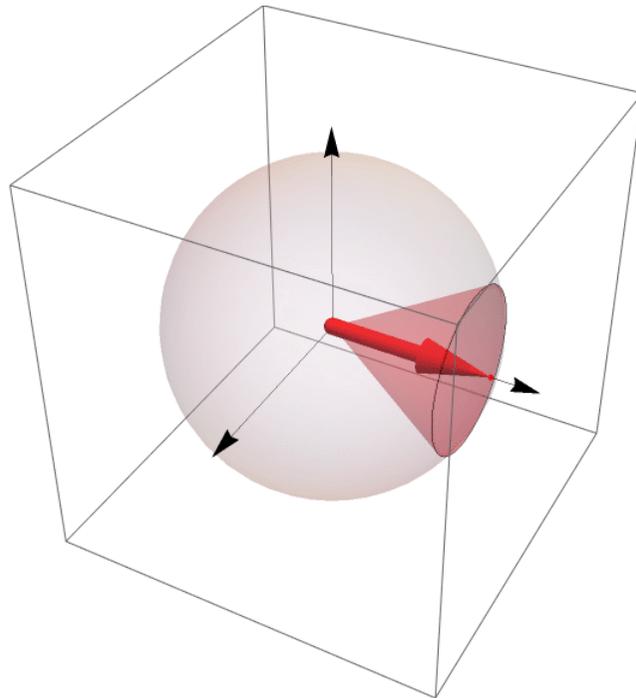
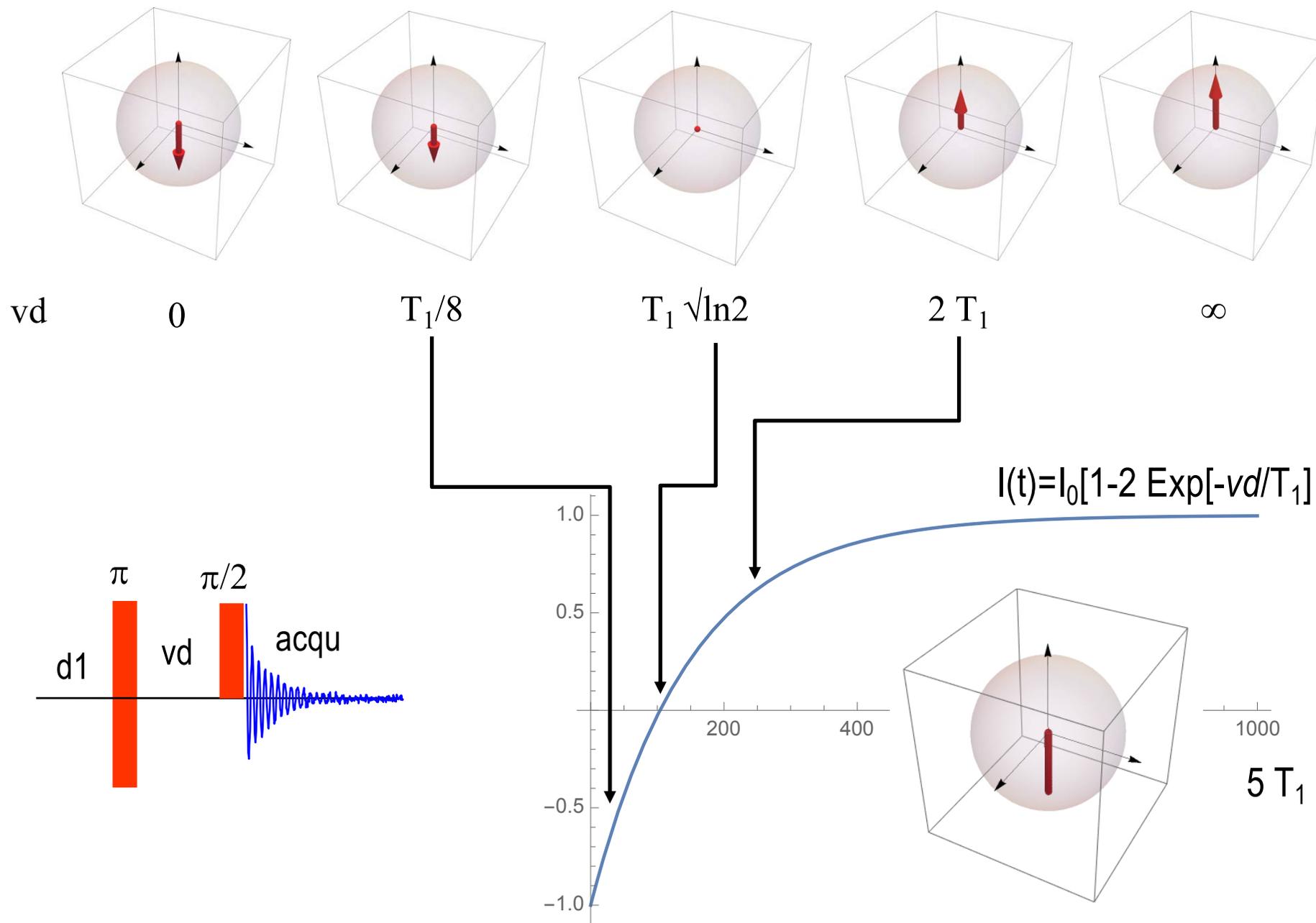
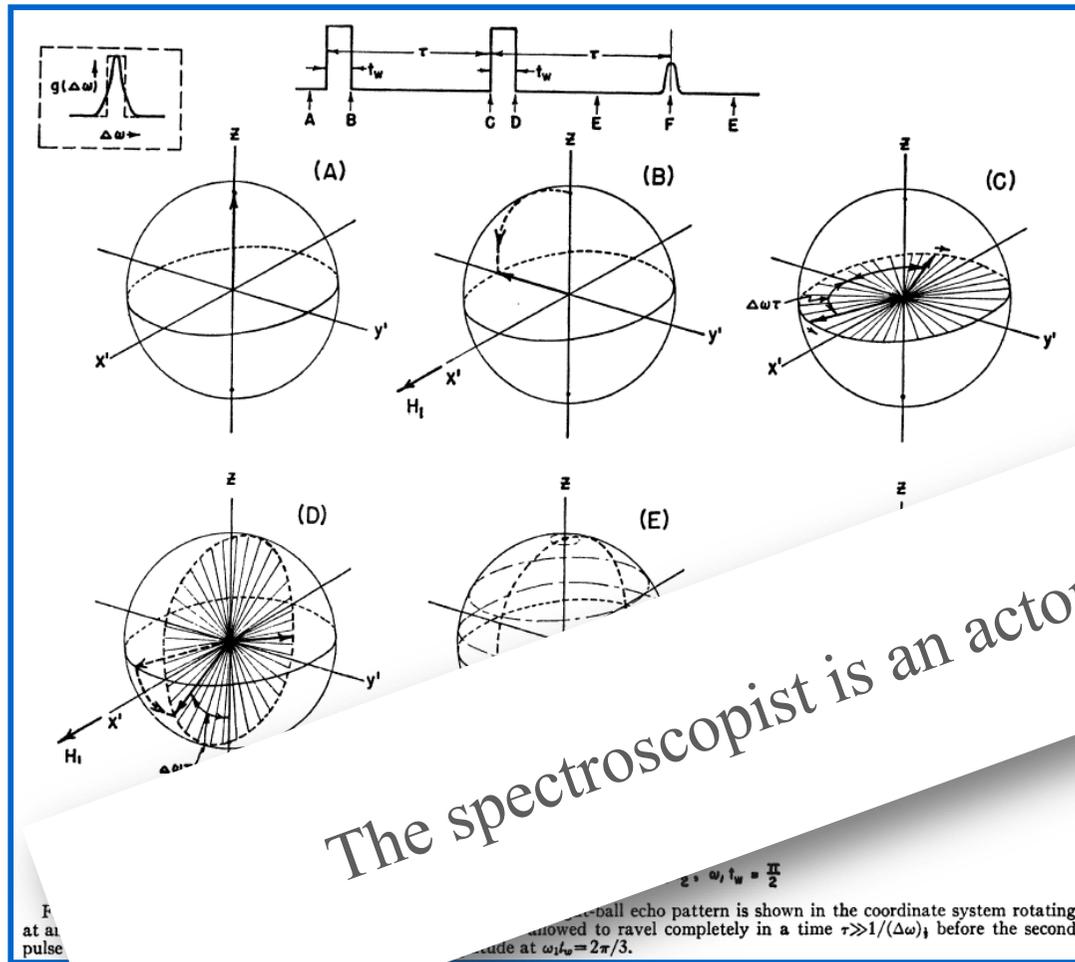


FIG. 2. The top trace indicates a beat note between an external r-f signal generator (near the Larmor frequency, loosely coupled to the inductive coil) and the nuclear signal shown alone (after detection) on the bottom trace. This beat note is identical in principle with the "wobble effect" (see reference 5) except that H_0 is held constant in this case.

E.L. Hahn, *Phys. Rev.*, **77**, 297 (1950)



Spin Echo – signal can be refocused

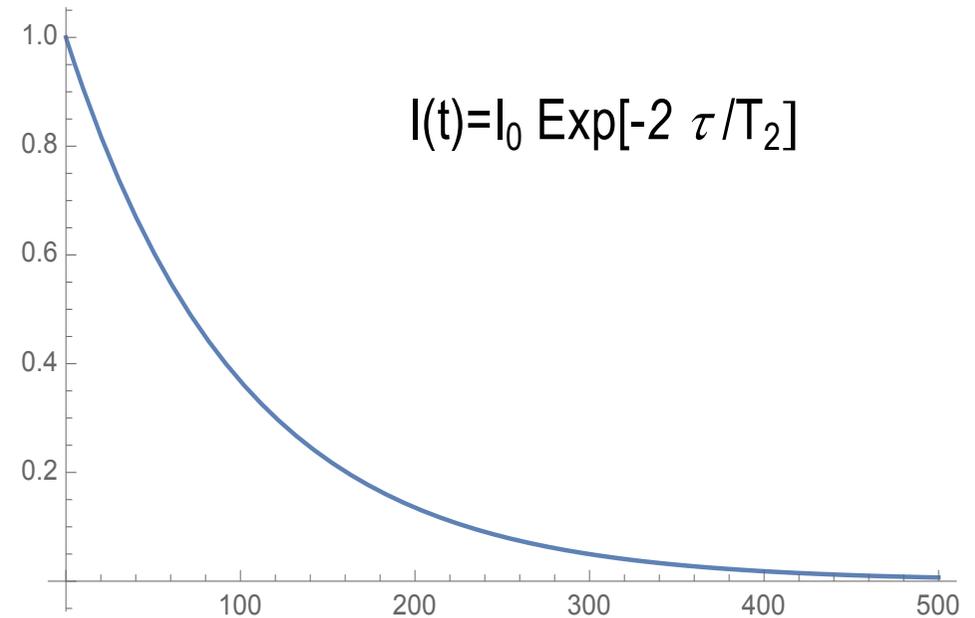
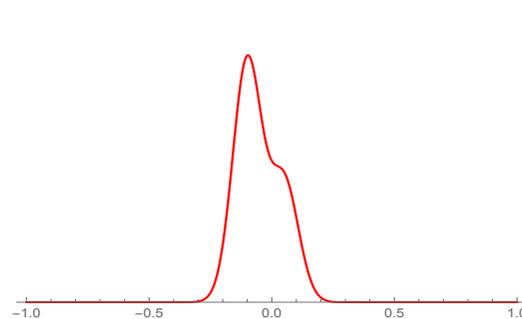
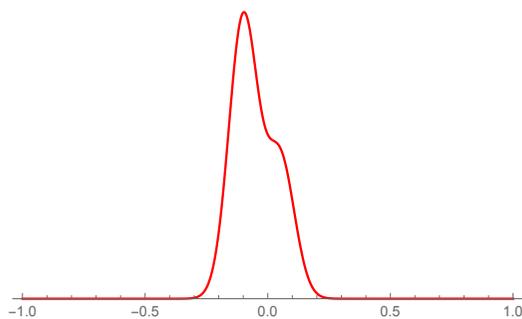
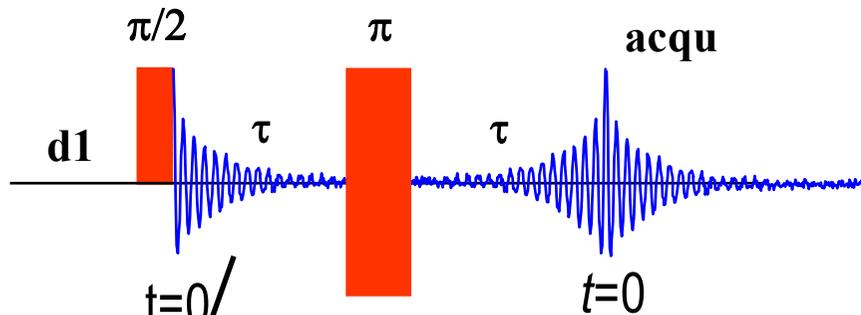


The spectroscopist is an actor of the experiment

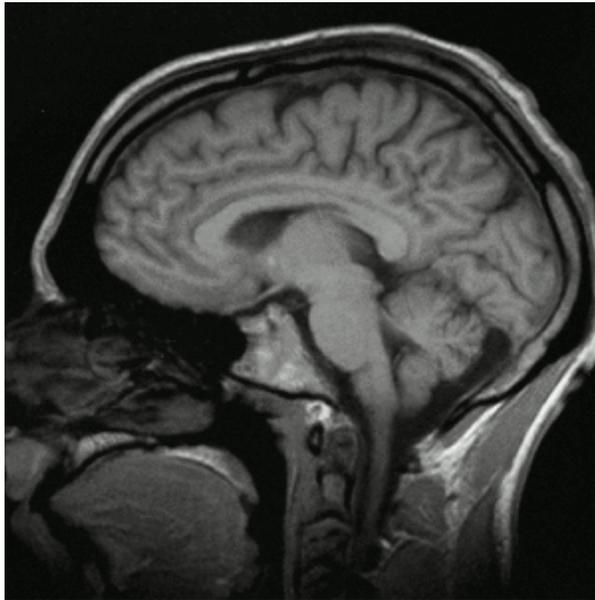


FIG. 5. Proton echo patterns in H_2O resulting from three applied r-f pulses. The pulses are visible in the upper two traces, and have a width $t_w \sim 0.5$ msec. In the upper trace $\tau = 0.008$ sec., $T = 0.067$ sec., and for the second trace $\tau = 0.046$ sec. and $T = 0.054$ sec. The bottom photograph shows a similar pattern for the case $T > 2\tau$ where induction decay signals can be seen following very short invisible r-f pulses. Saturation of a narrow band communications receiver, used in the case of the upper two traces, prevents the observation of these signals, whereas a wide band i.f. amplifier makes this observation possible in the bottom photograph.

E.L. Hahn, *Phys. Rev.*, **80**, 580 (1950)

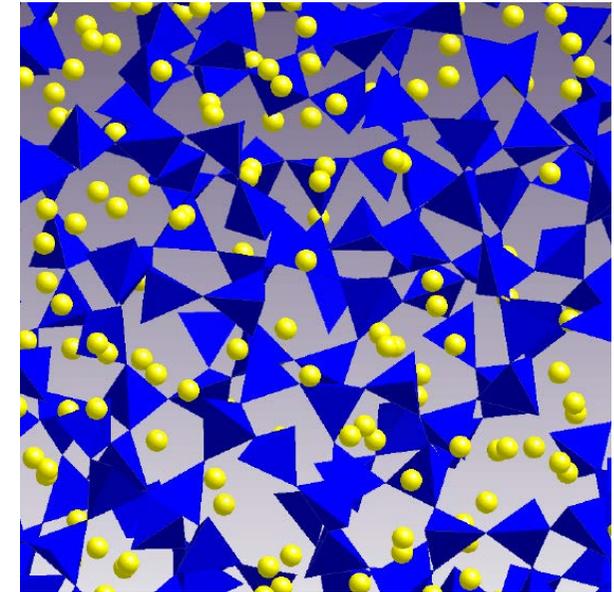
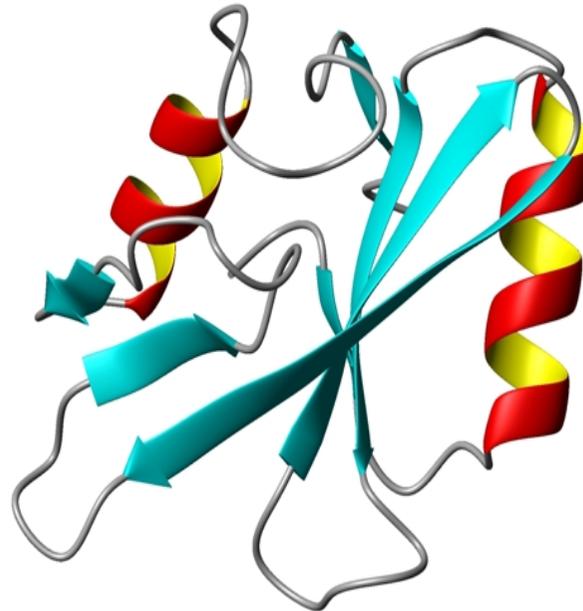


Measurement of T_2
 the life time of the coherence
 in the XY [transverse] plane



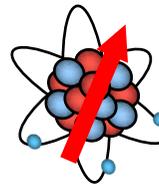
Spatial Encoding Field Gradients

Imaging
anatomy / function
Diffusion / Rheology



Local Field Homogeneous Principal Field (<ppm)

Spectroscopy
Chemistry



In a magnetic field - $\nu_0 = \gamma B_0 / 2\pi$

Image Formation by Induced Local Interactions: Examples Employing Nuclear Magnetic Resonance

AN image of an object may be defined as a graphical representation of the spatial distribution of one or more of its properties. Image formation usually requires that the object interact with a matter or radiation field characterized by a wavelength comparable to or smaller than the smallest features to be distinguished, so that the region of interaction may be restricted and a resolved image generated.

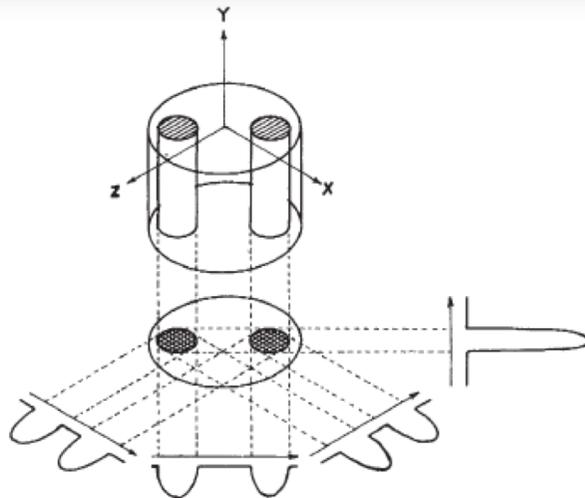
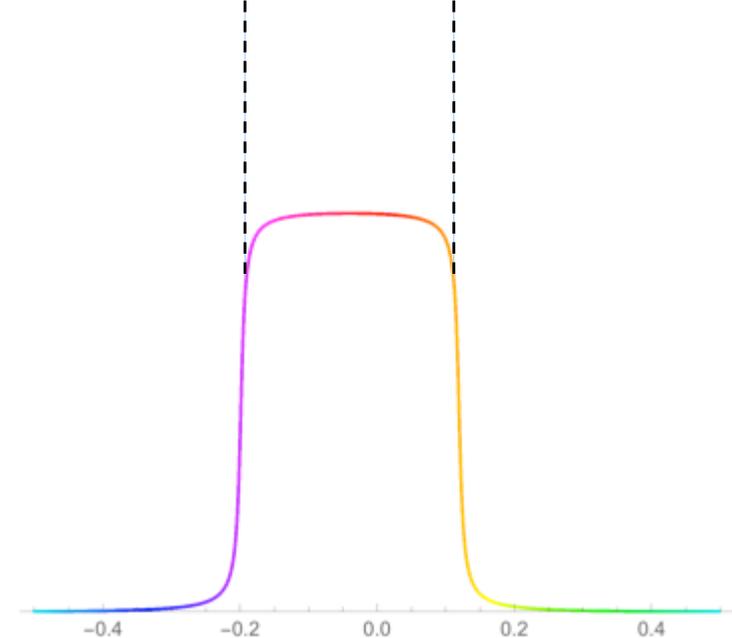
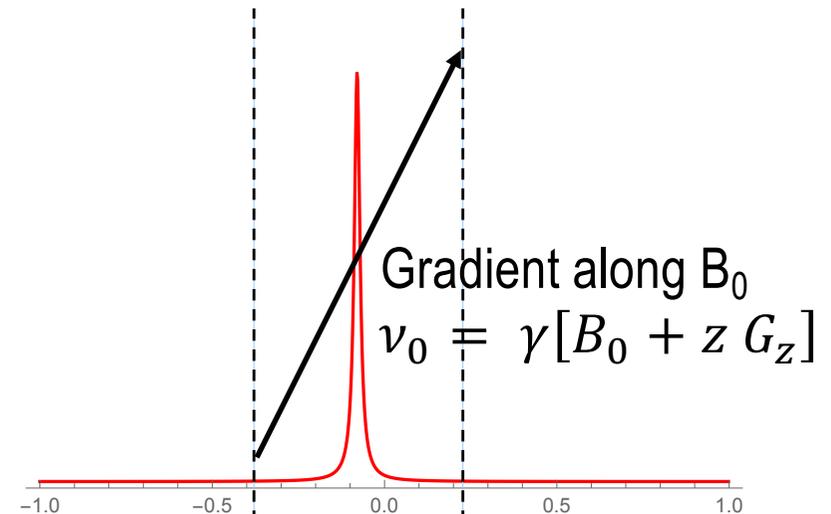


Fig. 1 Relationship between a three-dimensional object, its two-dimensional projection along the Y-axis, and four one-dimensional projections at 45° intervals in the XZ-plane. The arrows indicate the gradient directions.

Paul C. Lauterbur – Nature 1973
Nobel Physiology & Medicine 2003



Projection on the Field Gradient

The Discovery of Chemical Shift

The Dependence of a Nuclear Magnetic Resonance Frequency upon Chemical Compound*

W. G. PROCTOR AND F. C. YU
Department of Physics, Stanford University, Stanford, California
January 18, 1950

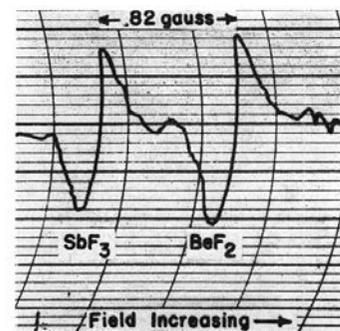
IN the course of measurements on N^{14} , mentioned in the previous letter, we made the surprising observation that its frequency of resonance, in liquid samples, depended strongly upon the chemical compound in which it was contained.^{1,2} This effect is strikingly demonstrated by the appearance of two resonances, separated by 1.6 kc in the neighborhood of 3300 kc, corresponding to a field of 10,500 gauss, using a solution of NH_4NO_3 in 2.0-molar $MnSO_4$ as a sample. These resonances presumably arise from the NH_4^+ and NO_3^- complexes, since samples of $NH_4C_2H_3O_2$ and HNO_3 separately give rise to two different resonances whose frequencies approximate those from the above sample. The separation is four times greater than the line widths measured between points of maximum slope.

Phys. Rev., 77, 716 (1950)

Dependence of the F^{19} Nuclear Resonance Position on Chemical Compound*

W. C. DICKINSON
Research Laboratory of Electronics, Massachusetts Institute of Technology, Cambridge, Massachusetts
January 9, 1950

MOST unexpectedly, it has been found that for F^{19} the value of the applied magnetic field H_0 for nuclear magnetic resonance at a fixed radiofrequency depends on the chemical compound containing the fluorine nucleus. The assumption has generally been made that the time average of all internal magnetic fields is zero, excluding of course the small diamagnetic field at the nucleus due to the Larmor precession of its atomic electrons in H_0 . Nuclear resonance shifts in metals,¹ interpreted as being due to the conduction electrons, are larger by about an order of magnitude than those reported here.²

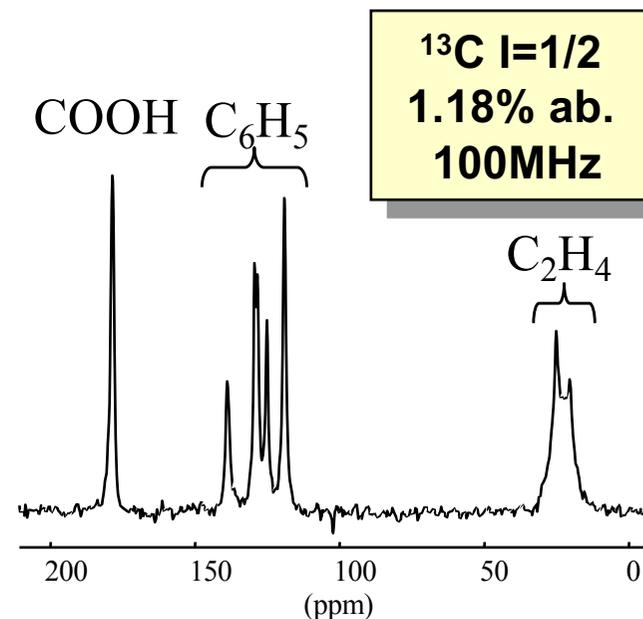
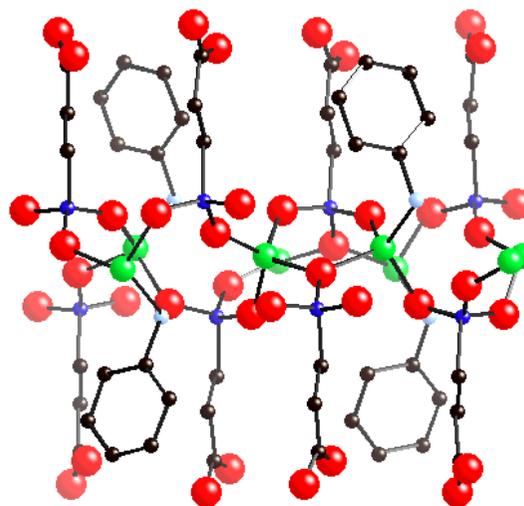
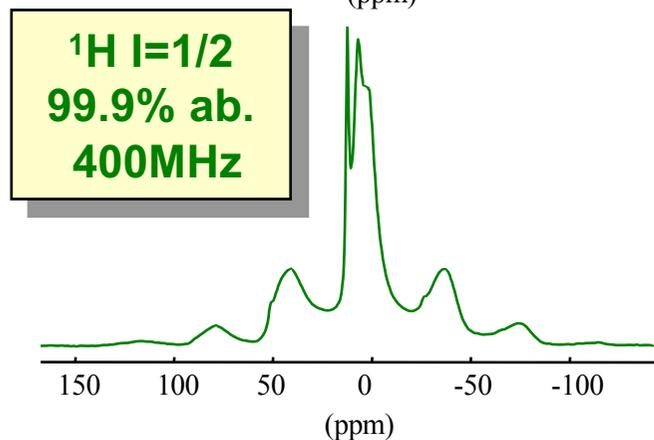
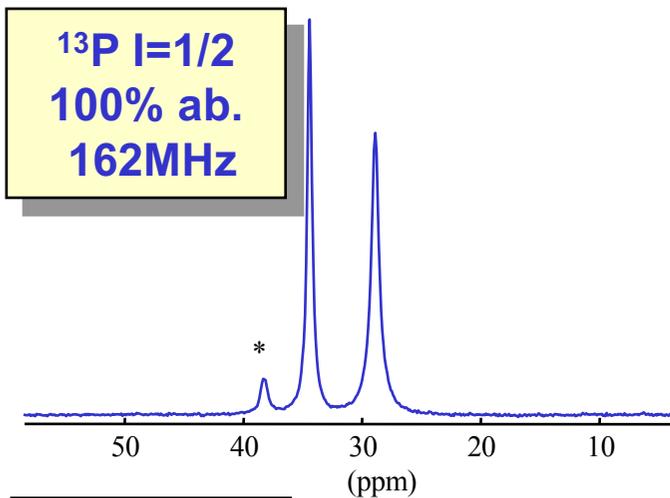


Phys. Rev., 77, 736 (1950)

FIG. 2. The nuclear resonances of F^{19} in a single sample containing a half and half mixture of SbF_3 and BeF_3 (saturated aqueous solutions). The applied resonance magnetic field is about 7000 gauss at a radiofrequency of 28.0 megacycles.



Phosphonate (B.Bujoli – Nantes)

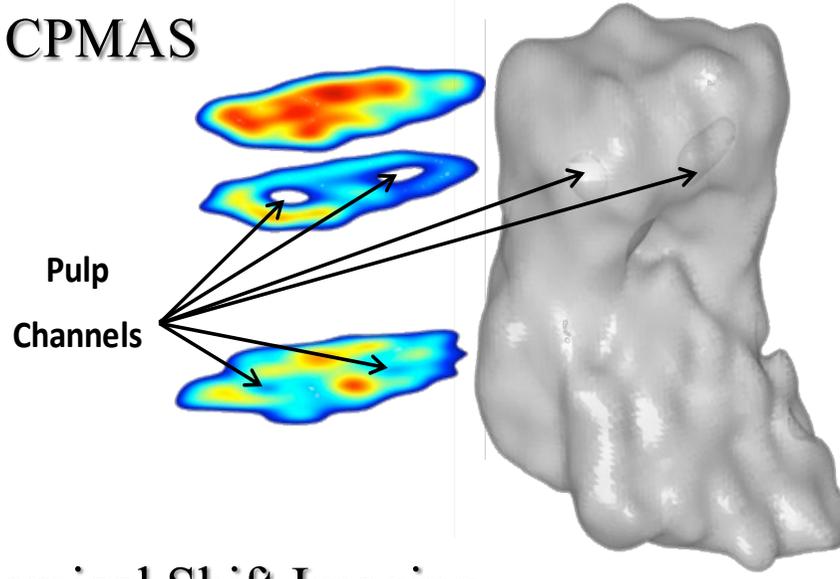


**^{17}O I=5/2
0.037% ab.
54 MHz**

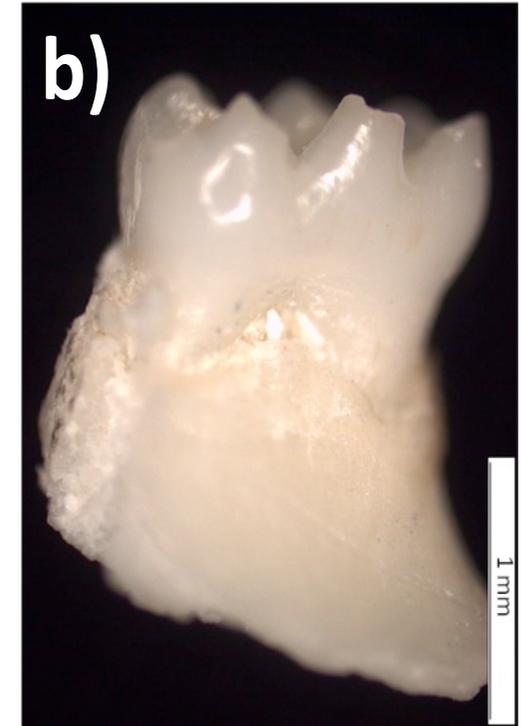
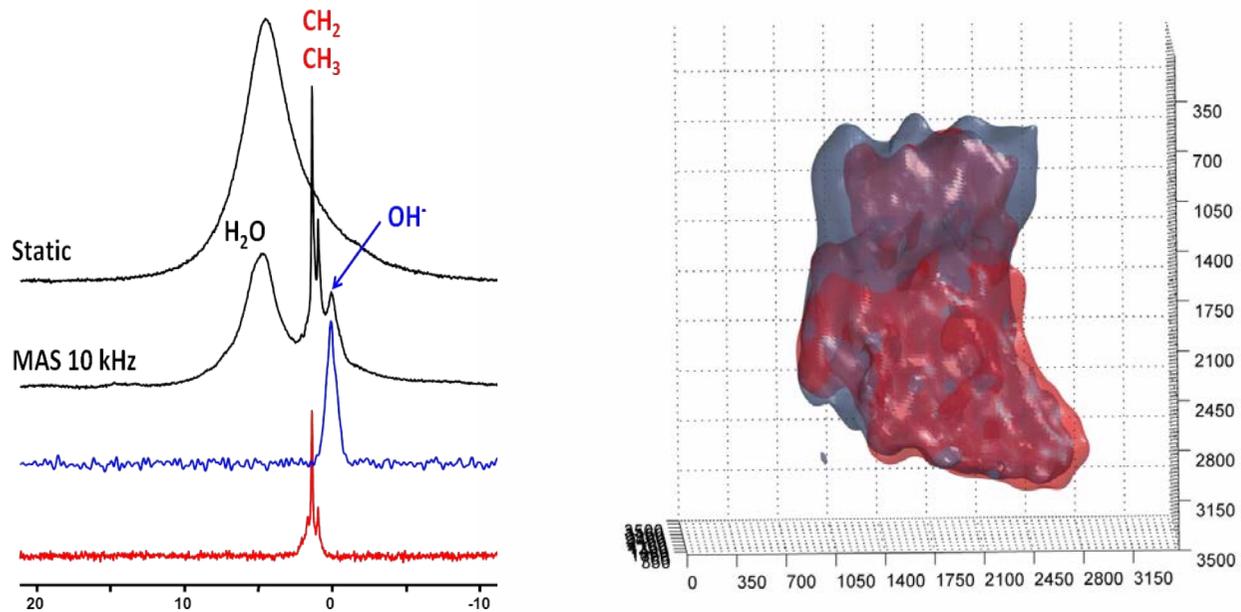
**^{14}N I=1
99.6 % ab.
29 MHz**

**^{67}Zn I=5/2
4.11% ab.
25 MHz**

$^1\text{H}/^{31}\text{P}$ CPMAS

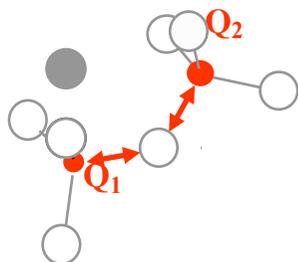
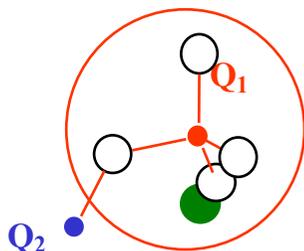


^1H Chemical Shift Imaging



Chemical Shift Anisotropy
 electronic shielding
 first shells, coordination and
 geometry

~10s of kHz

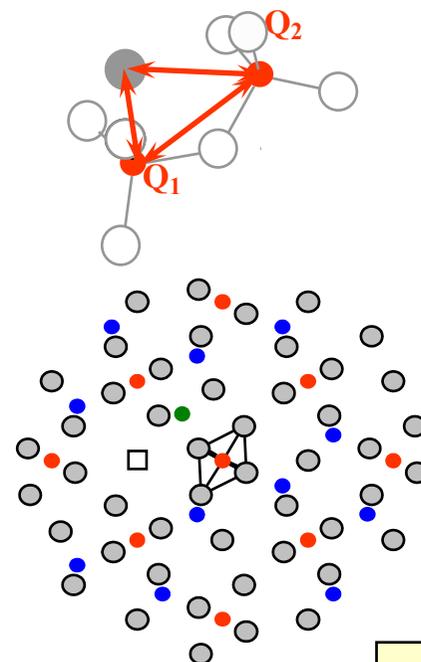


< 100s of Hz

Indirect J Coupling
 chemical bonding
 Connectivity

Dipolar interaction
 neighboring spins
 Distances

~kHz $1/r^3$

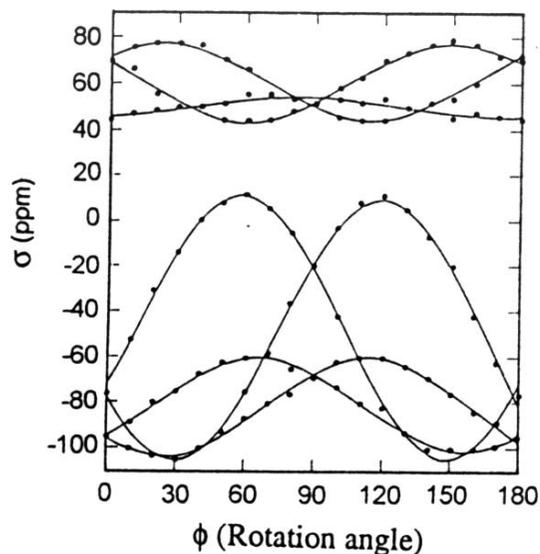


I > 1/2 up to MHz

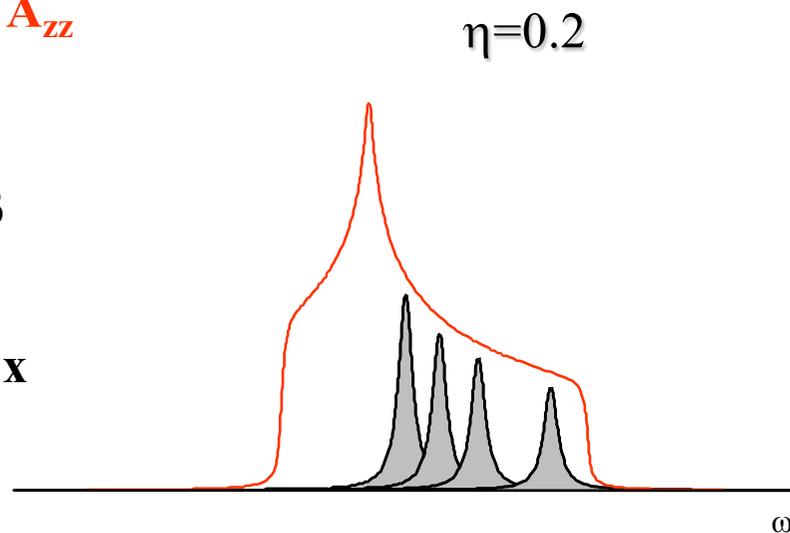
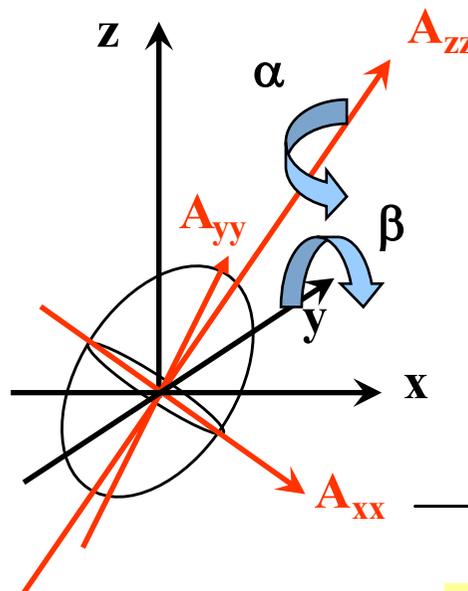
Quadrupolar interaction
 Electric Field Gradient
 Surrounding Charges,
 electrons and nuclei
 geometry

Interactions are anisotropic and take the following form (at 1st order):

$$\nu = \nu_0 + A \left(\frac{3\cos^2\beta - 1}{2} - \frac{\eta}{2} \sin^2\beta \cos 2\alpha \right)$$

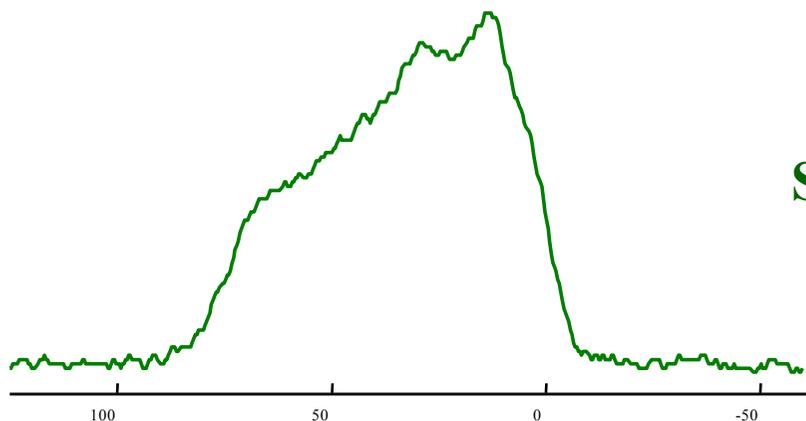


Single Crystal rotation pattern

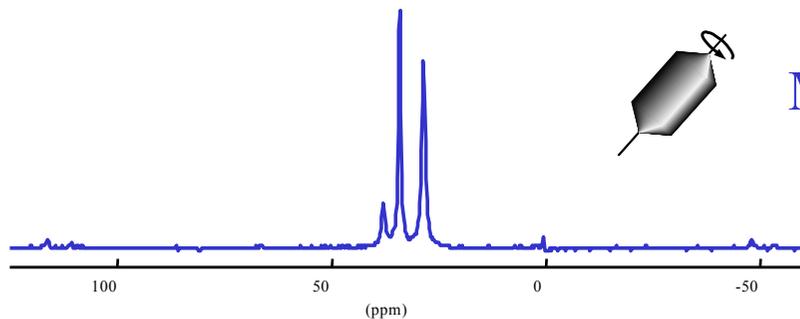
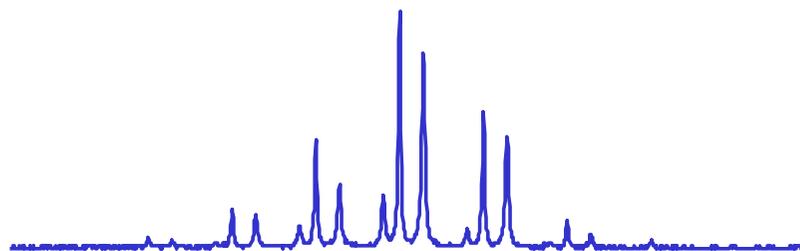


Crystalline Powder overlapping Broad lines

^{31}P Spin 1/2 CSA

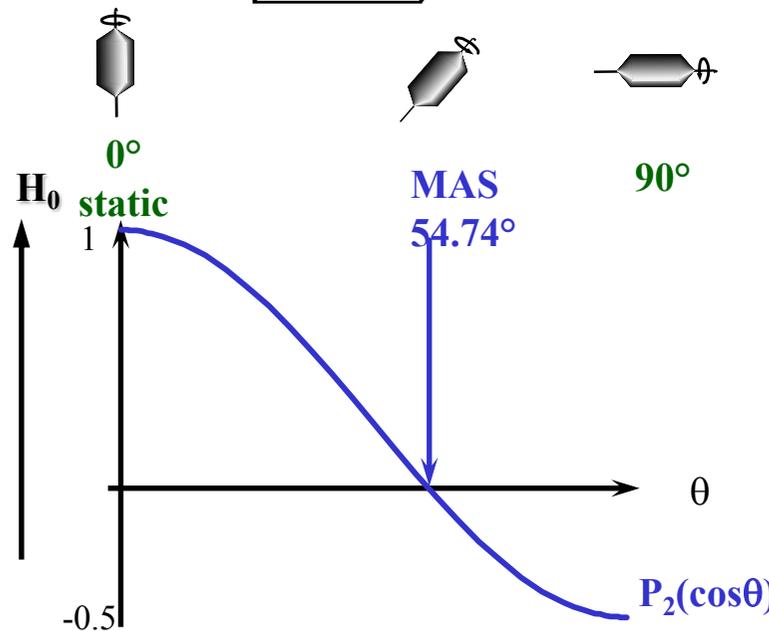
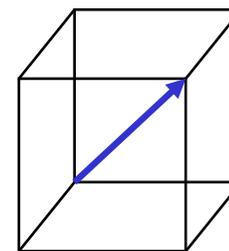


Static



MAS

Modulation into sharp lines



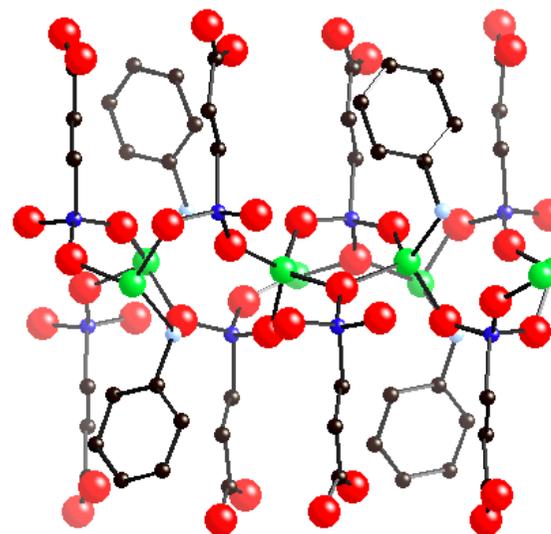
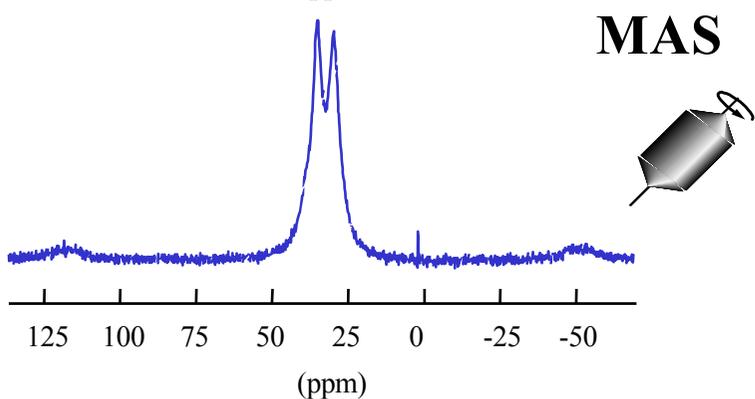
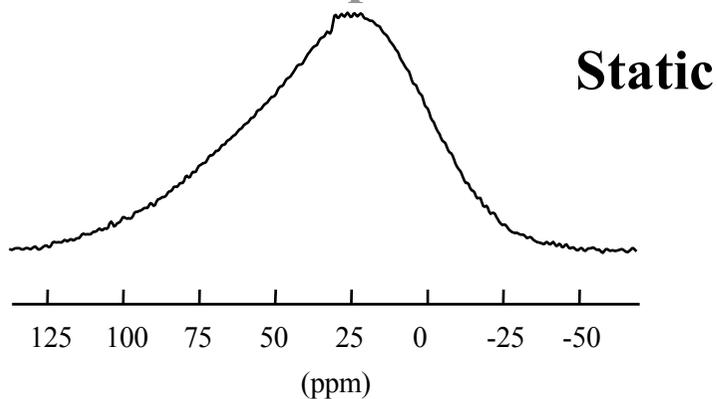
Dipolar $\rightarrow 0$
 Chem. Shift $\rightarrow \delta_{\text{iso}}$
 Jcoupling

Quad 1st $\rightarrow 0$
 Quad 2nd $\rightarrow \delta^{2\text{nd}}$

^{31}P

high speed MAS spinning attenuate dipolar coupling and modulates chemical shift anisotropy

unresolved spectrum
CSA + Dipolar



~~CSA + Dipolar~~

^{31}P

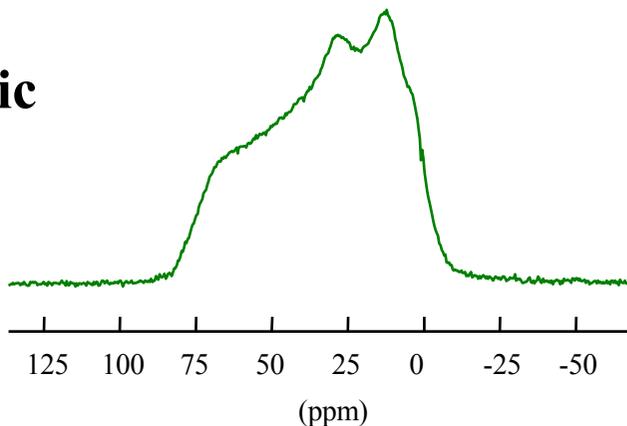
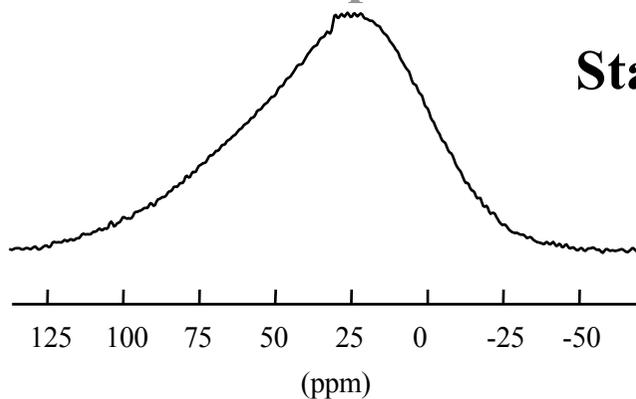
Proton decoupling averages strong ^{31}P - ^1H dipolar interaction



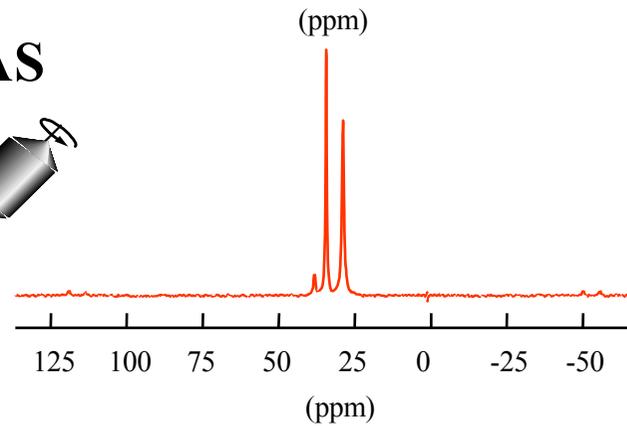
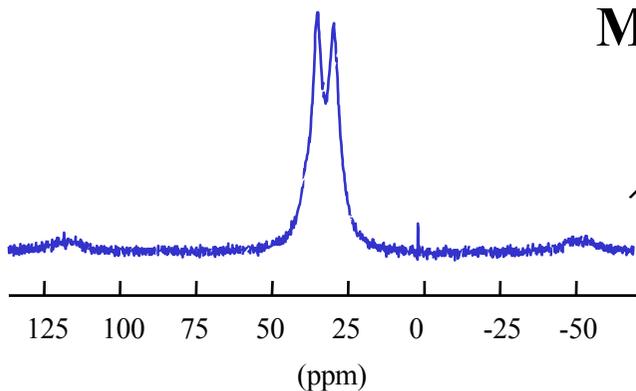
unresolved spectrum
CSA + Dipolar

~~CSA + Dipolar~~

Static



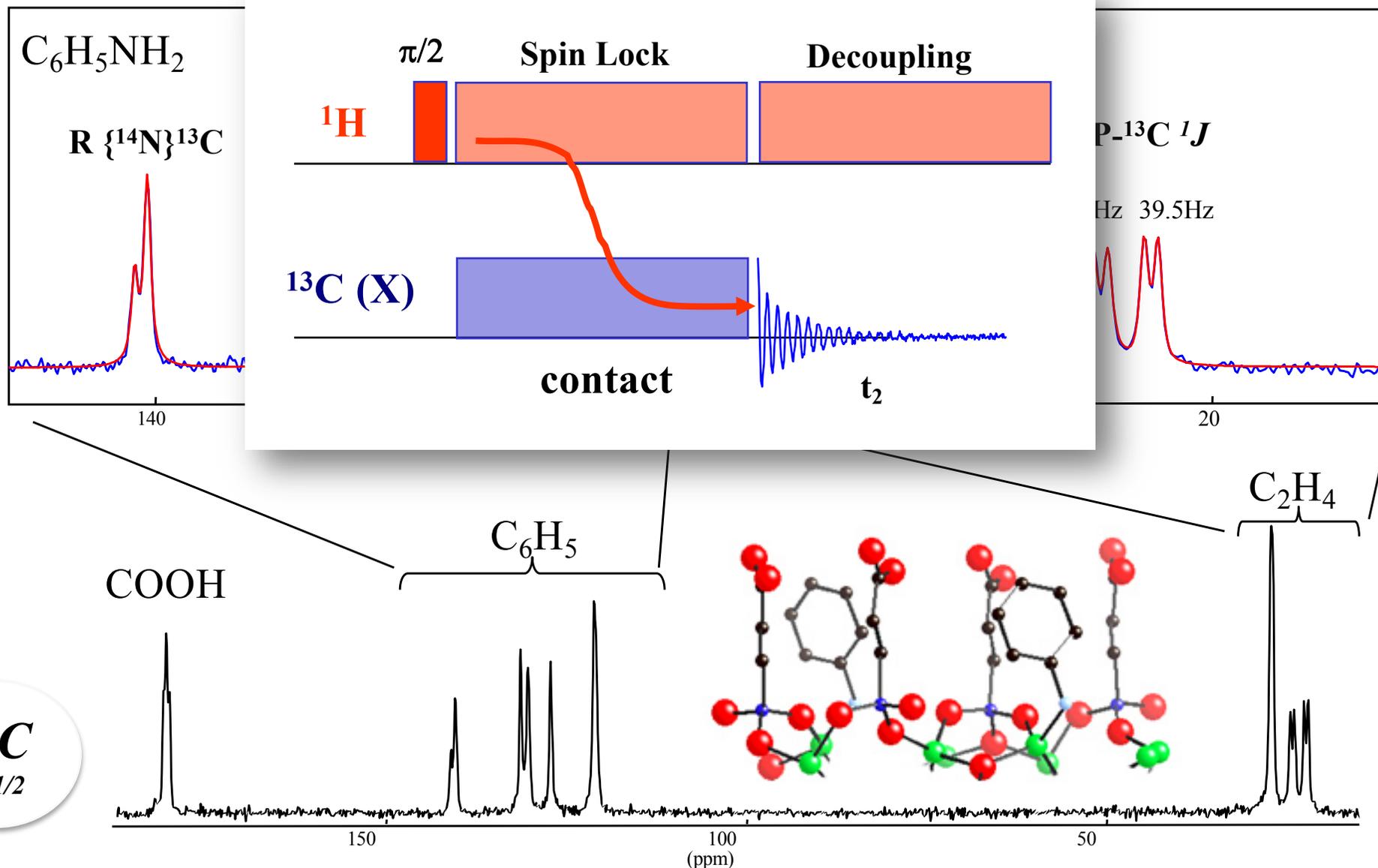
MAS



~~CSA + Dipolar~~

simplified resolved spectrum

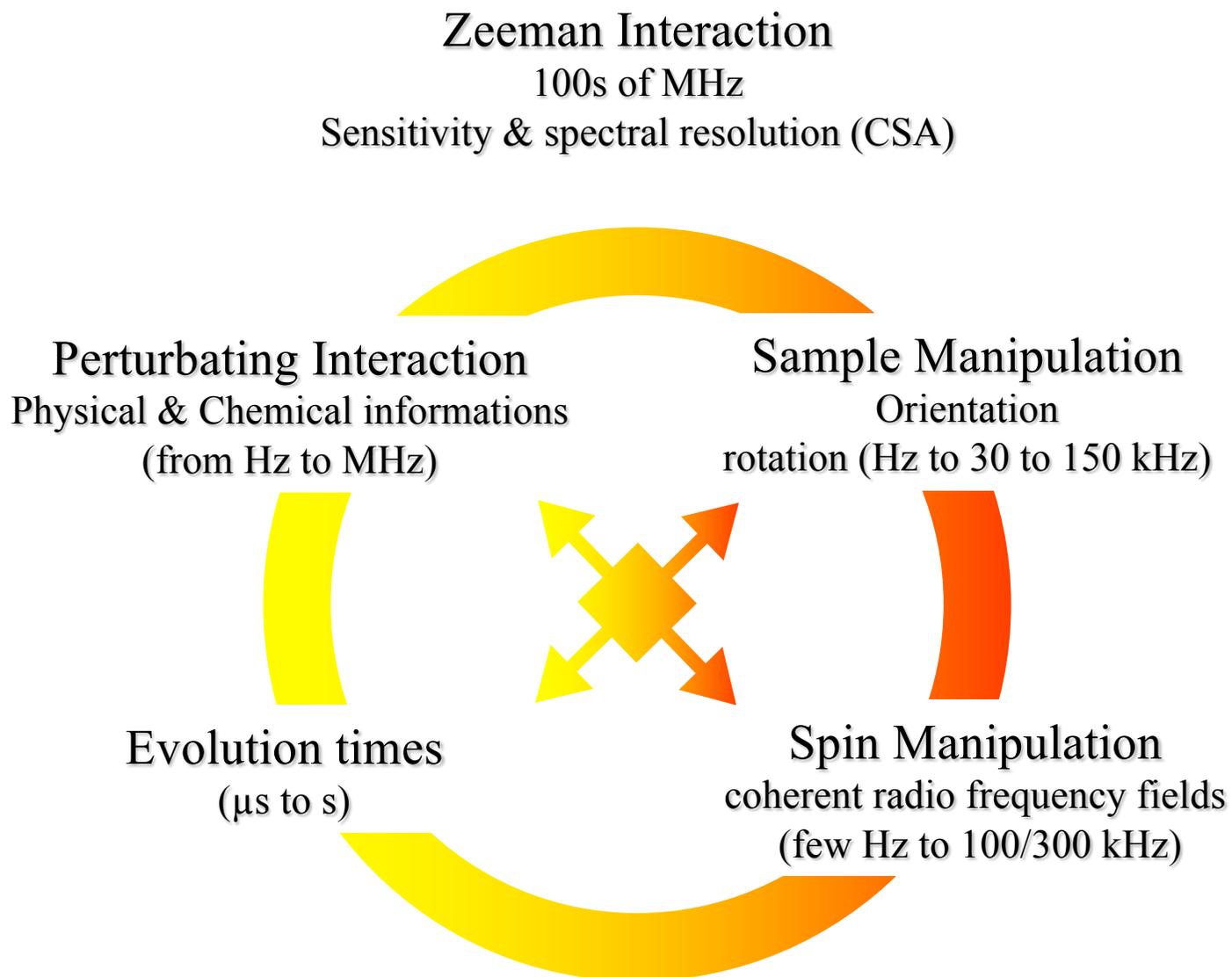
high speed MAS spinning attenuate dipolar coupling and modulates chemical shift anisotropy



D.Massiot, F.Fayon, M.Deschamps, S.Cadars, P.Florian, V.Montouillout, N.Pellerin, J.Hiet, A.Rakhmatullin, C.Bessada

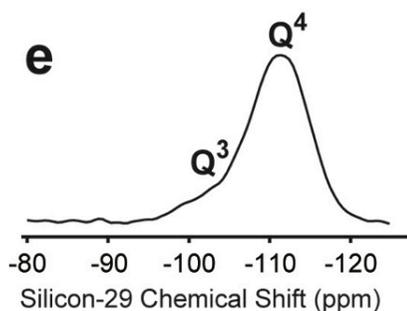
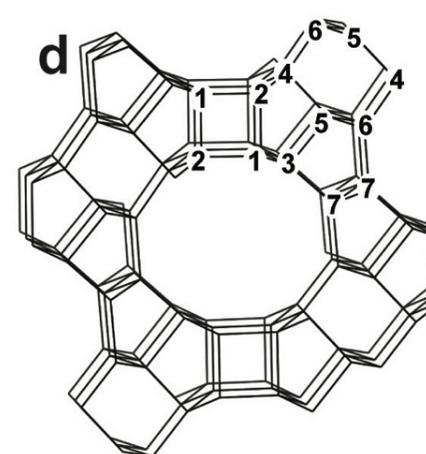
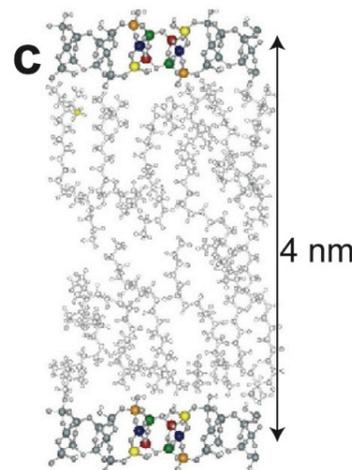
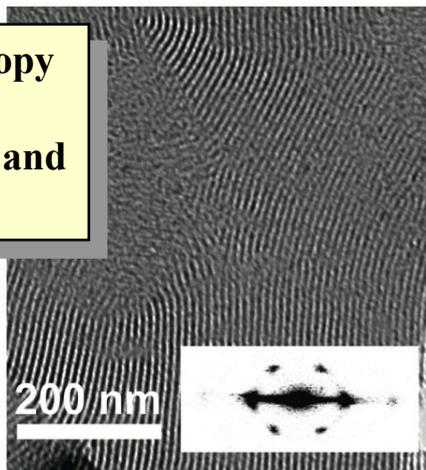
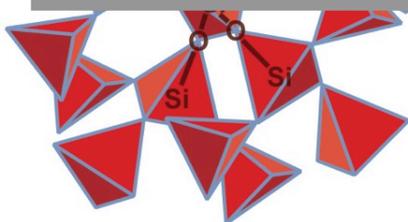
'Detection and use of small J couplings in solid state NMR experiments.'

[Comptes Rendus de Chimie](#) 13 117-129 2010

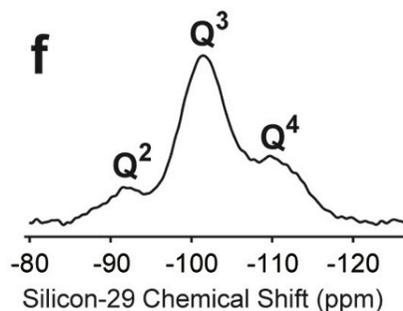


Long lasting relaxation times T1 up to 100s sec – T2 up to 100s ms

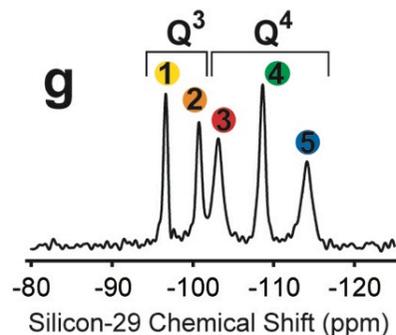
**Chemical Shift Anisotropy
electronic shielding
first shells, coordinence and
geometry**



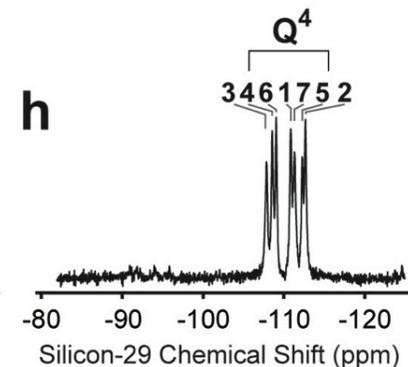
Glass



Mesoporous Silica

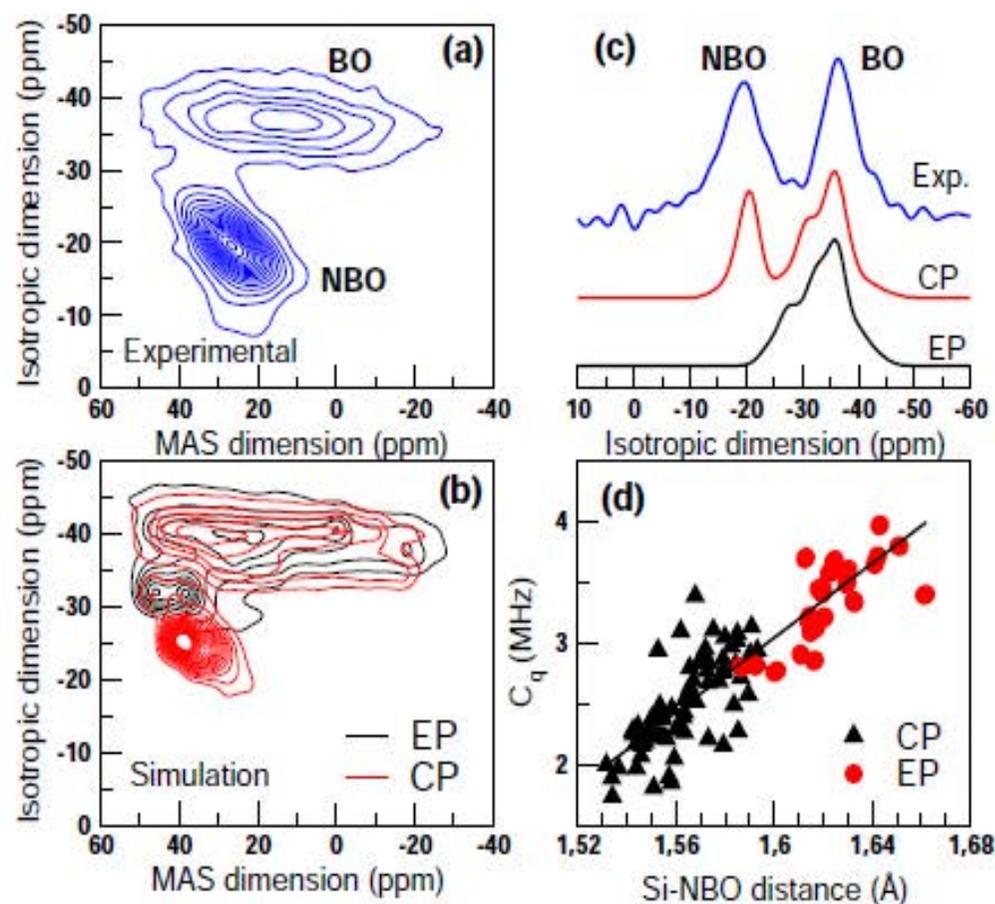
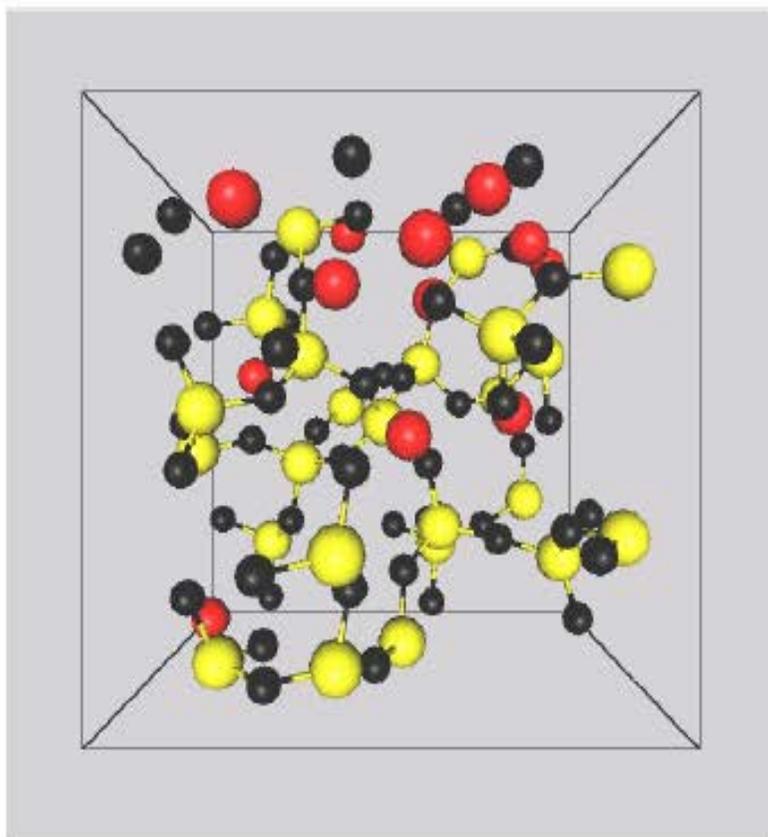


SiO_2 -surfactant
Mesophase

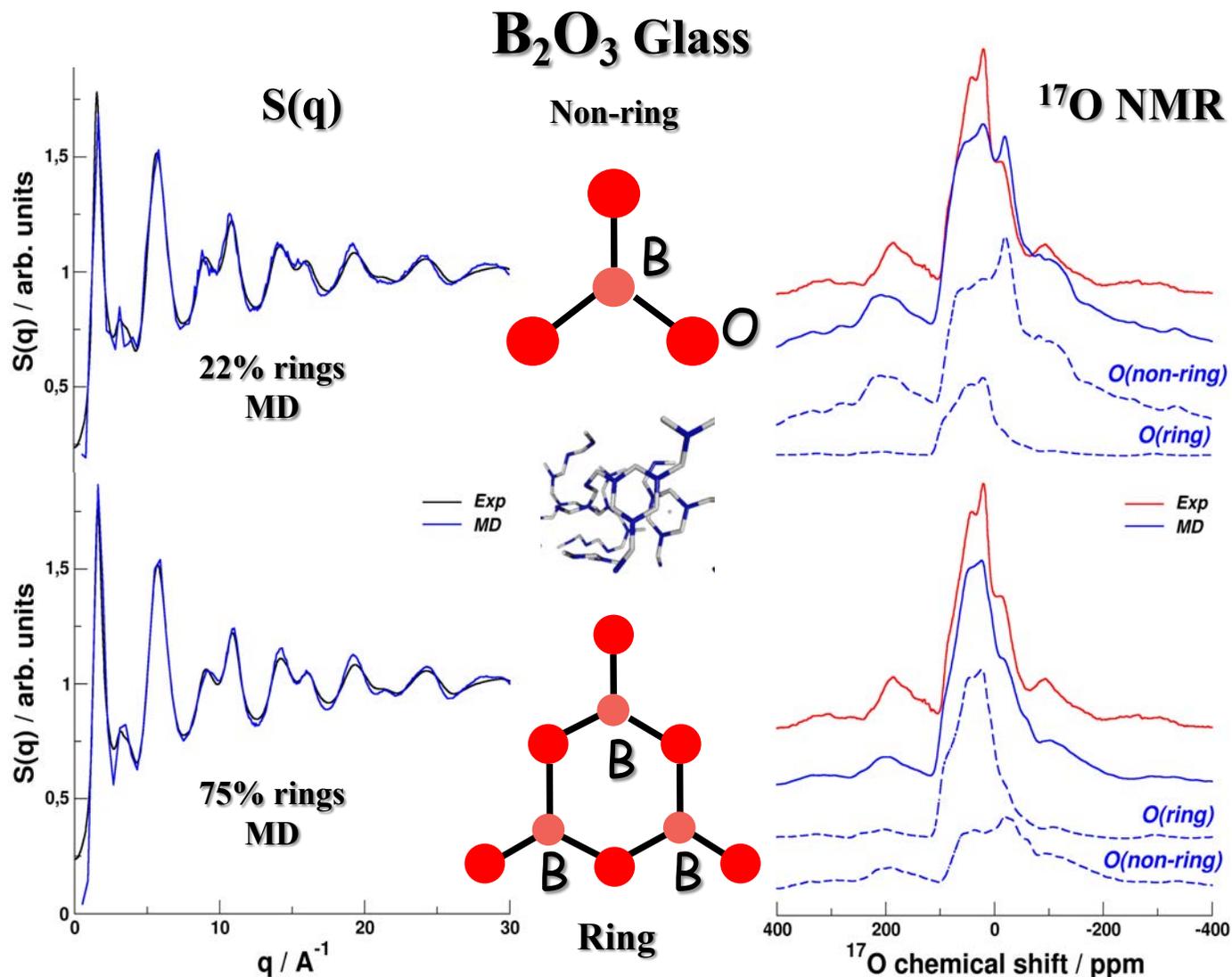


Zeolite

*Average Chemical Shift ~ local field
electronic shielding $\propto B_0 \rightarrow$ first shells, coordinence and geometry*

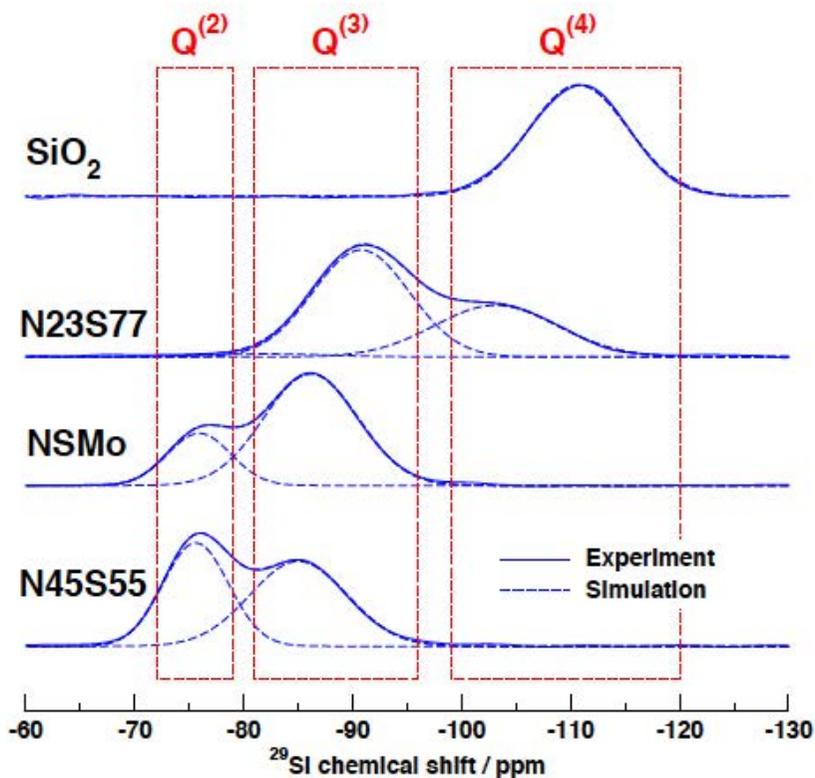
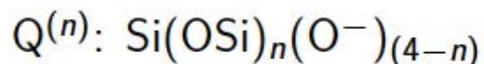


The combination of molecular dynamics simulations with first-principles (or *ab initio*) calculations with density functional theory (DFT) enable the prediction of NMR spectra from a structural model (here oxygen-17 MQMAS). Two spectra are compared, the first from classical molecular dynamics (EP, Effective Potential) and the second from *ab-initio* molecular dynamics (CP, Car-Parinello). Differences between the two models in Si-NBO distances induce a difference in the position of isotropic peaks of non-bridging oxygens. This can be explained by the high sensitivity of the quadrupole interaction (C_Q) to the Si-NBO distance

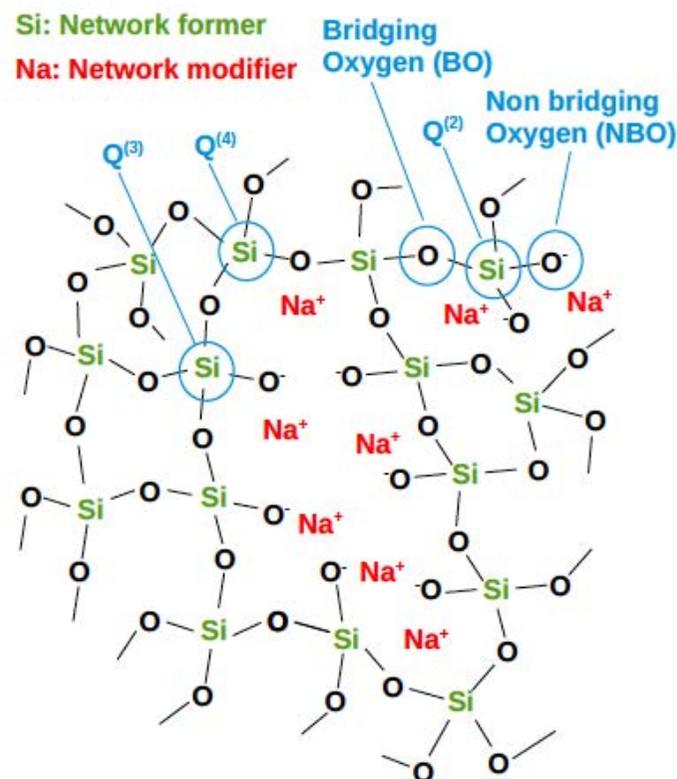


G.Ferlat, T.Charpentier, A.P.Seitsonen, A.Takada, M.Lazzeri, L.Cormier, G.Calas, F.Mauri
 "Boroxol Rings in Liquid and Vitreous B₂O₃ from First Principles"
 Phys. Rev. Lett. 101 065504 2008

^{29}Si MAS NMR



^{29}Si MAS NMR:
Direct access to silicon $Q^{(n)}$ speciation

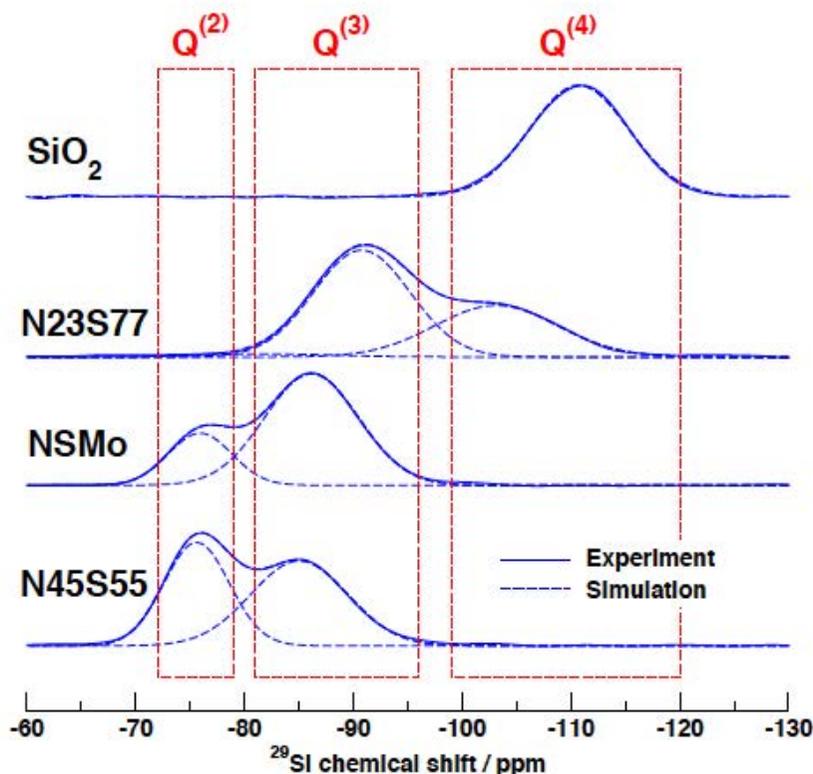
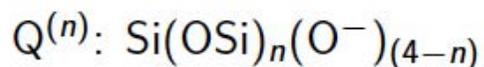


NMR peaks reflective of a Gaussian distribution of δ_{iso} ($I=1/2$)

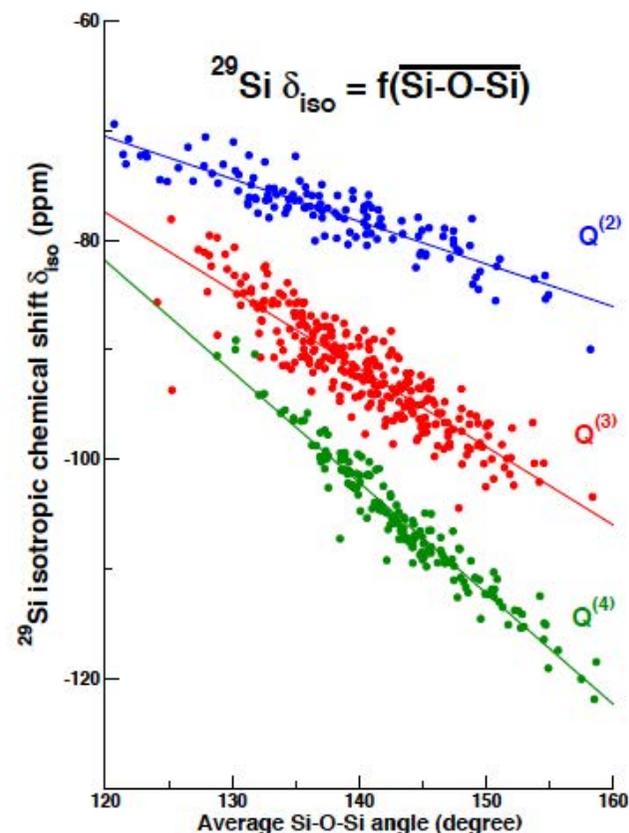
Note: $\delta_{iso} = -(\sigma_{ref} - \sigma_{iso})$

Angeli F., *et al.* Geochim. Cosmochim. Acta, 75, 2453-2469, 2011 - Ispas S., *et al.* Solid State Sci., 12, 183-192, 2010

^{29}Si MAS NMR: Direct access to silicon $Q^{(n)}$ speciation

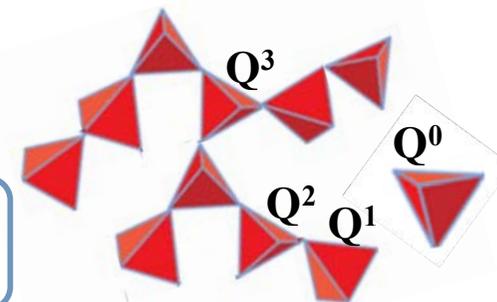


Binary $\text{Na}_2\text{O} - \text{SiO}_2$ glasses
 NSMo: N39S60 + 1 MoO_3

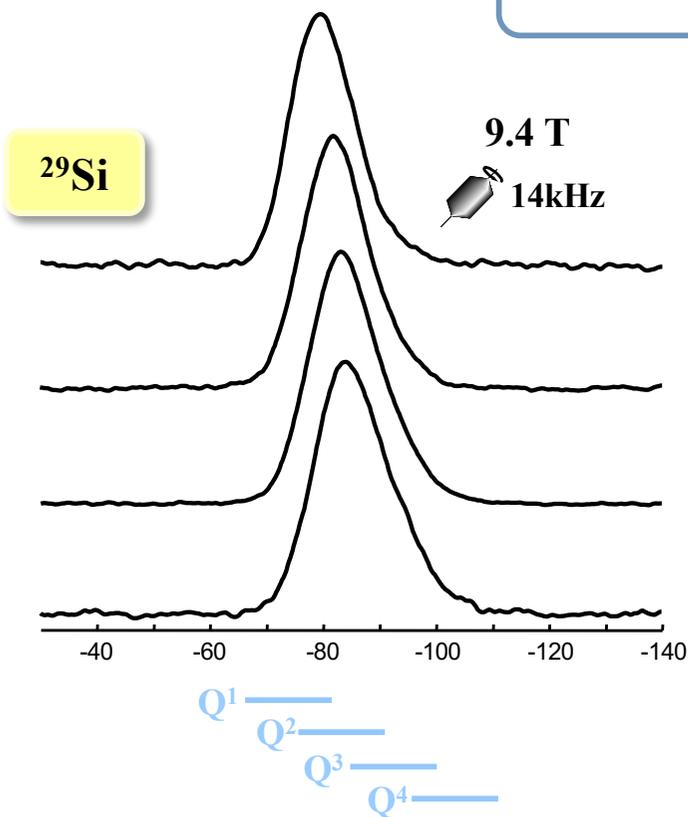


The NMR response of $Q^{(n)}$ species to disorder (bond angle distribution) is different.

Silicate (phosphate) network: Qⁿ units
(SiO₄, PO₄ tetrahedra with n bridging oxygen atoms)



1D MAS spectra of CaO-SiO₂-P₂O₅ glasses

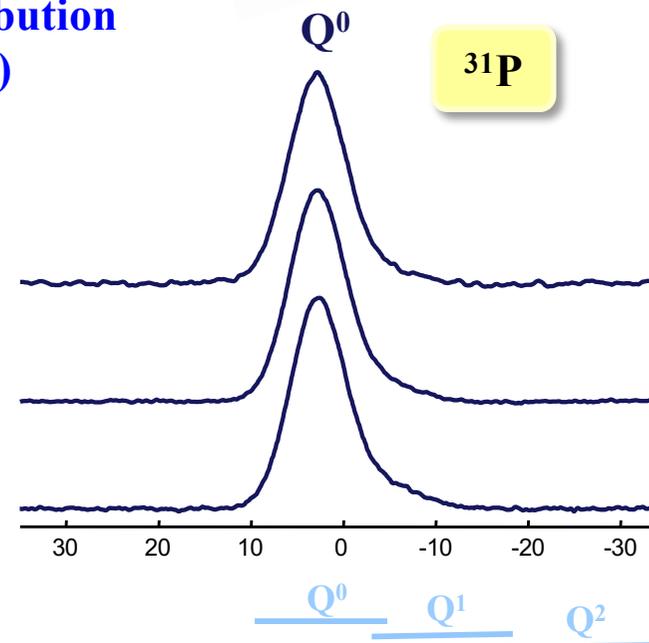


Broad chemical shift distribution
(disordered materials)

Ca/Si = 1.11

↓ P₂O₅

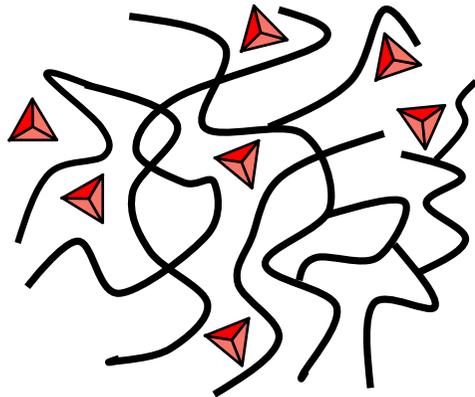
2.6 P₂O₅
3.8 P₂O₅
5.0 P₂O₅



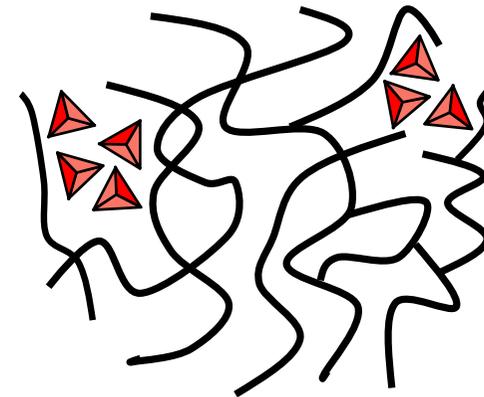
Lack of resolution: Qⁿ units
quantification ?

Mainly orthophosphate units
(Q⁰ : PO₄³⁻)

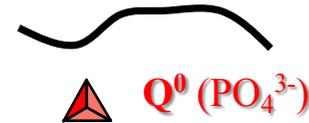
Chemical homogeneity ? (Random distribution)



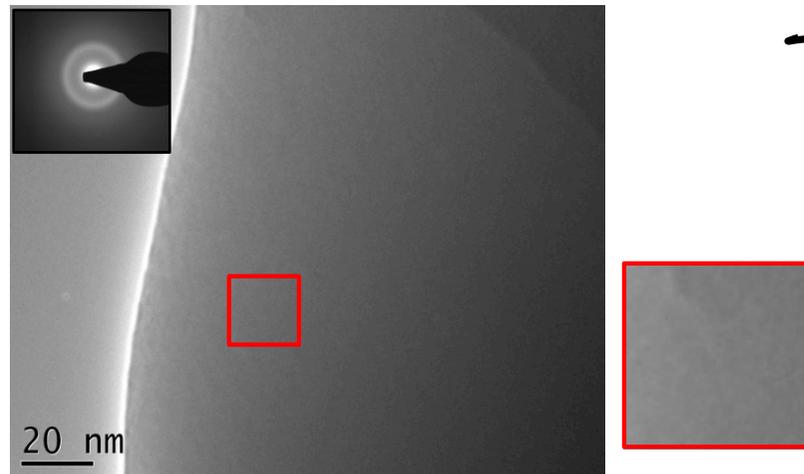
Phosphate clusters ?



Silicate chains



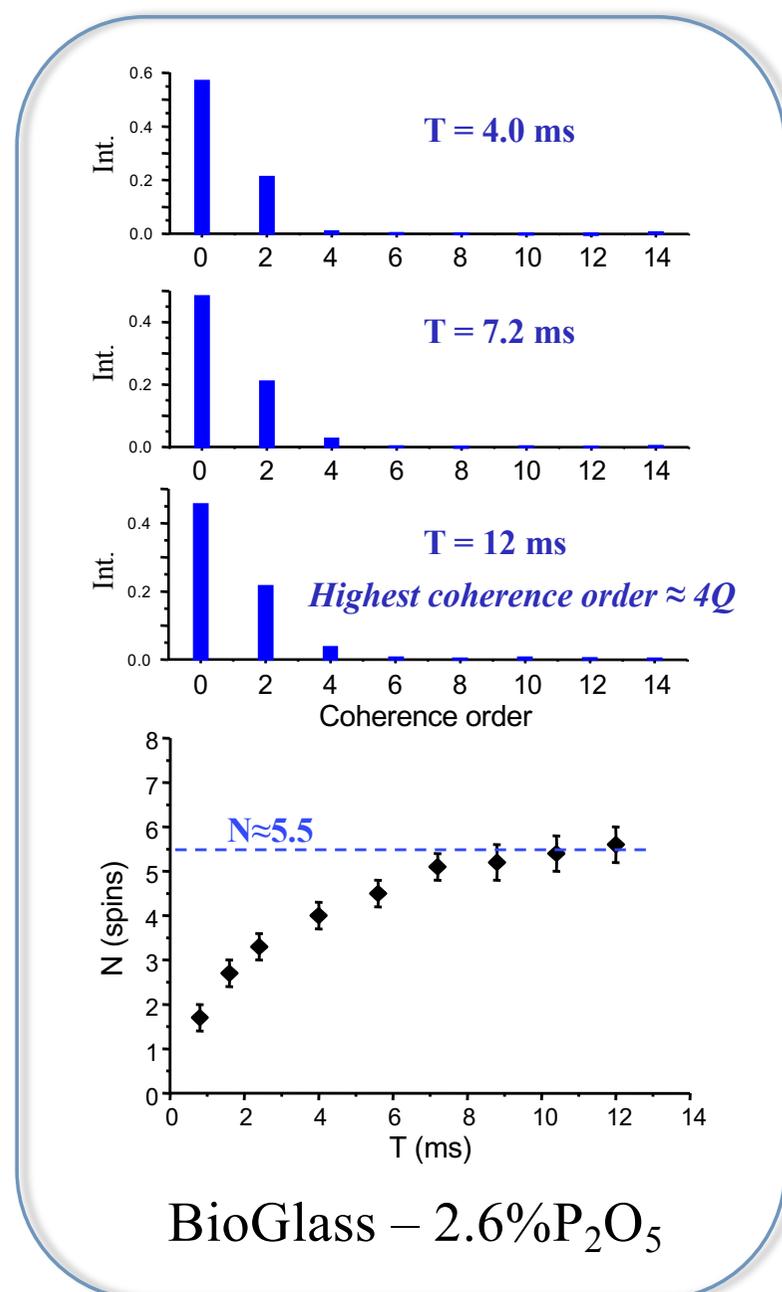
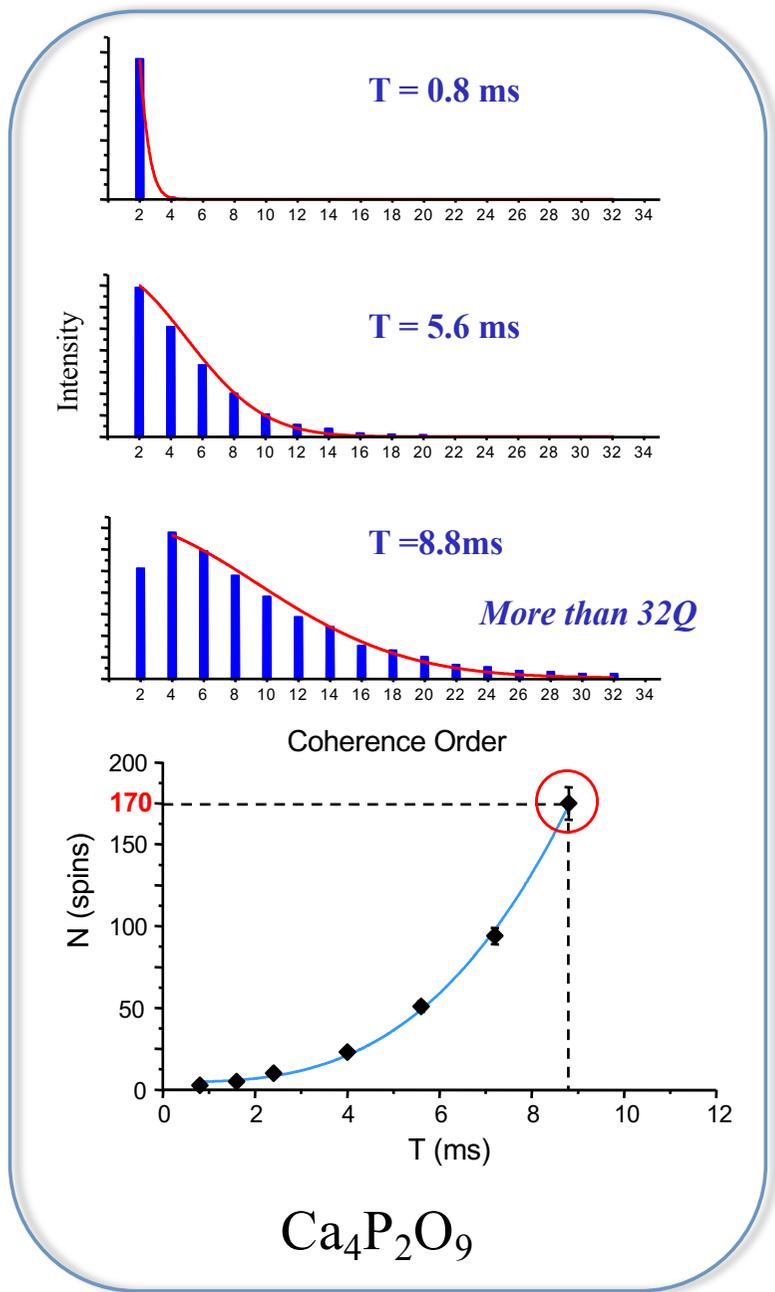
$\text{CaO-SiO}_2\text{-P}_2\text{O}_5$

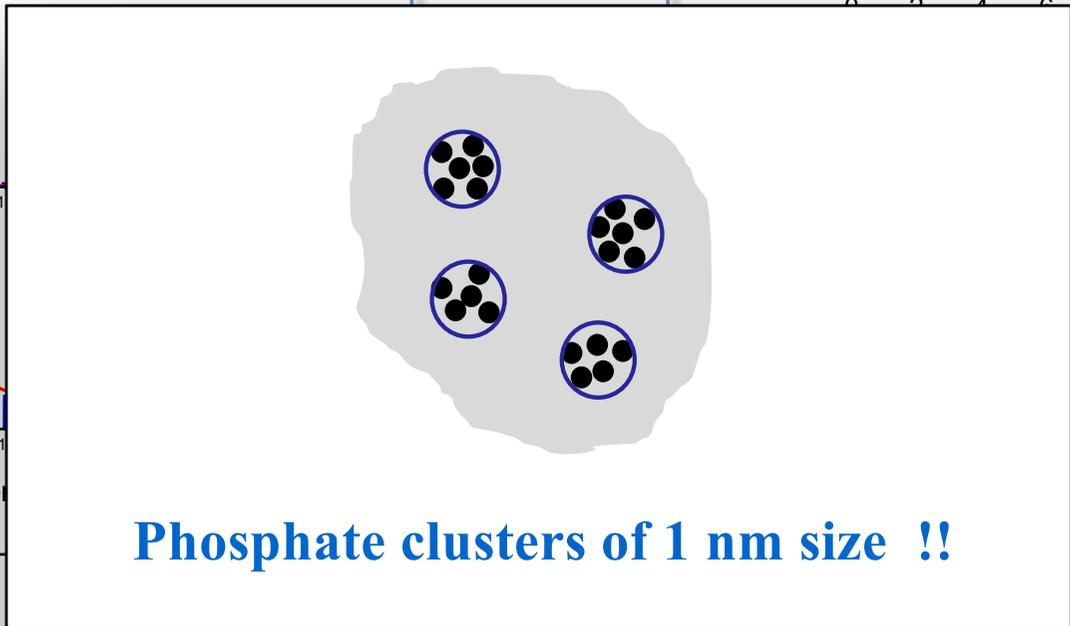
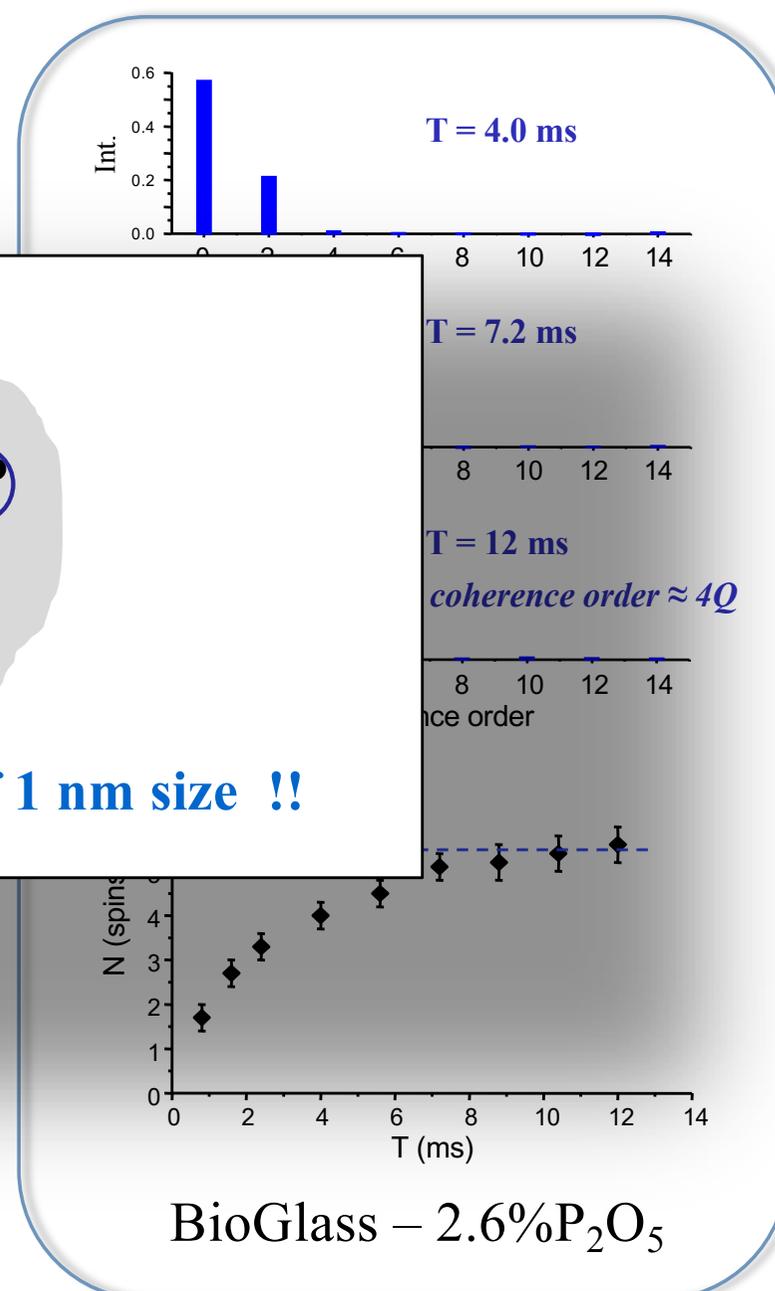
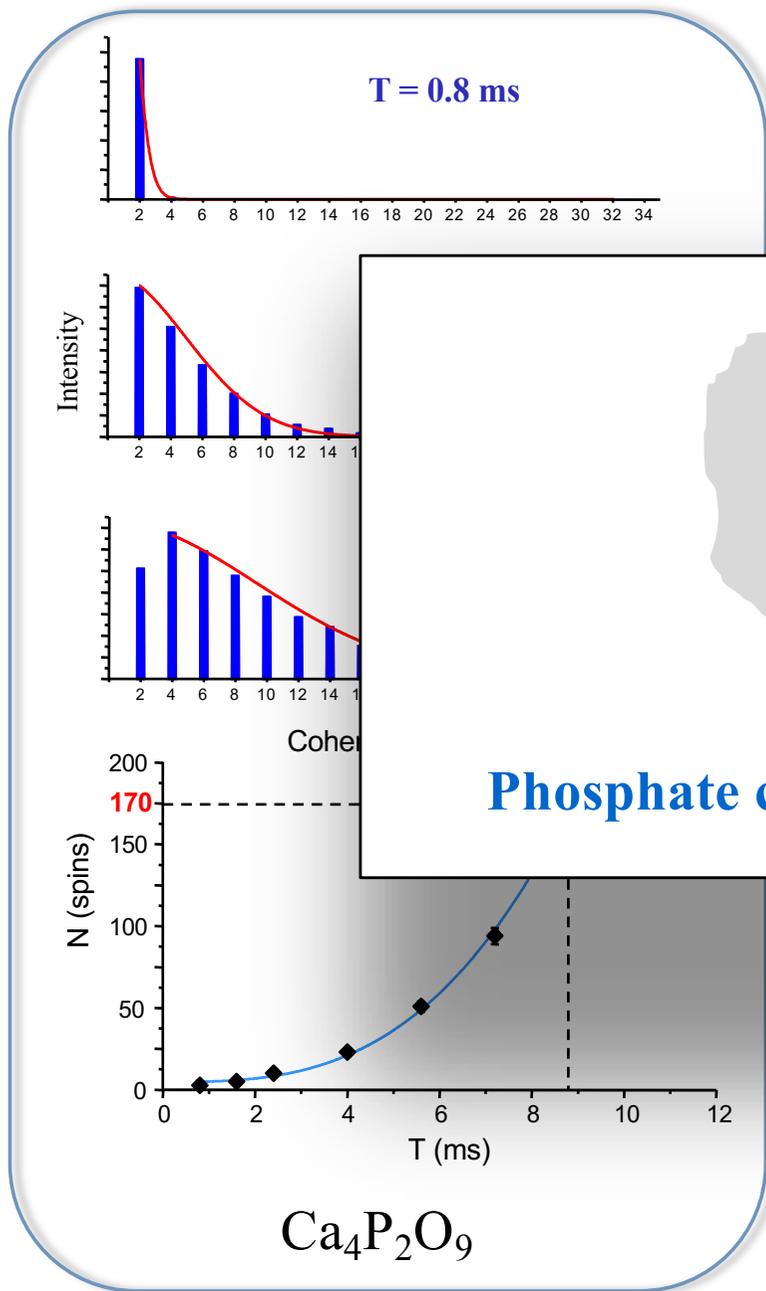


TEM (5.0 mol.% P_2O_5)

- Very weak Z-contrast between Si and P: **No information from SAXS and TEM**

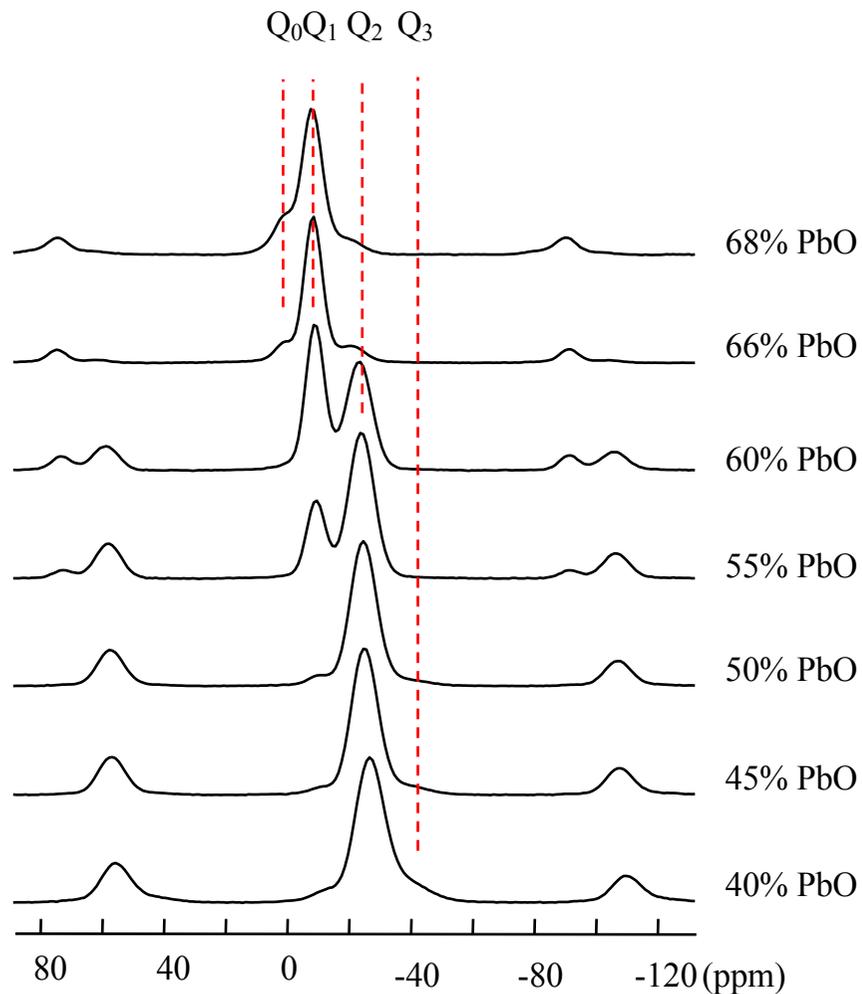
Counting ^{31}P neighbors by spin-counting dipolar NMR ?





Fayon et al., *JACS* 119 (1997), *JNCS* 232-234 (1998), *JNCS* 243 (1999), *JMR* 137 (1999), *Inorg. Chem.* 38 (1999).

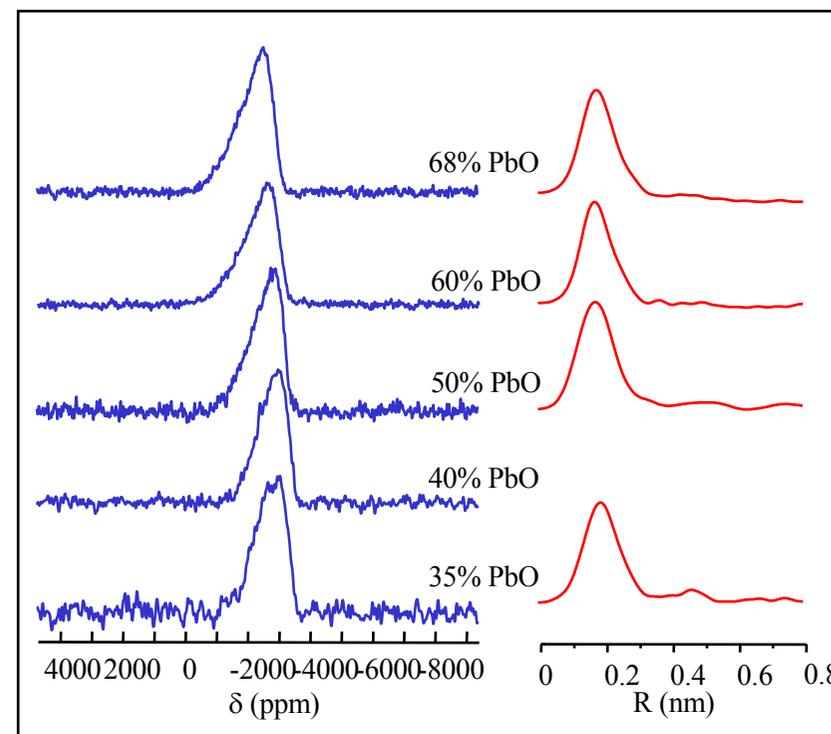
^{31}P MAS NMR spectra



^{207}Pb static NMR spectra

NMR
static or MAS

XAFS



$\text{P}_2\text{O}_5\text{-PbO}$

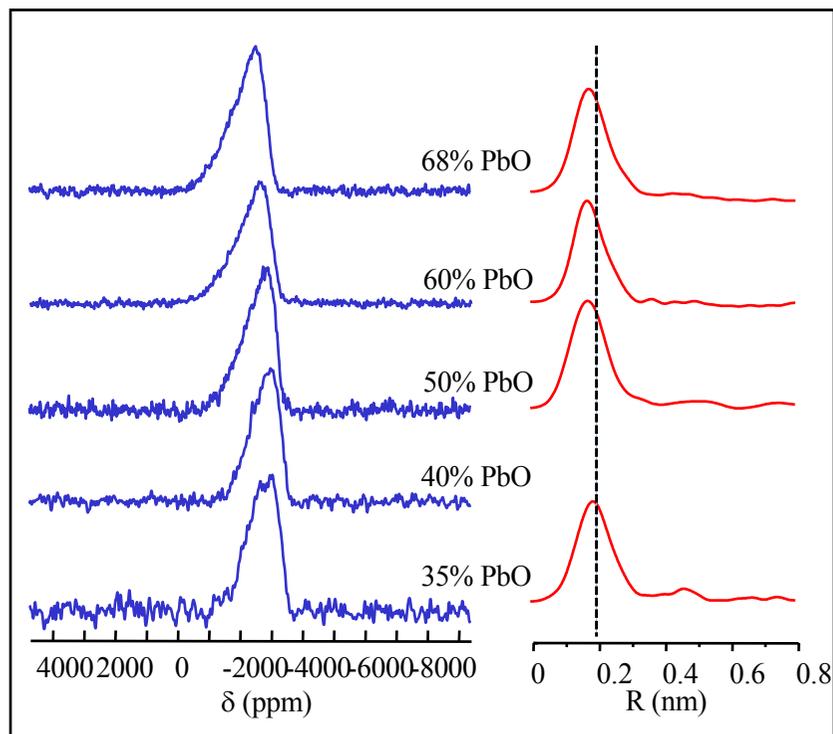
↪ CN ~ 9 to 7

↪ d(Pb-O) ~ 0.26-7nm

Fayon et al., *JACS* 119 (1997), *JNCS* 232-234 (1998), *JNCS* 243 (1999), *JMR* 137 (1999), *Inorg. Chem.* 38 (1999).

NMR
static or MAS

XAFS

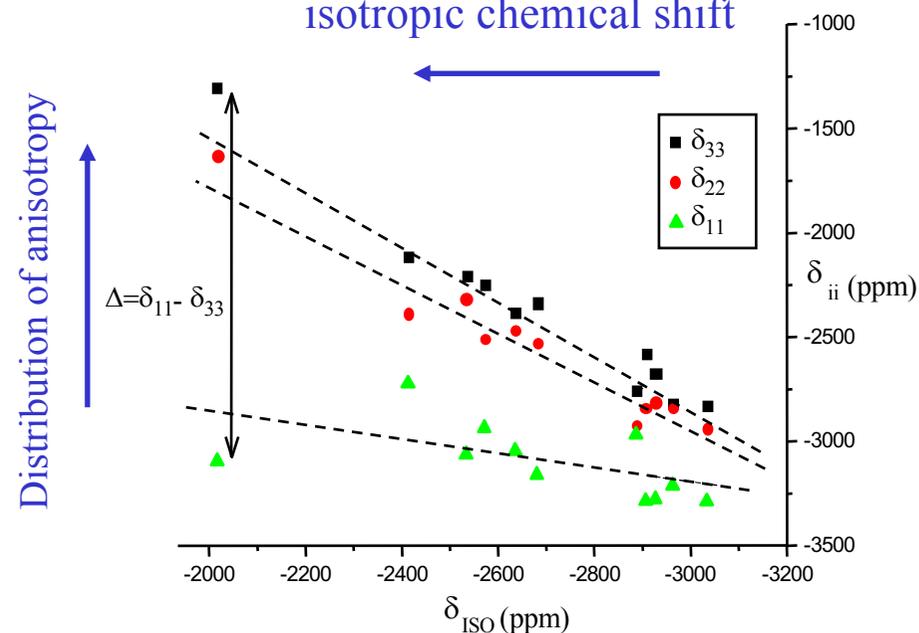


$\text{P}_2\text{O}_5\text{-PbO}$

↪ CN ~ 9 to 7

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Distribution of
isotropic chemical shift

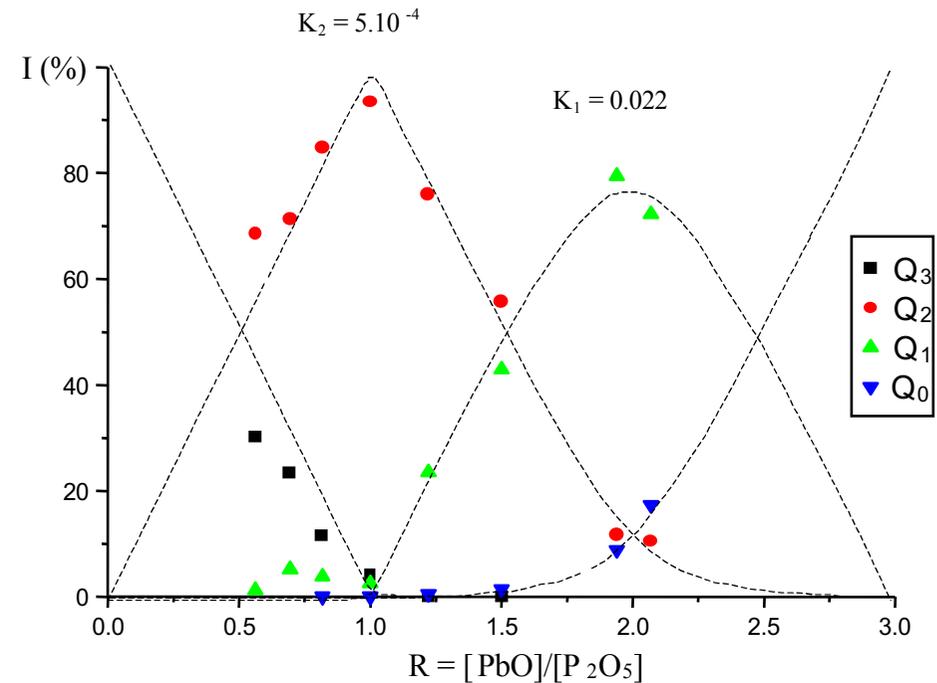
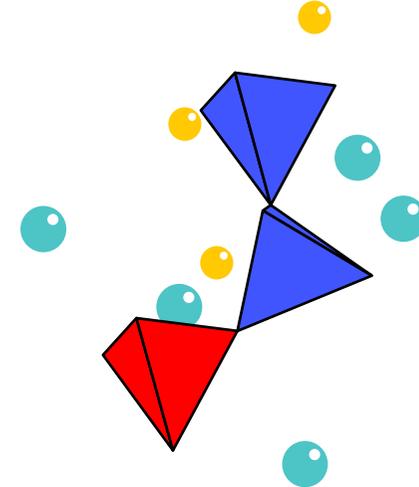
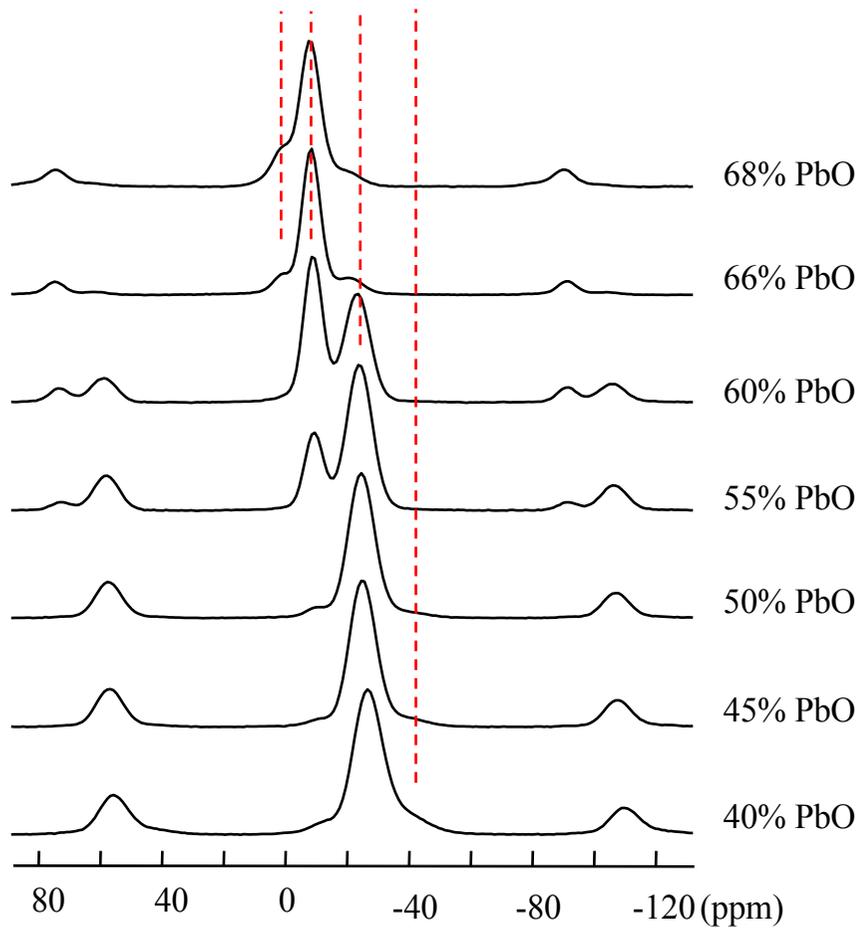


**Distribution of ionic to covalent
chemical binding of Pb in the network**

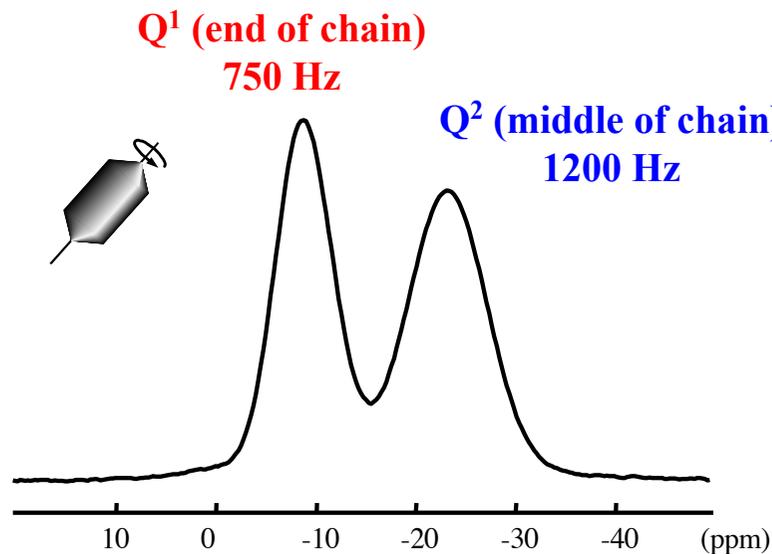
Fayon et al., *JACS* 119 (1997), *JNCS* 232-234 (1998), *JNCS* 243 (1999), *JMR* 137 (1999), *Inorg. Chem.* 38 (1999).

^{31}P MAS NMR spectra

$Q_0 Q_1 Q_2 Q_3$

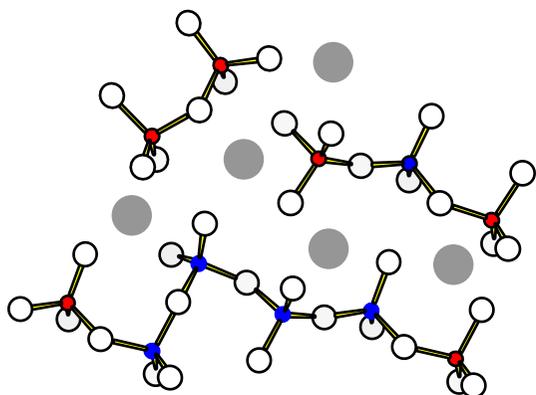


$(\text{PbO})_{0.61}(\text{P}_2\text{O}_5)_{0.39}$ glass



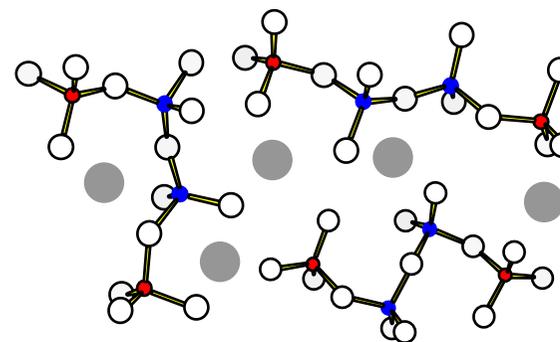
$$[\text{Q}^1] = [\text{Q}^2]$$

Average chain length
 $N_{\text{av.}} \sim 4$



Chain length distribution?
Chemical disorder

? Nature of disorder at the nanometric scale ?



Chain geometries?
Topological or geometrical disorder

P-O-P chemical bonds can be viewed from J based P-P experiments

2D NMR

The unpublished Baško Polje (1971) lecture notes about two-dimensional NMR spectroscopy

J. Jeener

Faculté des Sciences (CPI-232), Campus Plaine, Université Libre de Bruxelles, B-1050 Brussels, Belgium

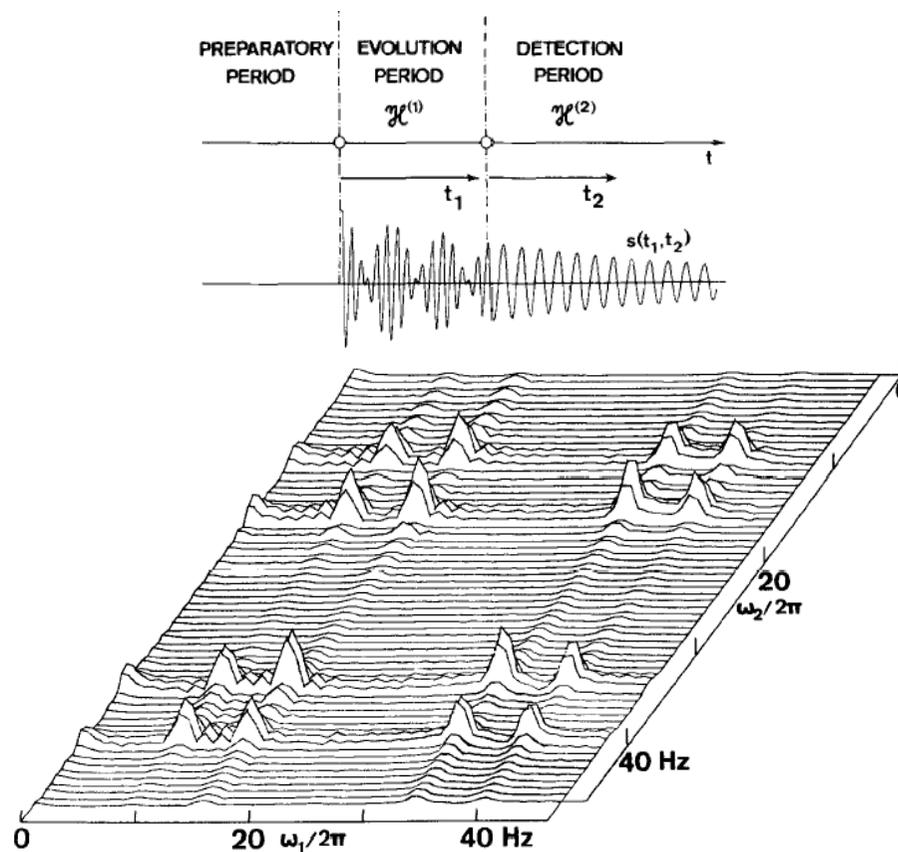
Abstract. — The main part of this paper is a reproduction of (previously unpublished) lecture notes, which were circulated in 1971, and which are often cited as the initiation of two-dimensional NMR spectroscopy. A brief discussion follows, about the way of handling dates and durations in time-dependent quantum mechanics, and about the use of diagrams in NMR pulse spectroscopy in the usual or the superoperator formalisms.

In « *NMR and More* », M. Goldman et M. Porneuf Eds (1994)

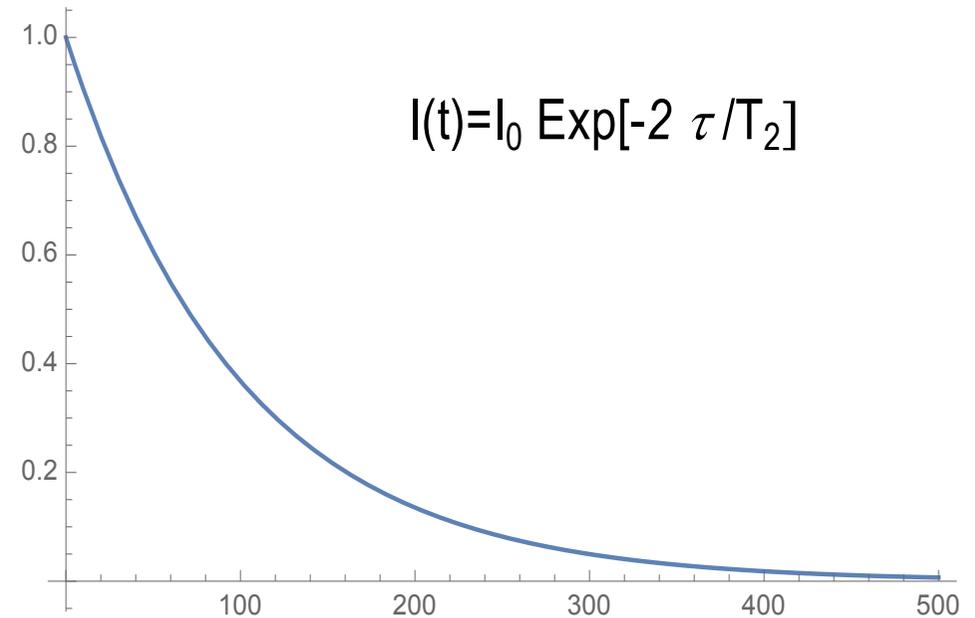
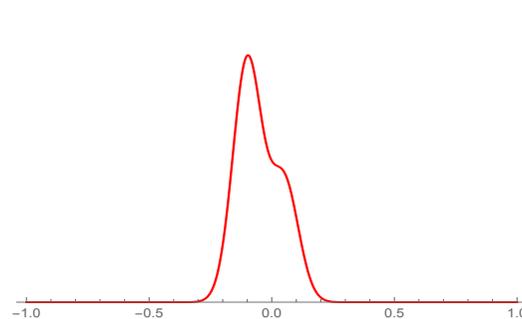
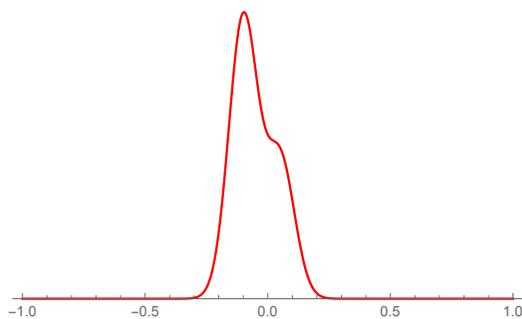
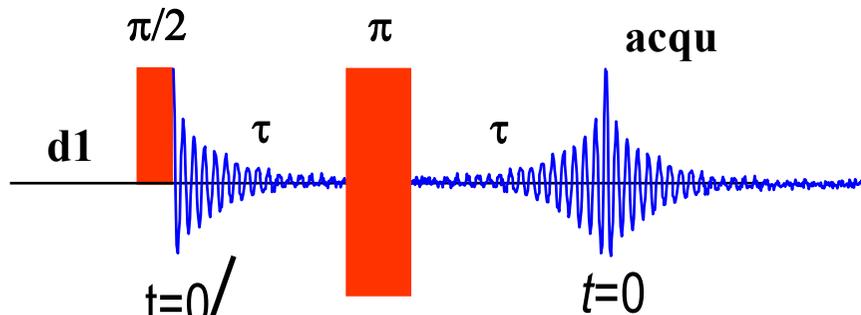
Two-dimensional spectroscopy. Application to nuclear magnetic resonance

W. P. Aue, E. Bartholdi, and R. R. Ernst

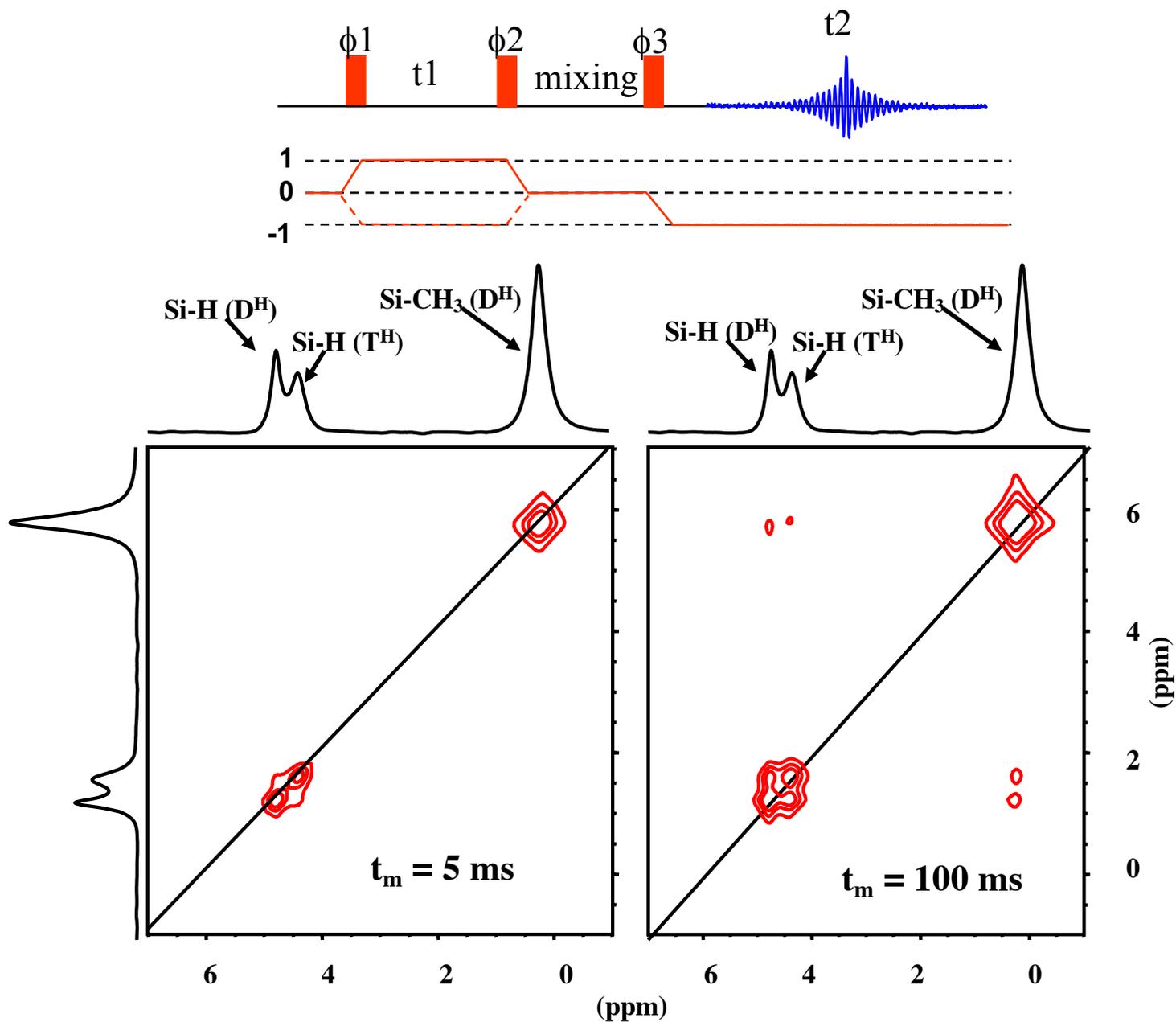
Laboratorium für physikalische Chemie, Eidgenössische Technische Hochschule, 8006 Zürich, Switzerland
(Received 13 November 1975)

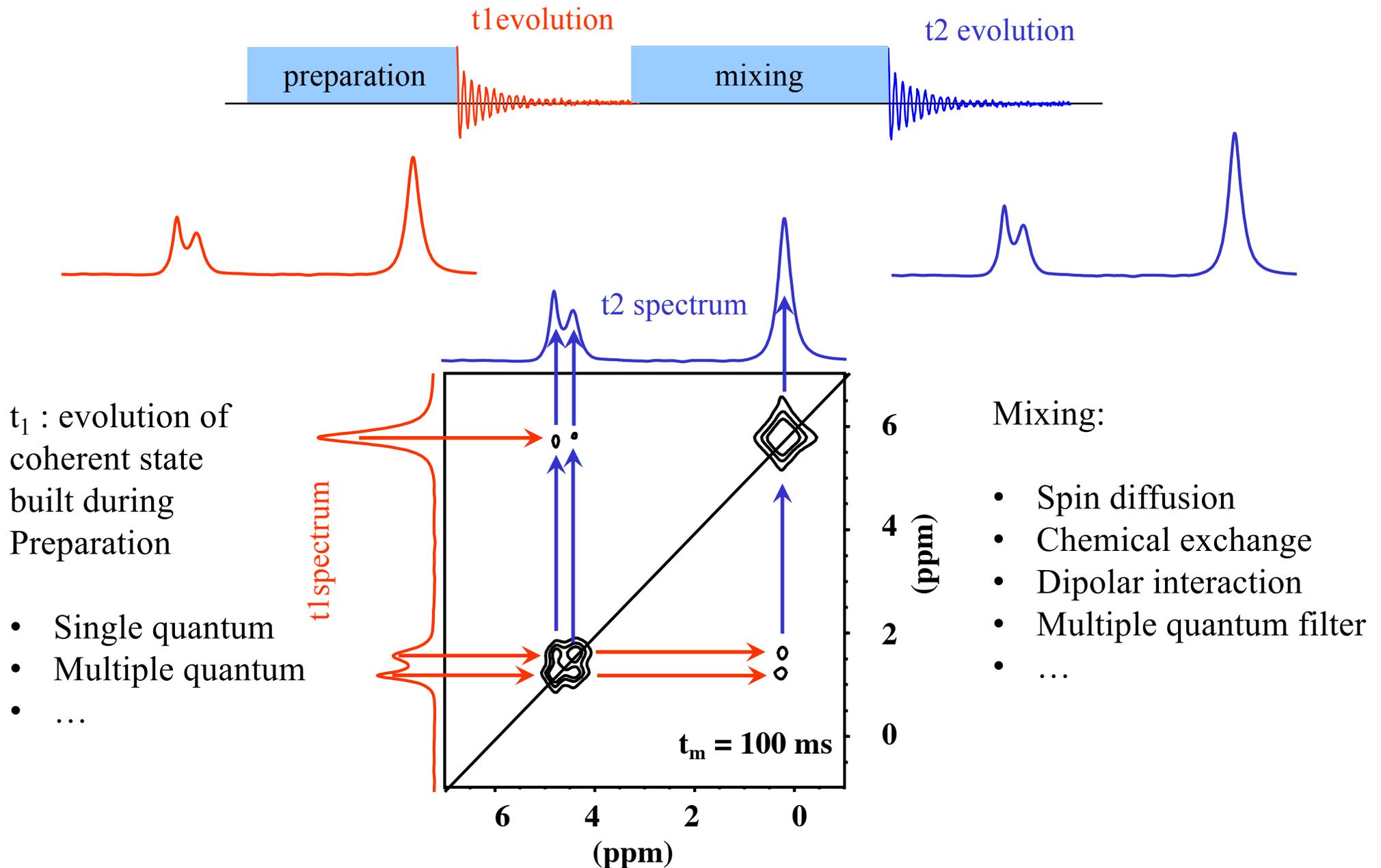


J. Chem. Phys., **64**, 2229 (1974)

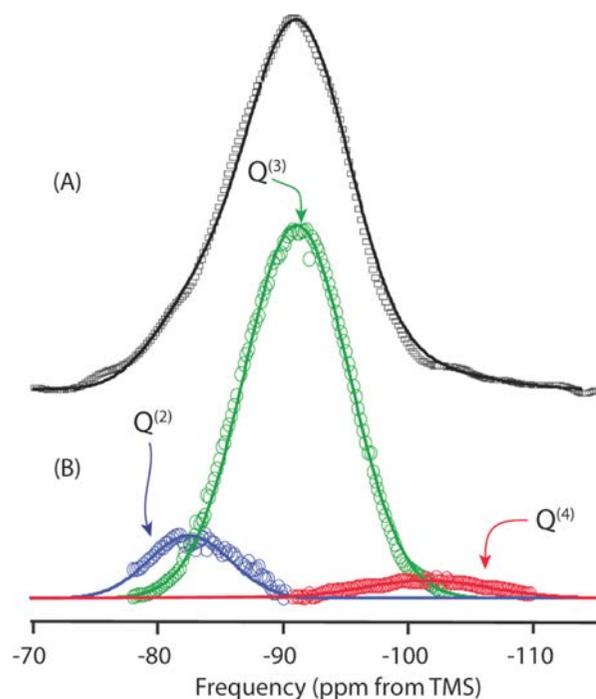
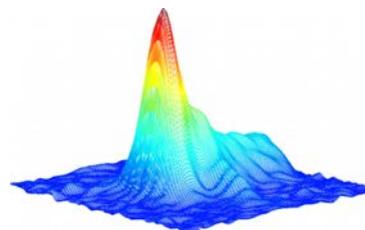


Measurement of T_2
 the life time of the coherence
 in the XY [transverse] plane

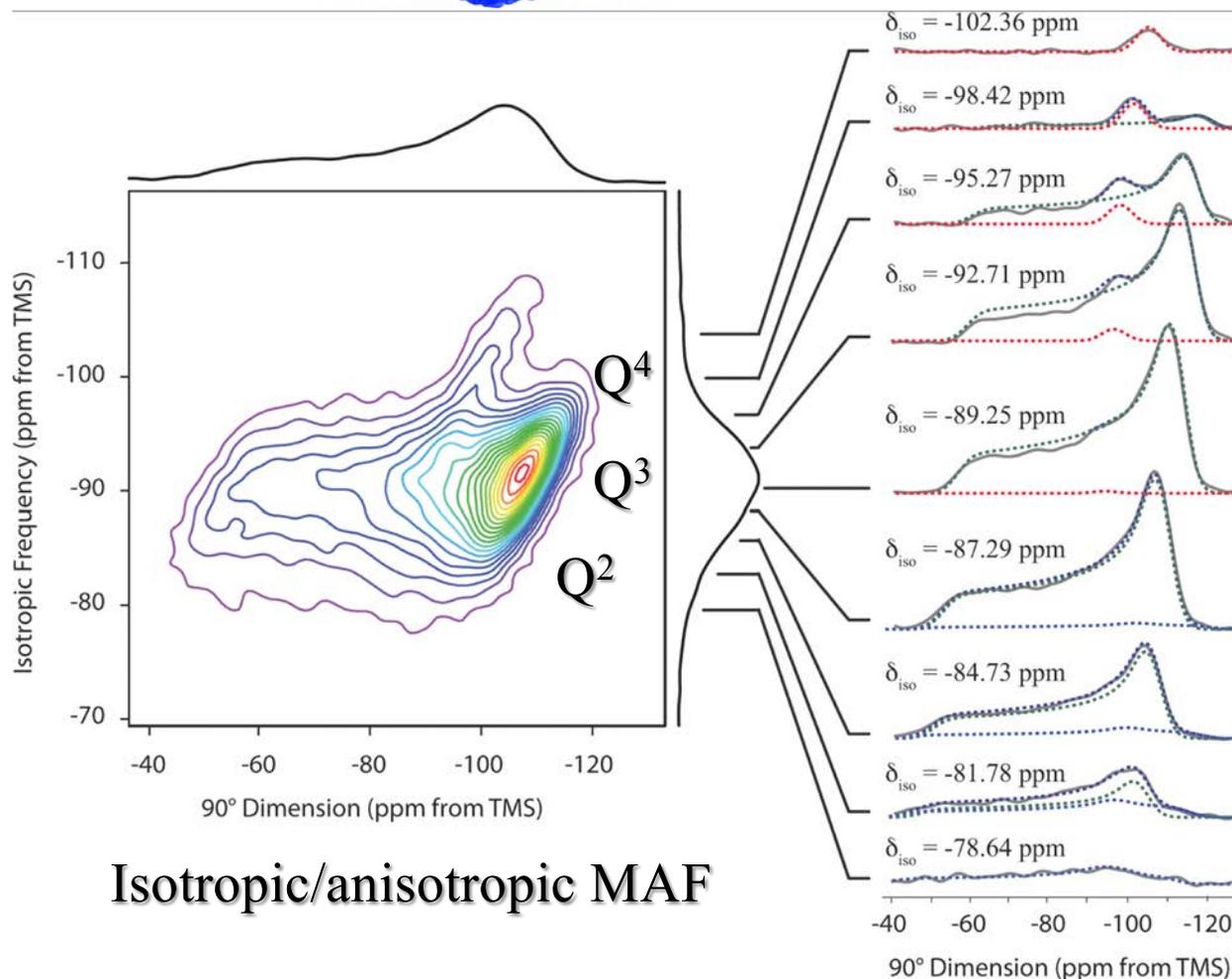




^{29}Si MAF Spectrum of $\text{K}_2\text{O}-2\text{SiO}_2$



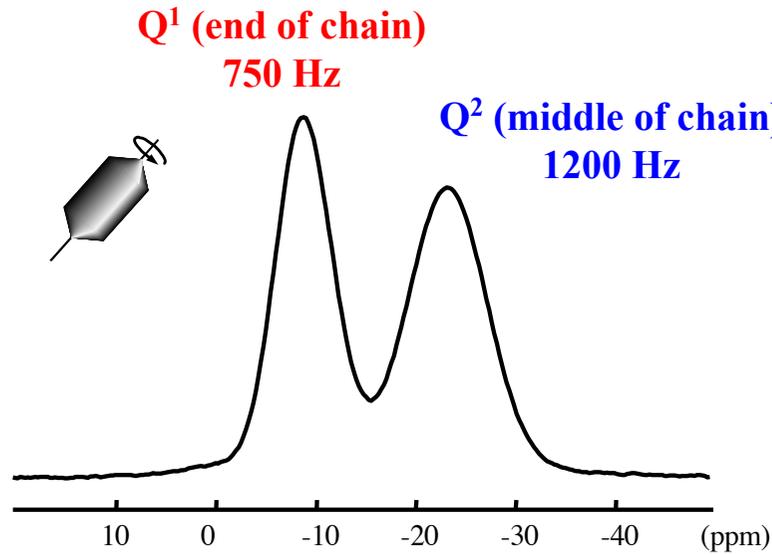
Isotropic MAS



Isotropic/anisotropic MAF

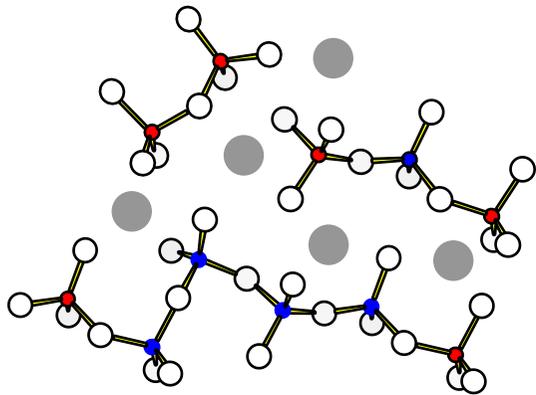
Phosphate Glasses : ^{31}P MAS NMR

$(\text{PbO})_{0.61}(\text{P}_2\text{O}_5)_{0.39}$ glass



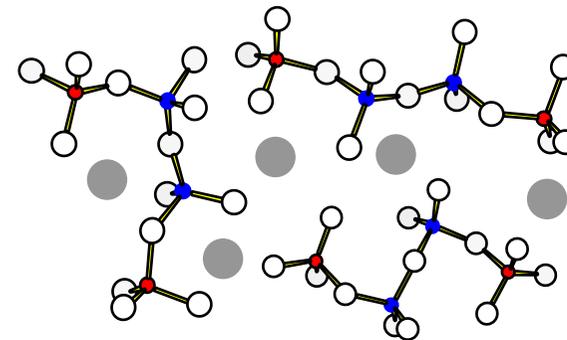
$$[Q^1] = [Q^2]$$

Average chain length
 $N_{\text{av.}} \sim 4$



Chain length distribution?
Chemical disorder

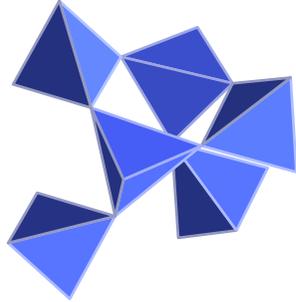
? Nature
of disorder at
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scale ?



Chain geometries?
Topological or geometrical disorder

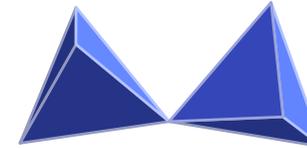
P-O-P chemical bonds can be viewed from J based P-P experiments

Individual ^{31}P

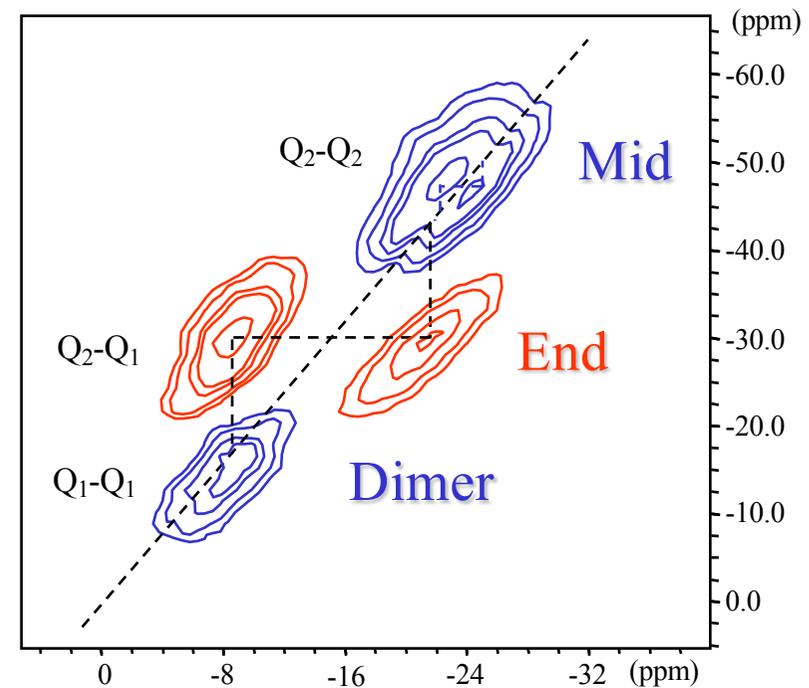
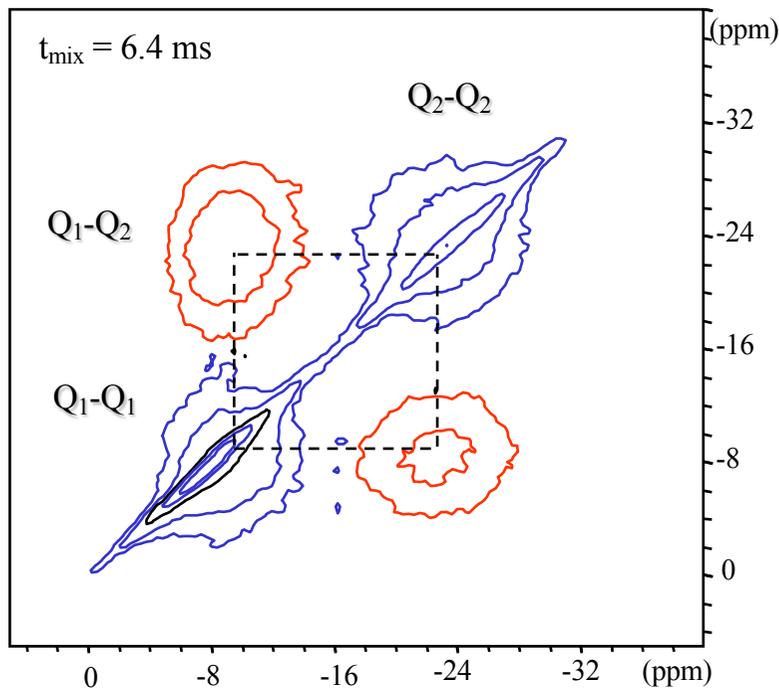


"Spin diffusion"
Spatial proximity

Pairs of ^{31}P

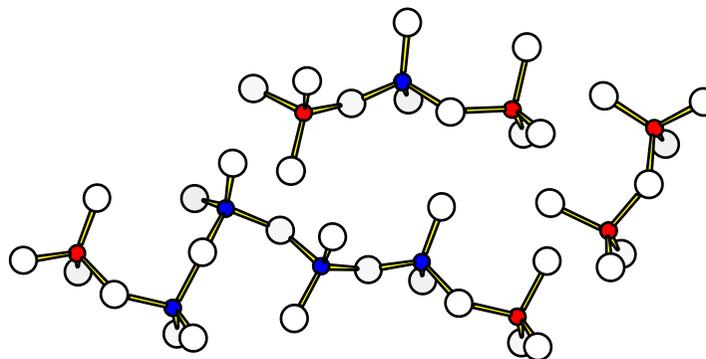
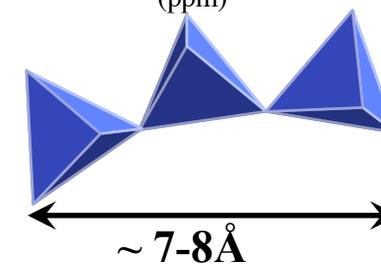
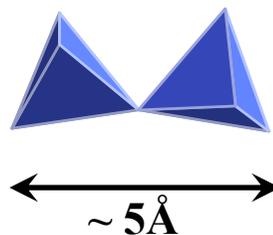
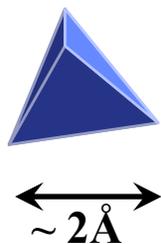
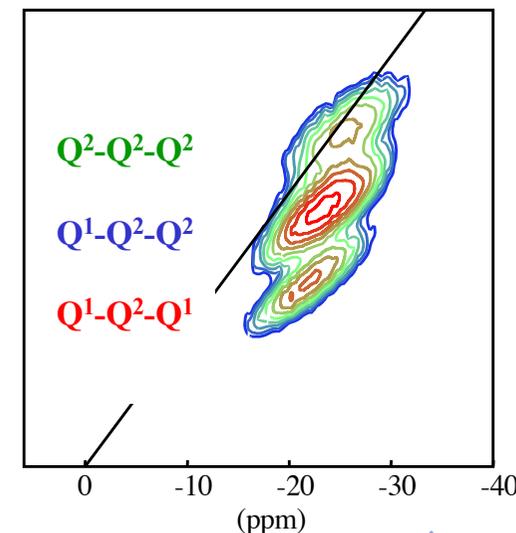
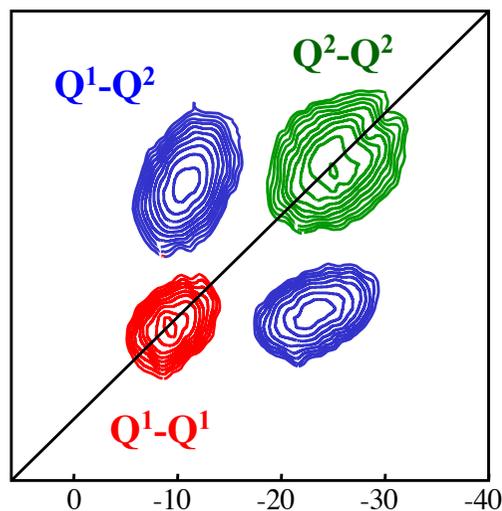
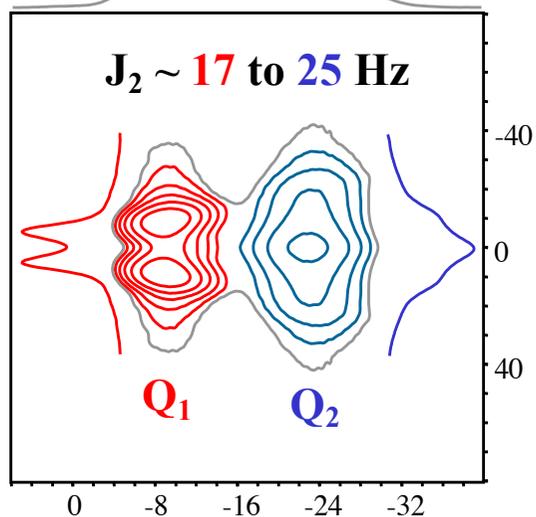


Double Quantum edition
Dipolar : spatial proximity
Jcoupling : chemical bond

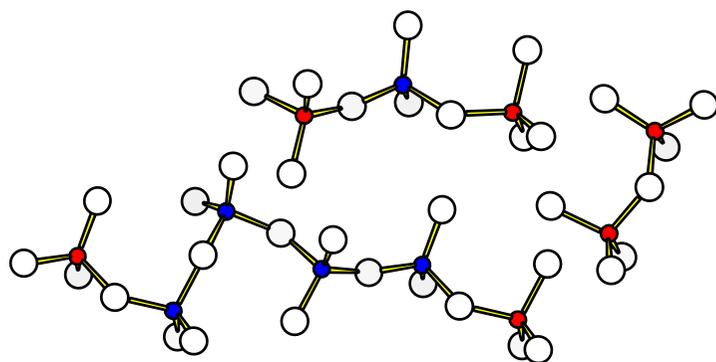
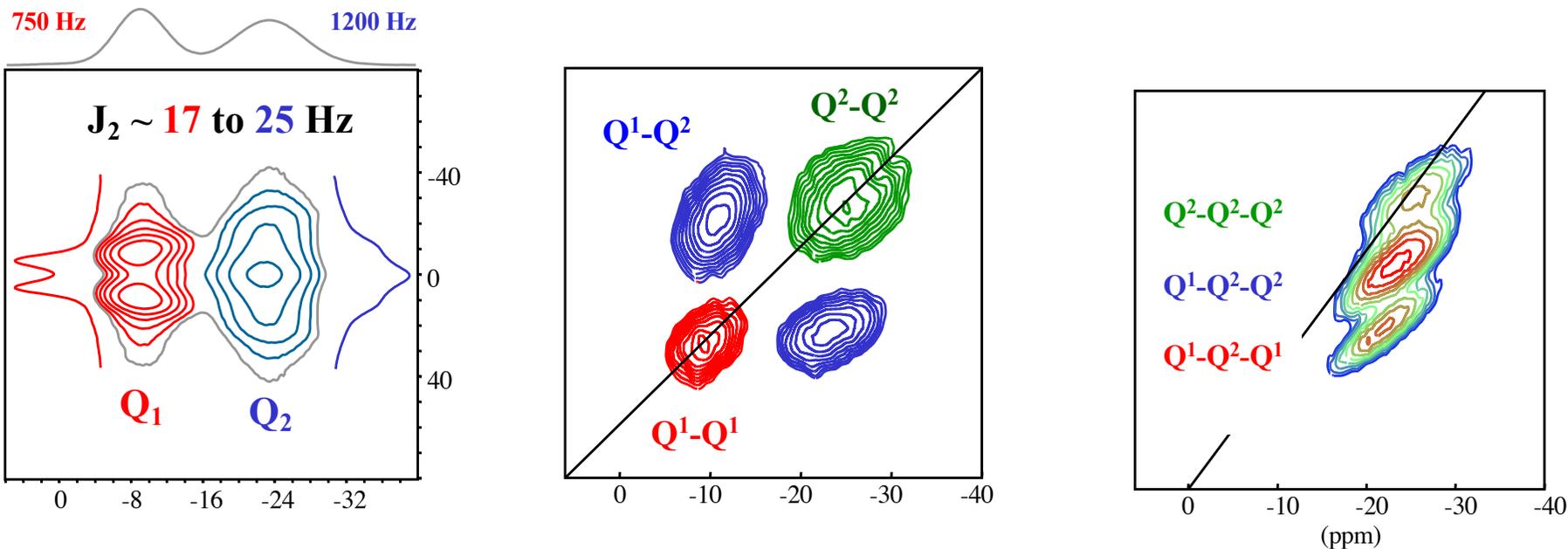


$(\text{PbO})_{0.61}(\text{P}_2\text{O}_5)_{0.39}$ glass

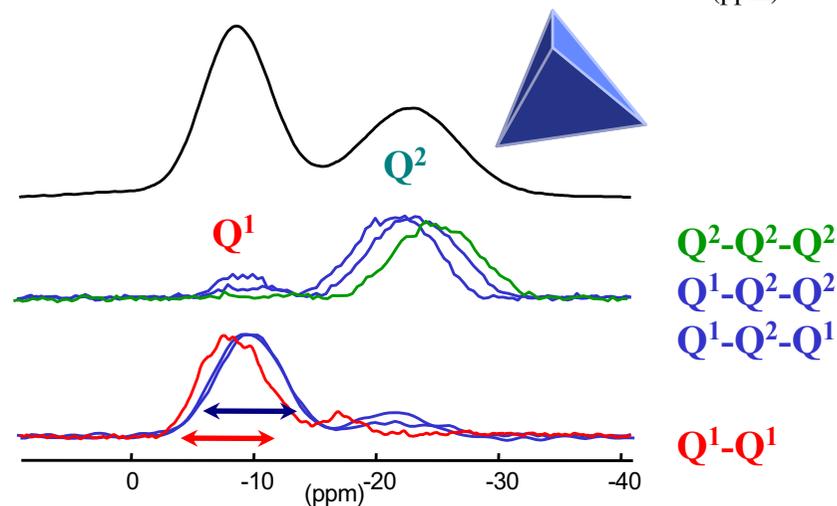
750 Hz 1200 Hz



$(\text{PbO})_{0.61}(\text{P}_2\text{O}_5)_{0.39}$ glass

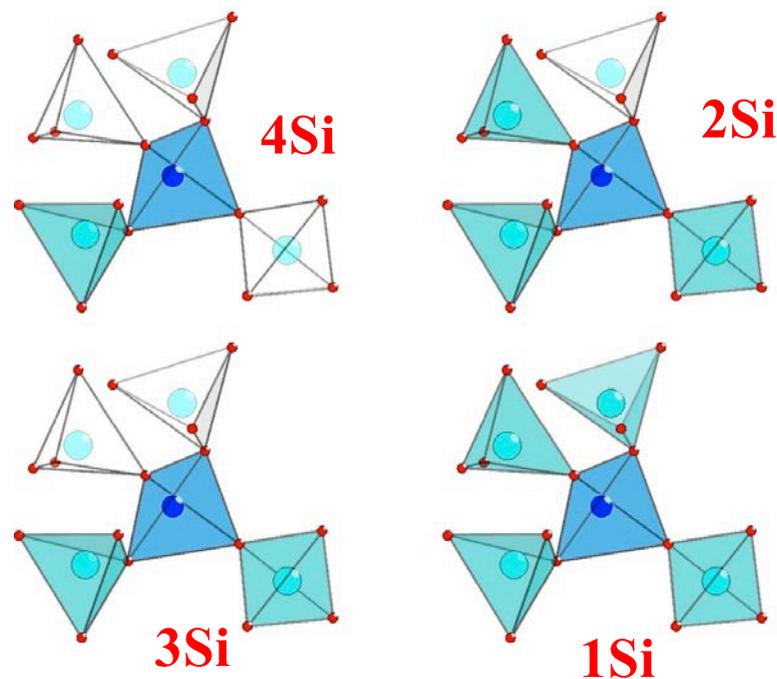
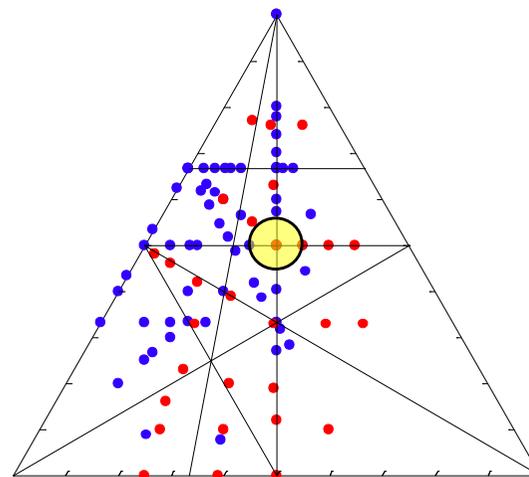
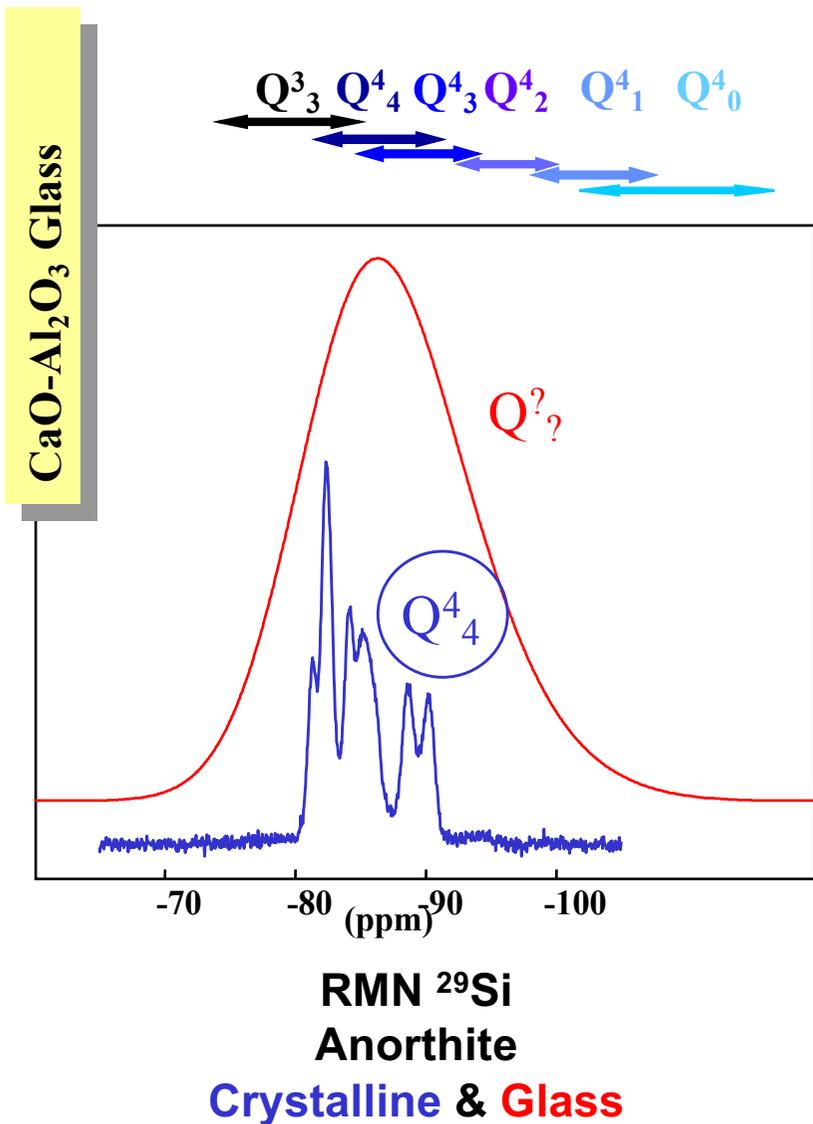


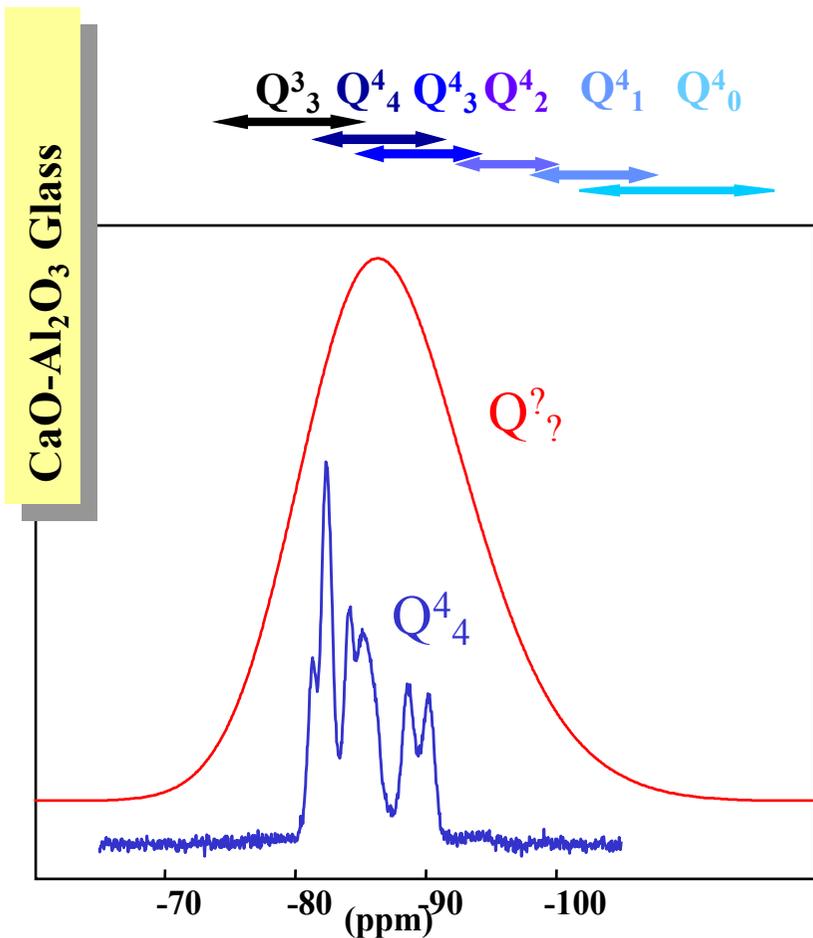
Topology



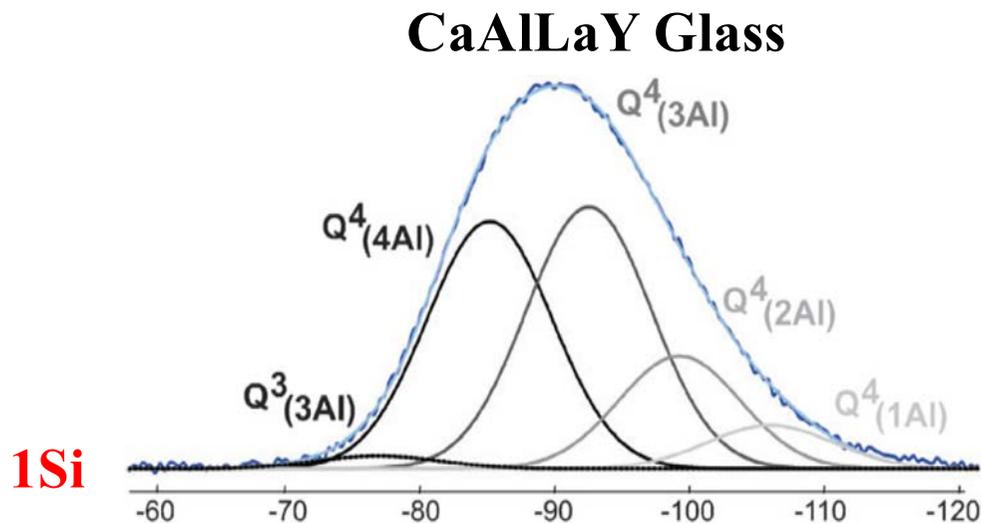
Geometry

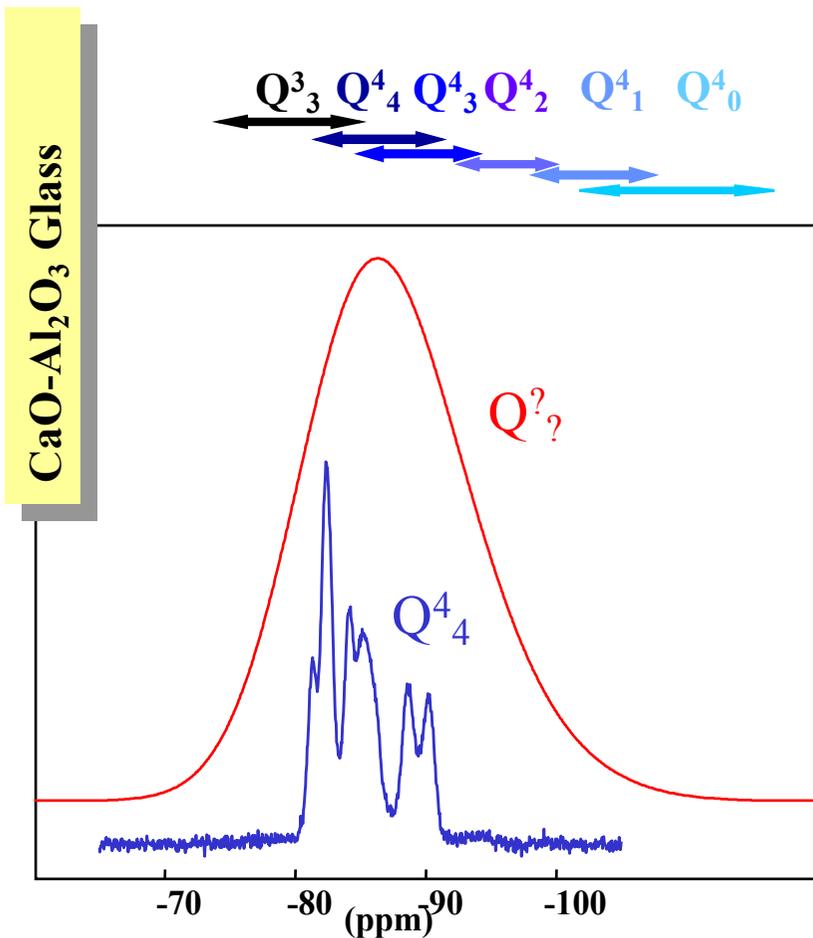
“Chemistry”



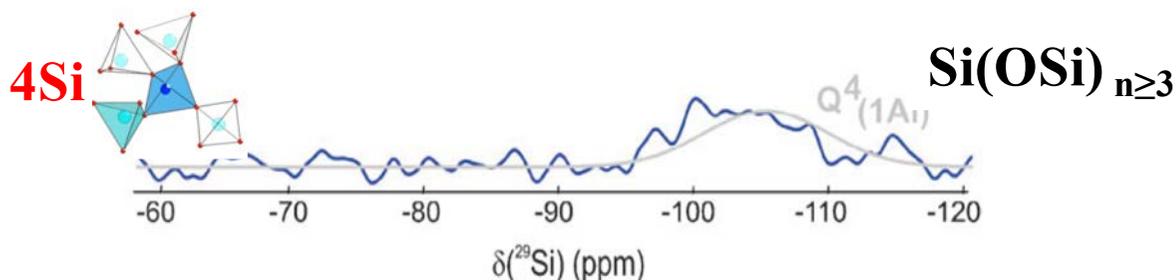
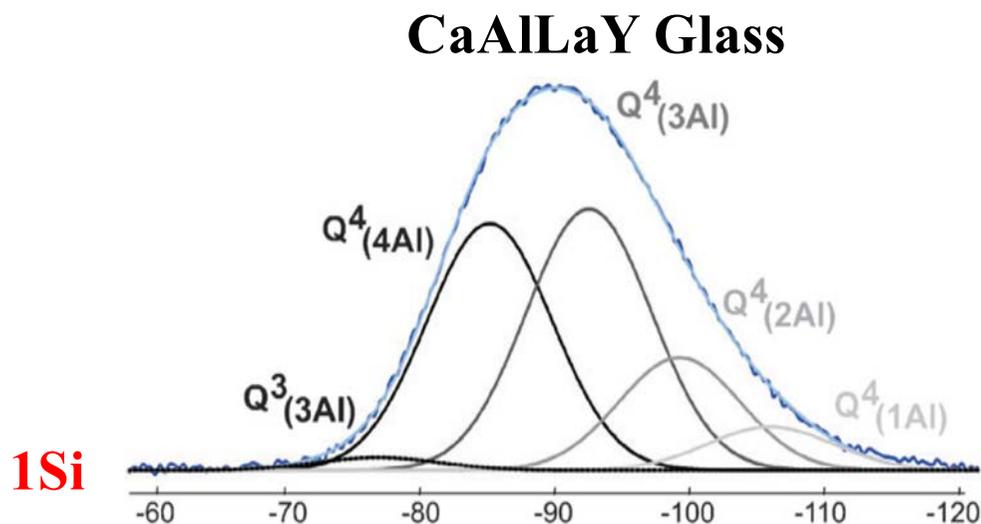


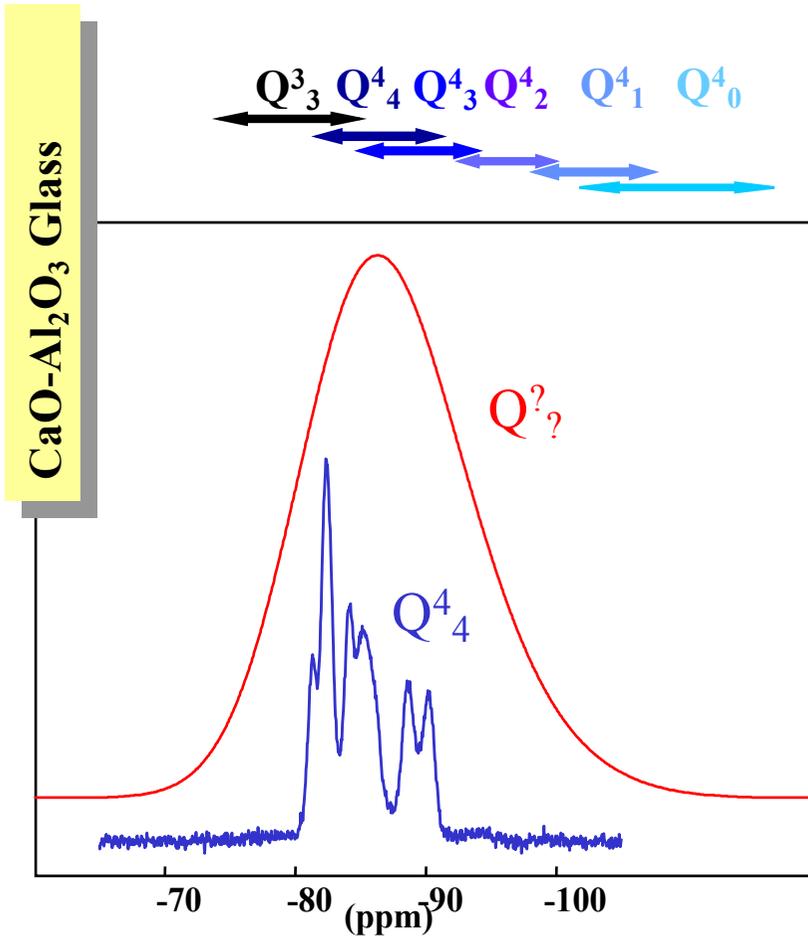
NMR²⁹Si
Anorthite
Crystalline & Glass





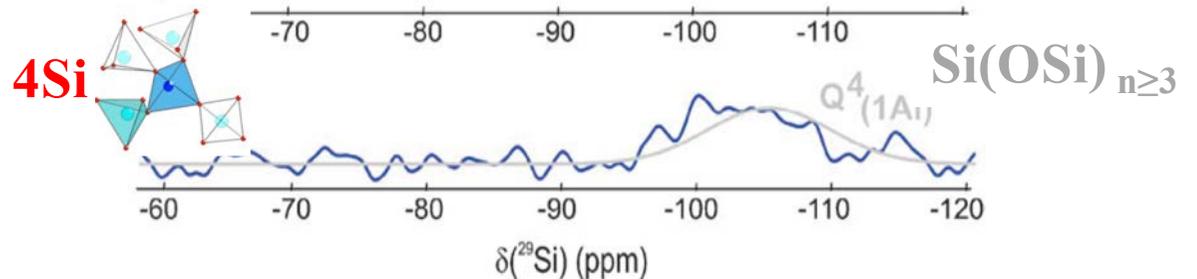
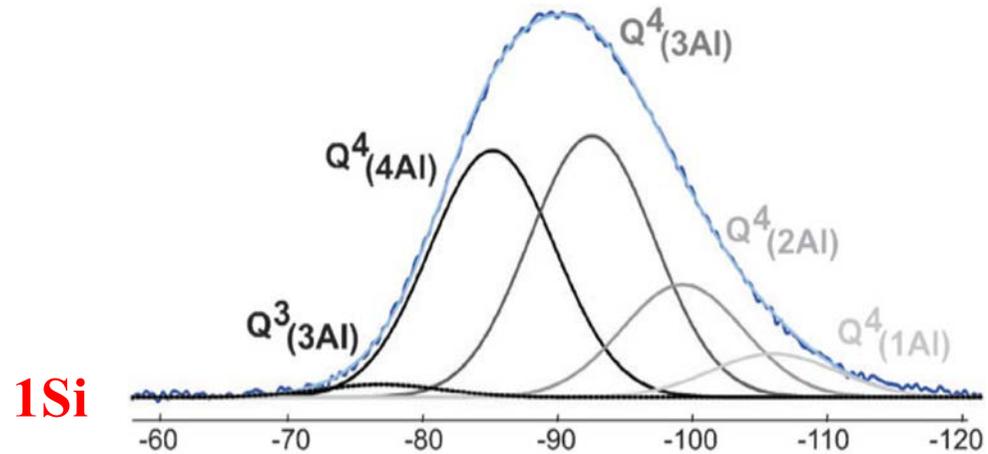
NMR²⁹Si
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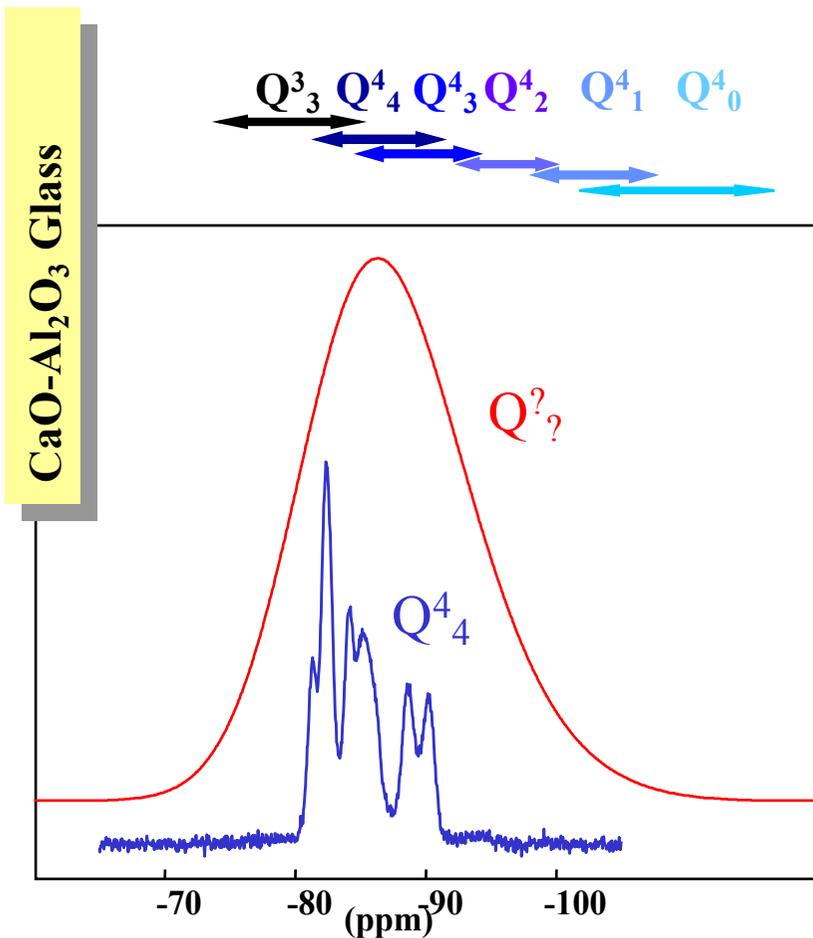




NMR²⁹Si
Anorthite
Crystalline & Glass

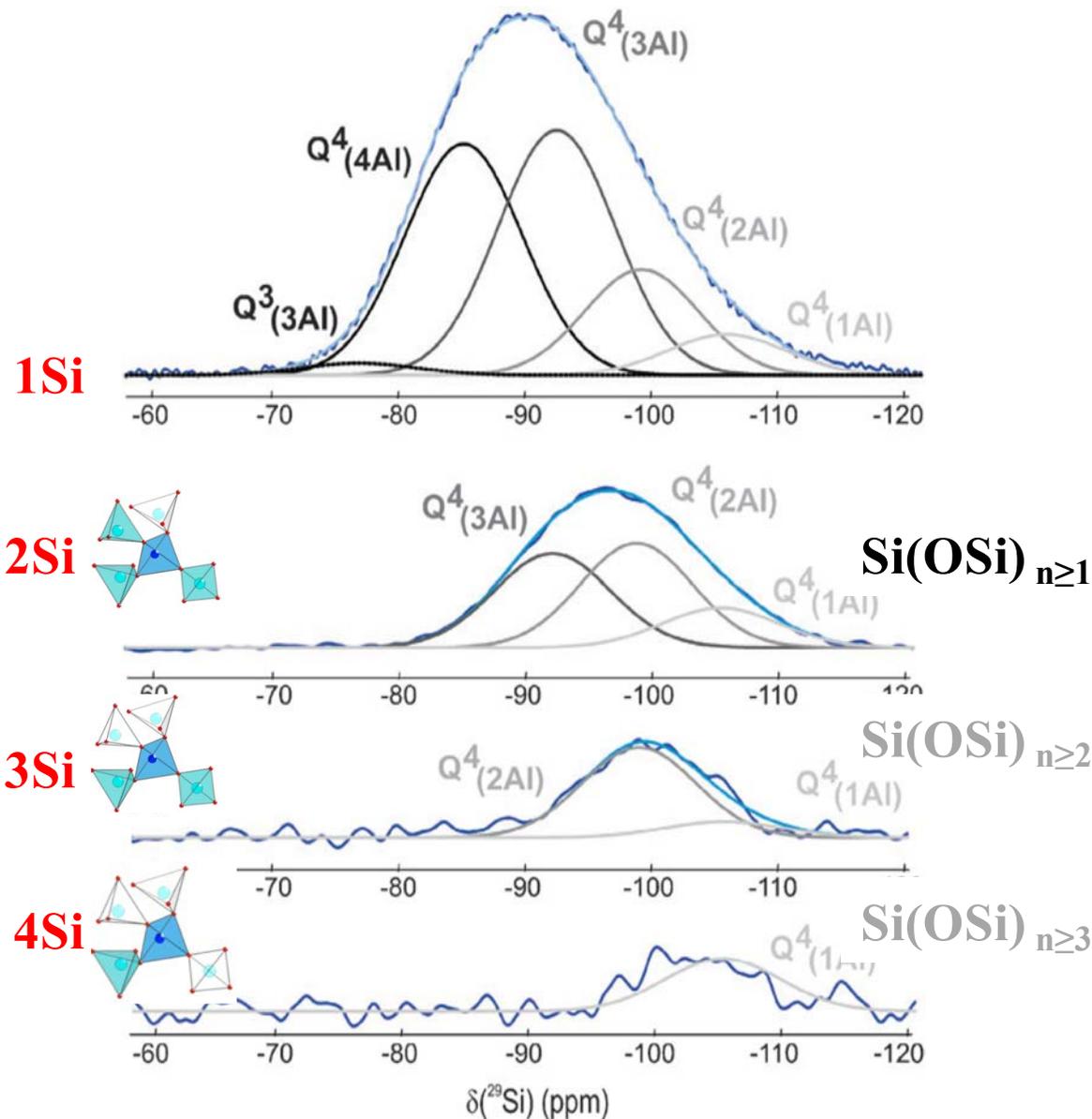
CaAlLaY Glass

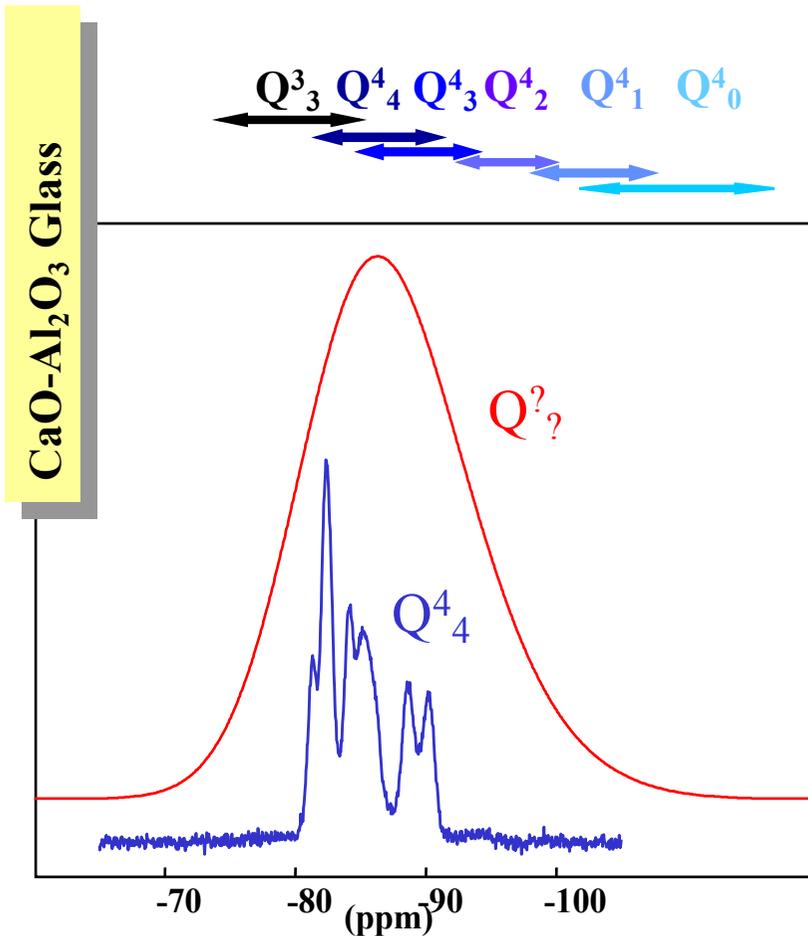




NMR²⁹Si
Anorthite
Crystalline & Glass

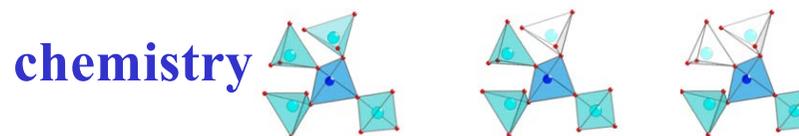
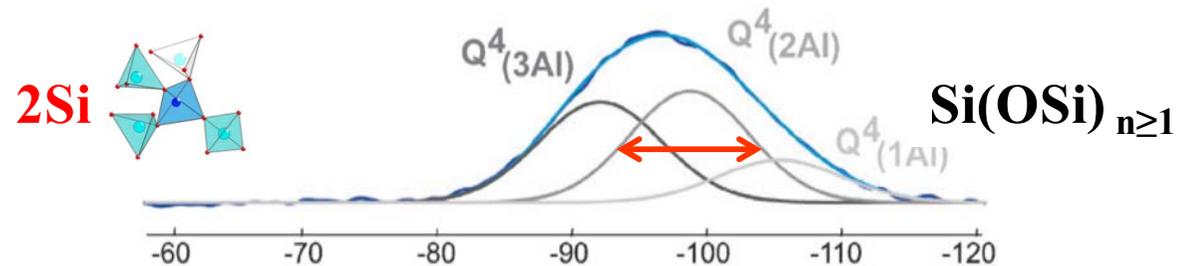
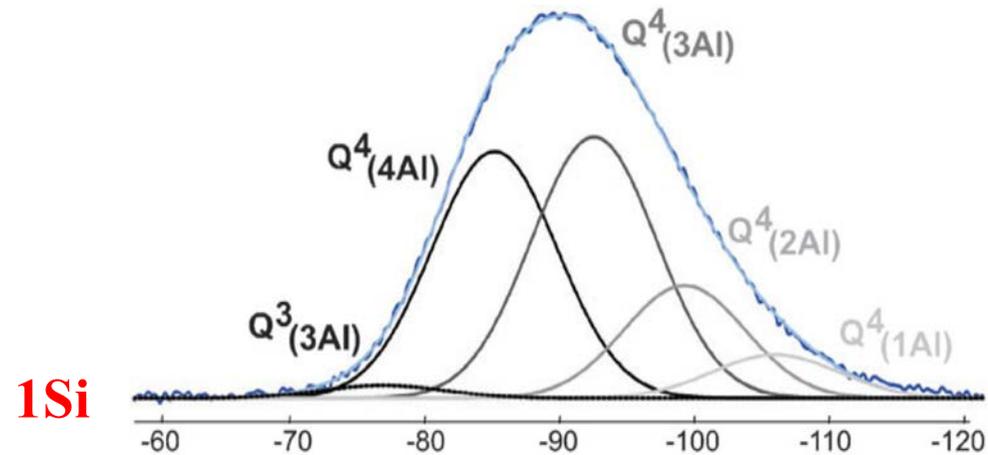
CaAlLaY Glass





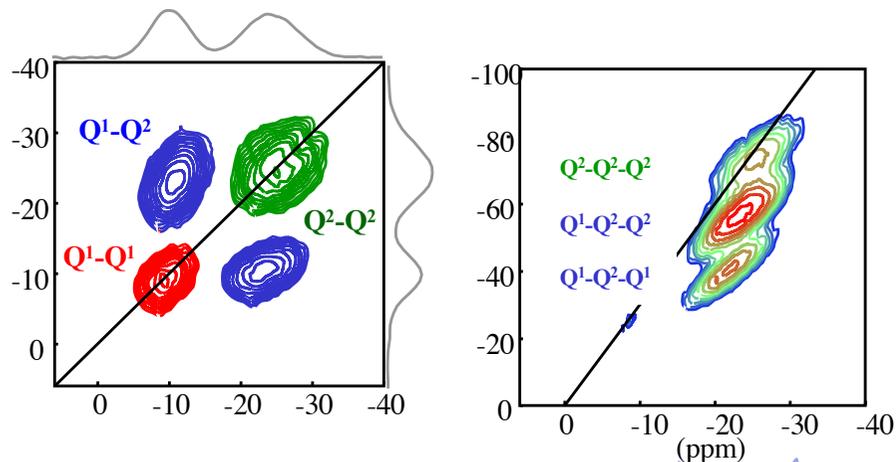
NMR²⁹Si
Anorthite
Crystalline & Glass

CaAlLaY Glass

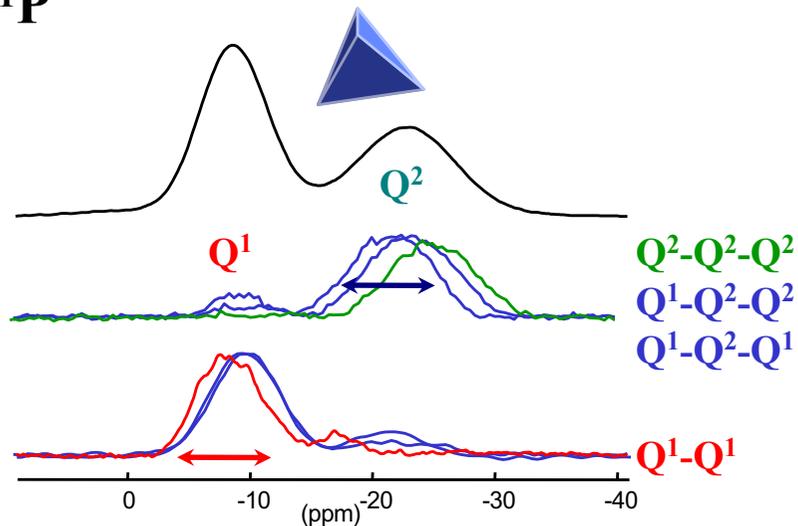


geometry





³¹P

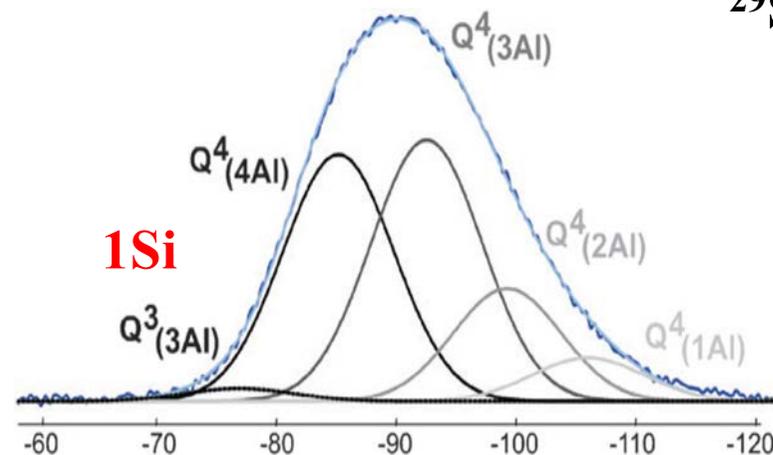


“Chemistry”

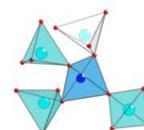
Lead Phosphate glasses

F.Fayon, et al. Journal of Magnetic Resonance 179 50-58 (2006)

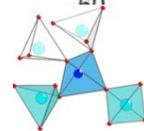
²⁹Si



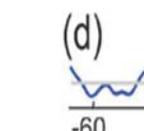
1Si



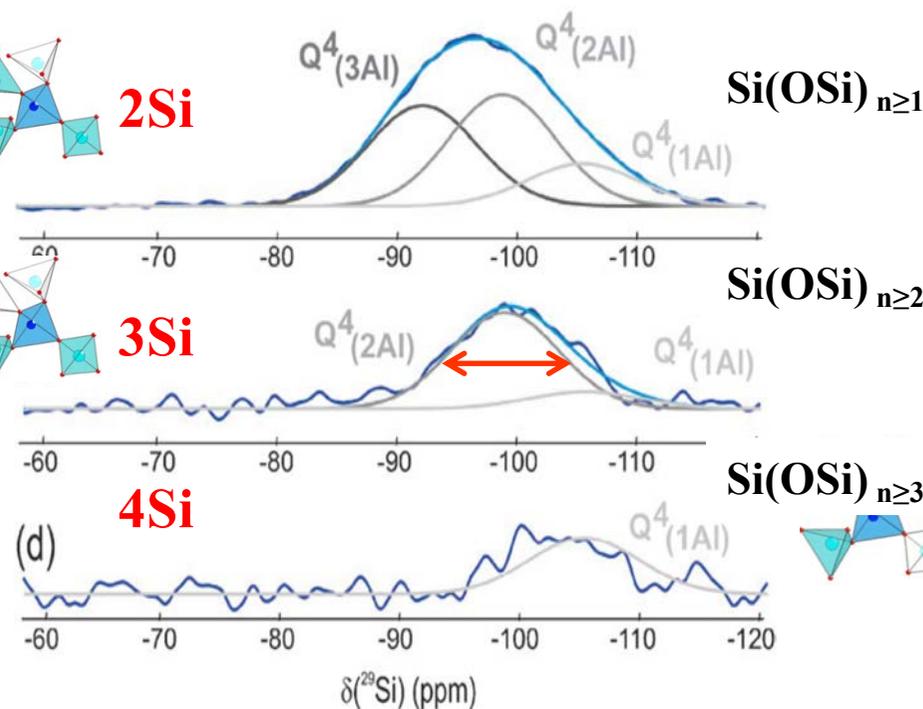
2Si



3Si

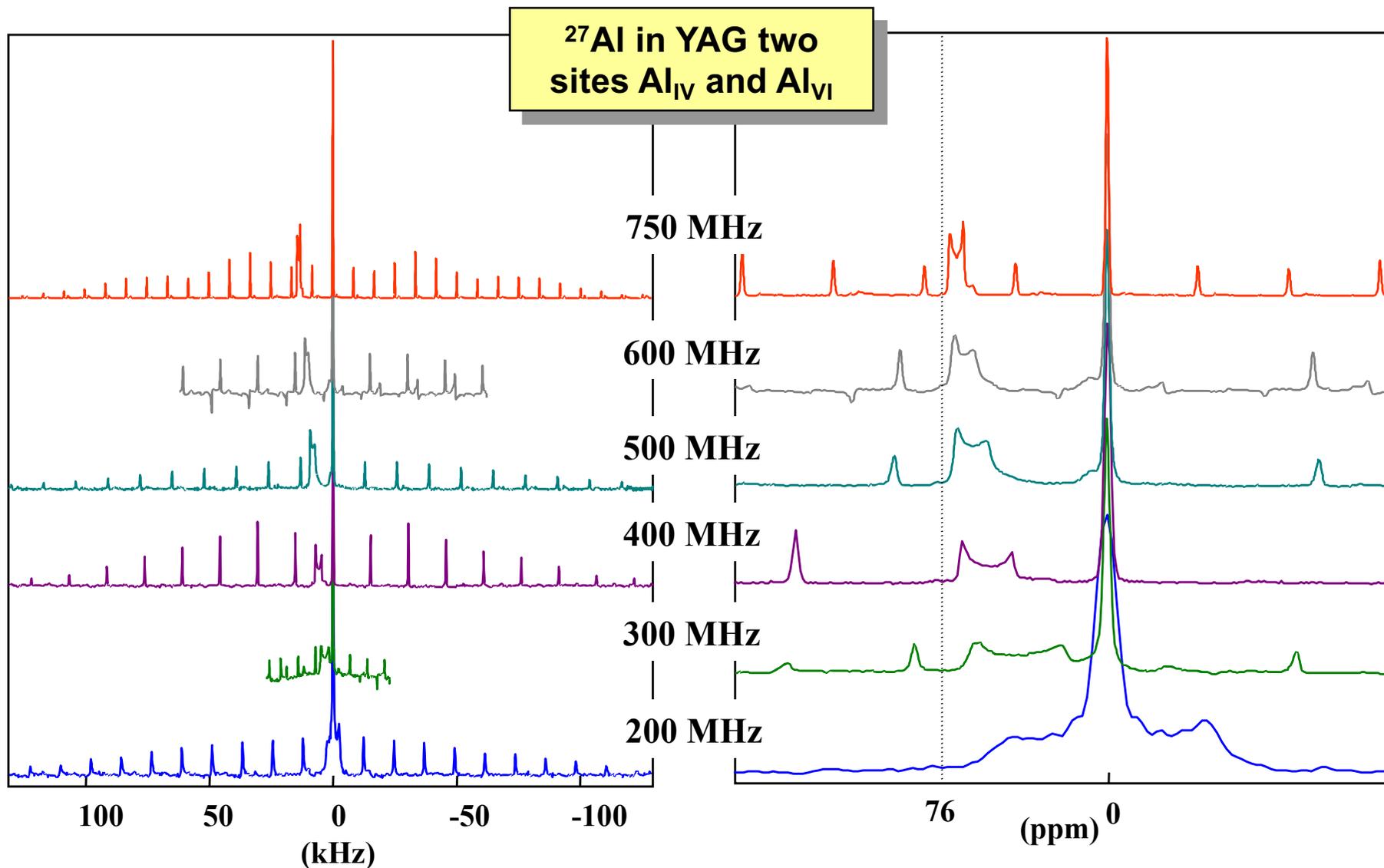


4Si



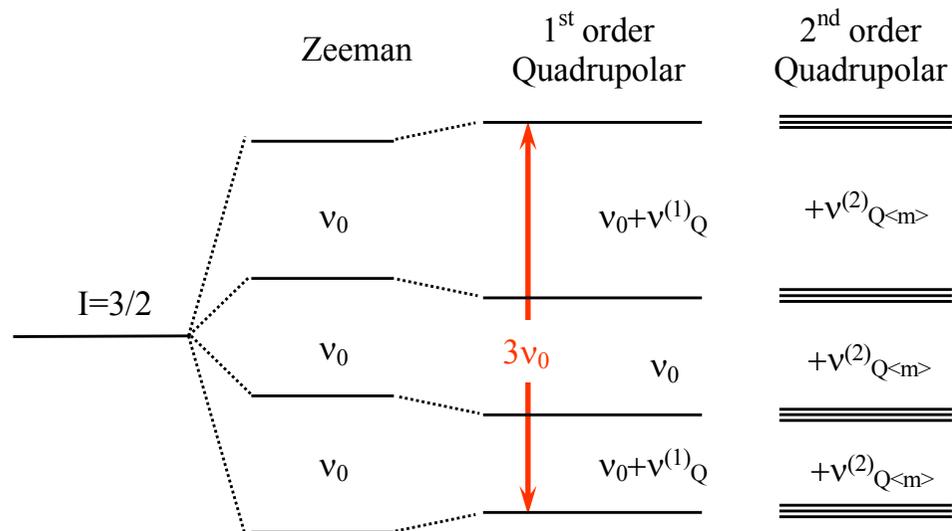
Calcium REE-Alumino-Silicate glasses

J.Hiet, et al. Phys. Chem. Chem. Phys. 11 6935-6940 2009



Second order shift and width goes with ν_Q^2/ν_0 (Hz)

2I+1 energy levels, 2I single quantum transitions



1st order quadrupolar interaction

$$H_Q^{(1)} = C_Q \sqrt{\frac{1}{6}} A_{20}^Q [3I_z^2 - I(I+1)]$$

2nd order quadrupolar interaction

$$H_Q^{(2)} = \frac{C_Q^2}{2\omega_0} \left[A_{2-1}^Q A_{21}^Q \{4I(I+1) - 8I_z^2 - 1\} + A_{2-2}^Q A_{22}^Q \{2I(I+1) - 2I_z^2 - 1\} \right]$$

$$v_{\langle m, m-1 \rangle} = v_0 + v_{\langle m, m-1 \rangle}^{\text{iso}(2)} + v_{\langle m, m-1 \rangle}^{(2)}(\alpha, \beta, \gamma)$$

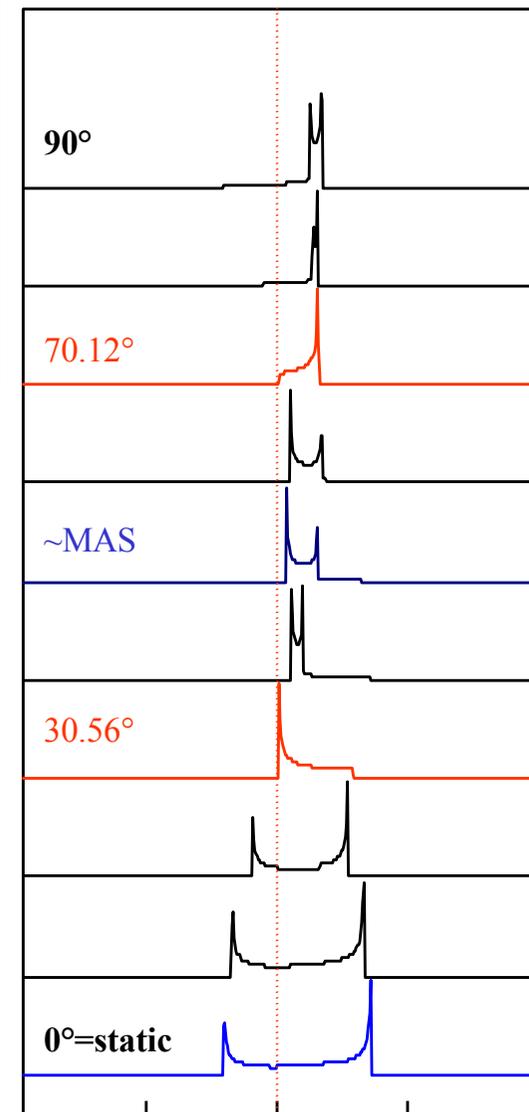
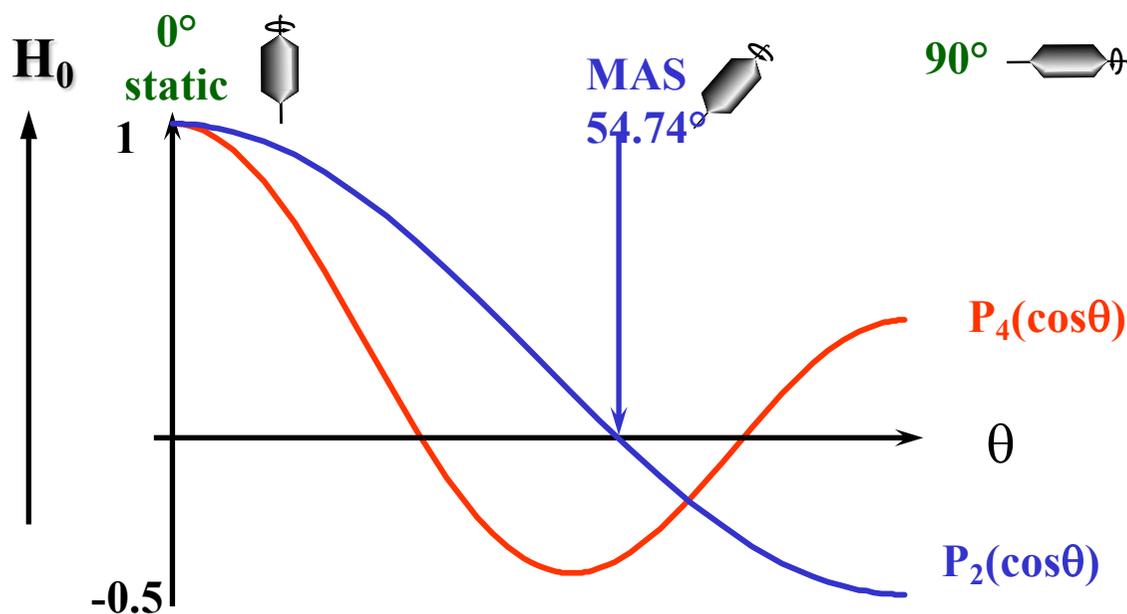
$$v_{\langle m, m-1 \rangle}^{\text{iso}(2)} = -\frac{I(I+1) - 3 - 9m(m-1)}{30} \frac{v_Q^2}{v_0} \sqrt{1 + \frac{\eta^2}{3}}$$

First order $H_I \ll H_Z$

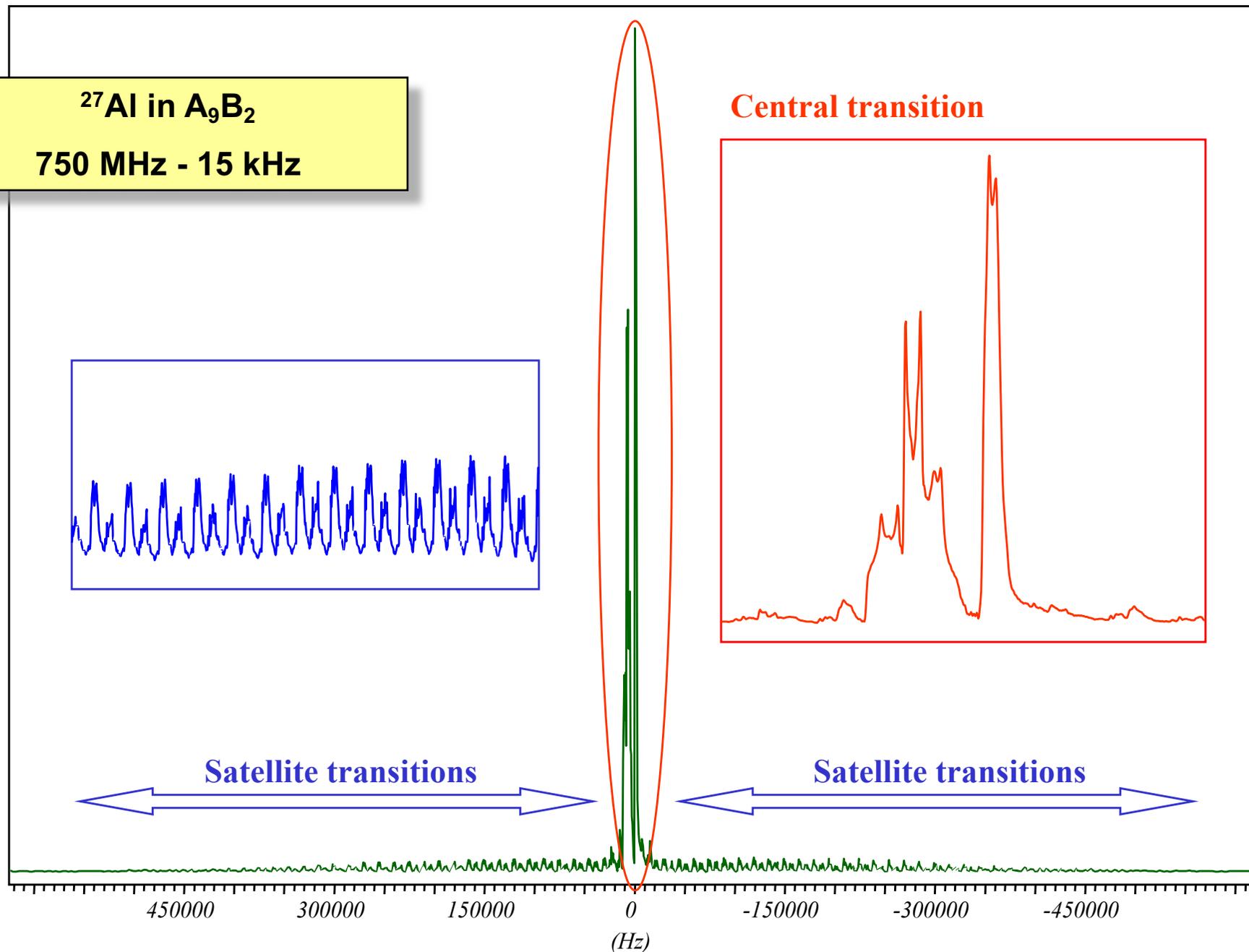
$$\nu_{\langle m, m-1 \rangle} = \nu_0 + P_2(\cos\theta) \frac{(1-2m)}{2} \nu_Q C_{\alpha, \beta} + \sum_{n=1}^2 A_n e^{-i(n\omega_r t + \phi)}$$

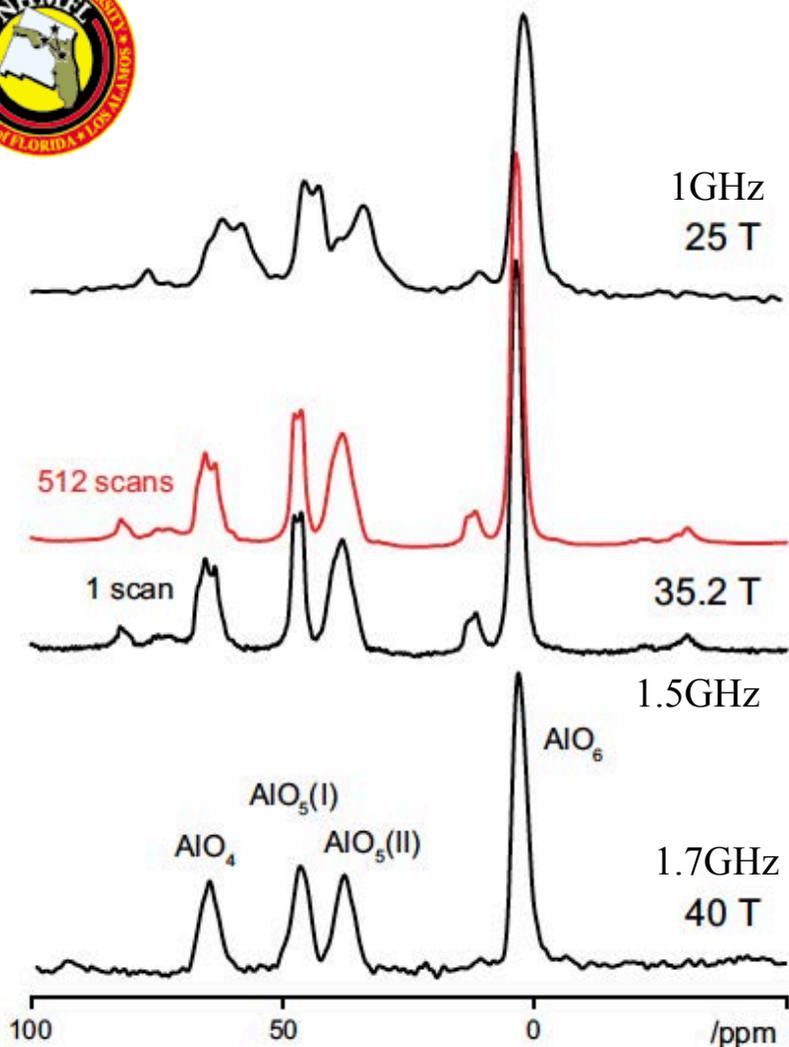
Second order $H_I < H_Z$

$$\nu_{\langle m, m-1 \rangle} = \nu_0 + \sum_{l=0,2,4} C_{m,l}^1 B_l^{(2)} P_l(\cos\theta) + \sum_{n=1}^4 A_n^{(2)} e^{-i(n\omega_r t + \phi)}$$



^{27}Al in A_9B_2
750 MHz - 15 kHz

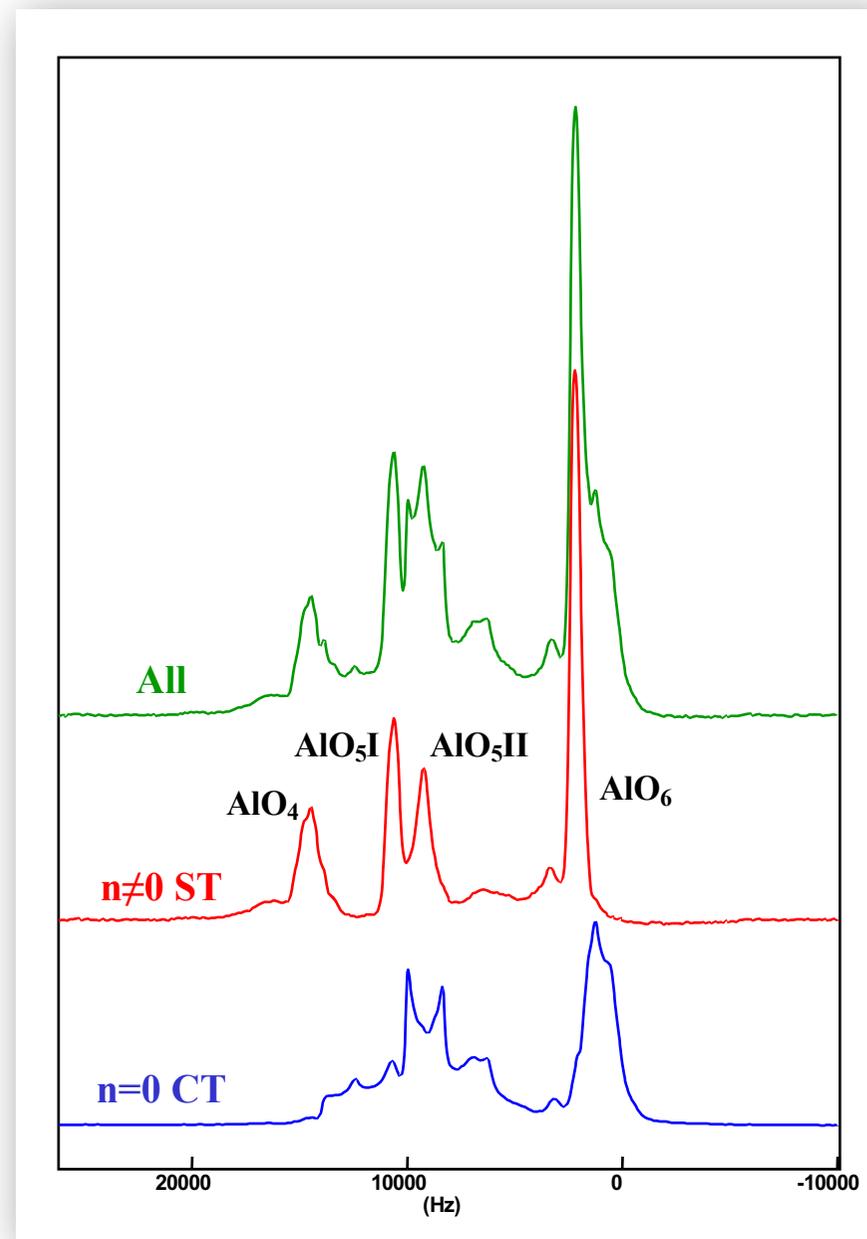


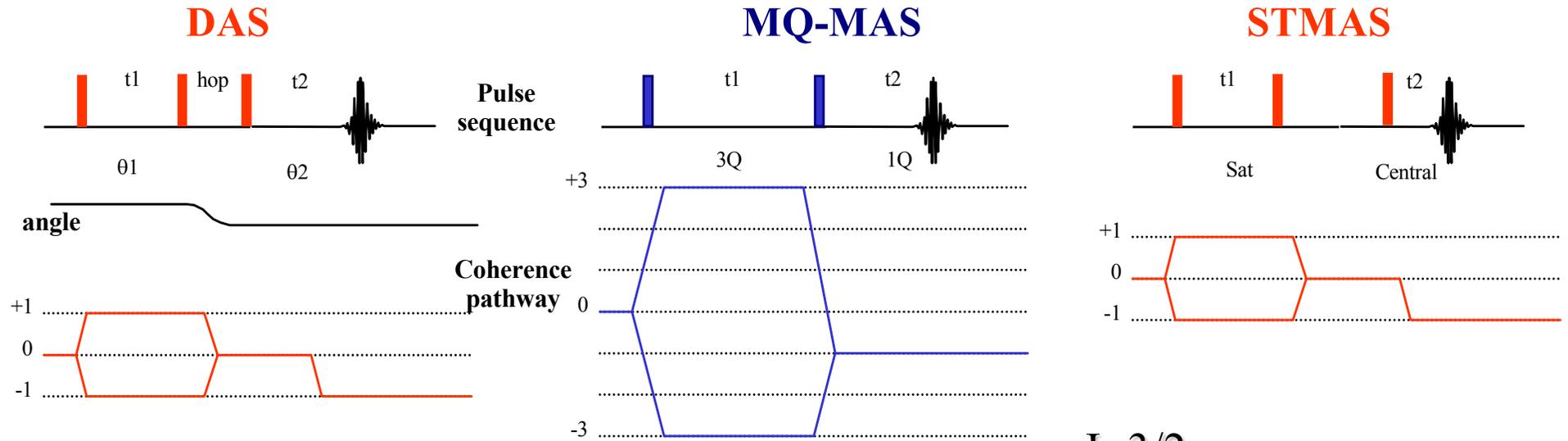


.. APM Kentgens.. *Solid State NMR* 5 175-180 **1995**

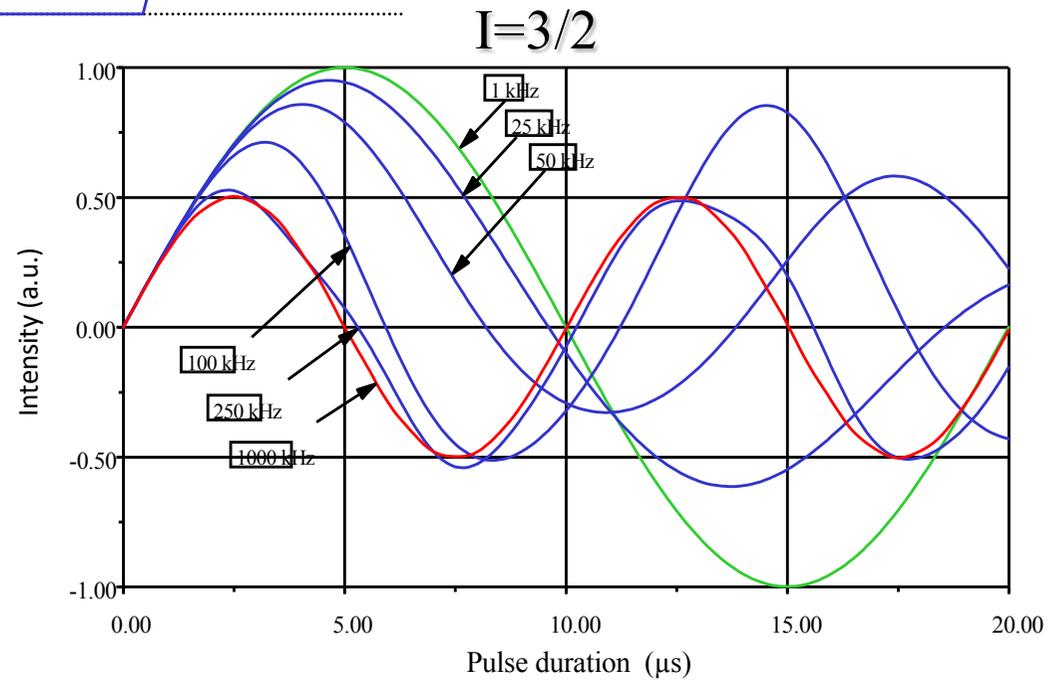
.. Z. Gan .. *J. Am. Chem. Soc.* 124 5634-5635 **2002**

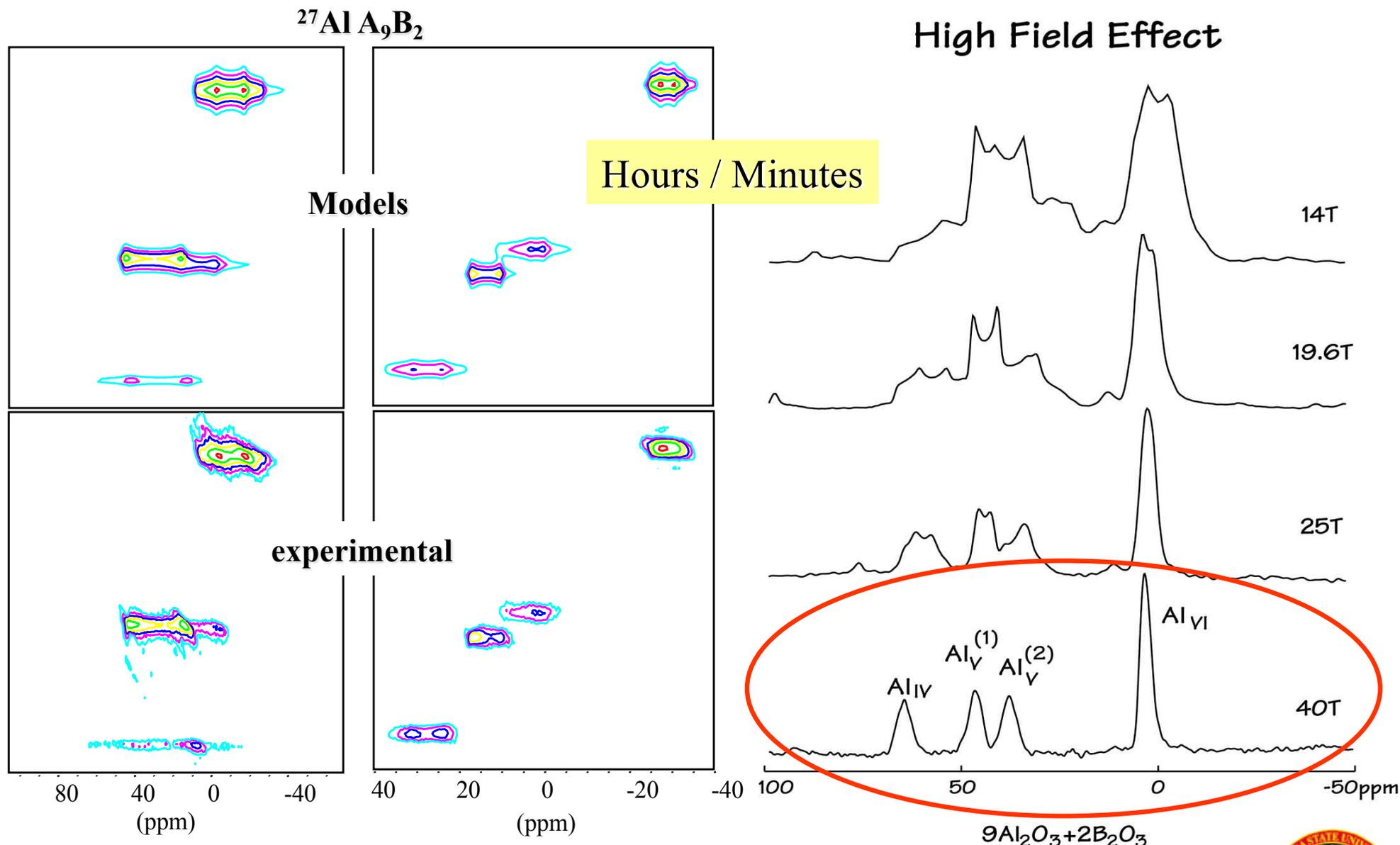
Z.Gan... *JMR* **2017**





$$\Omega_p = - \sum_{l=0,2,4} C_l^p A_l(\theta, \varphi) P_l(\cos \beta)$$





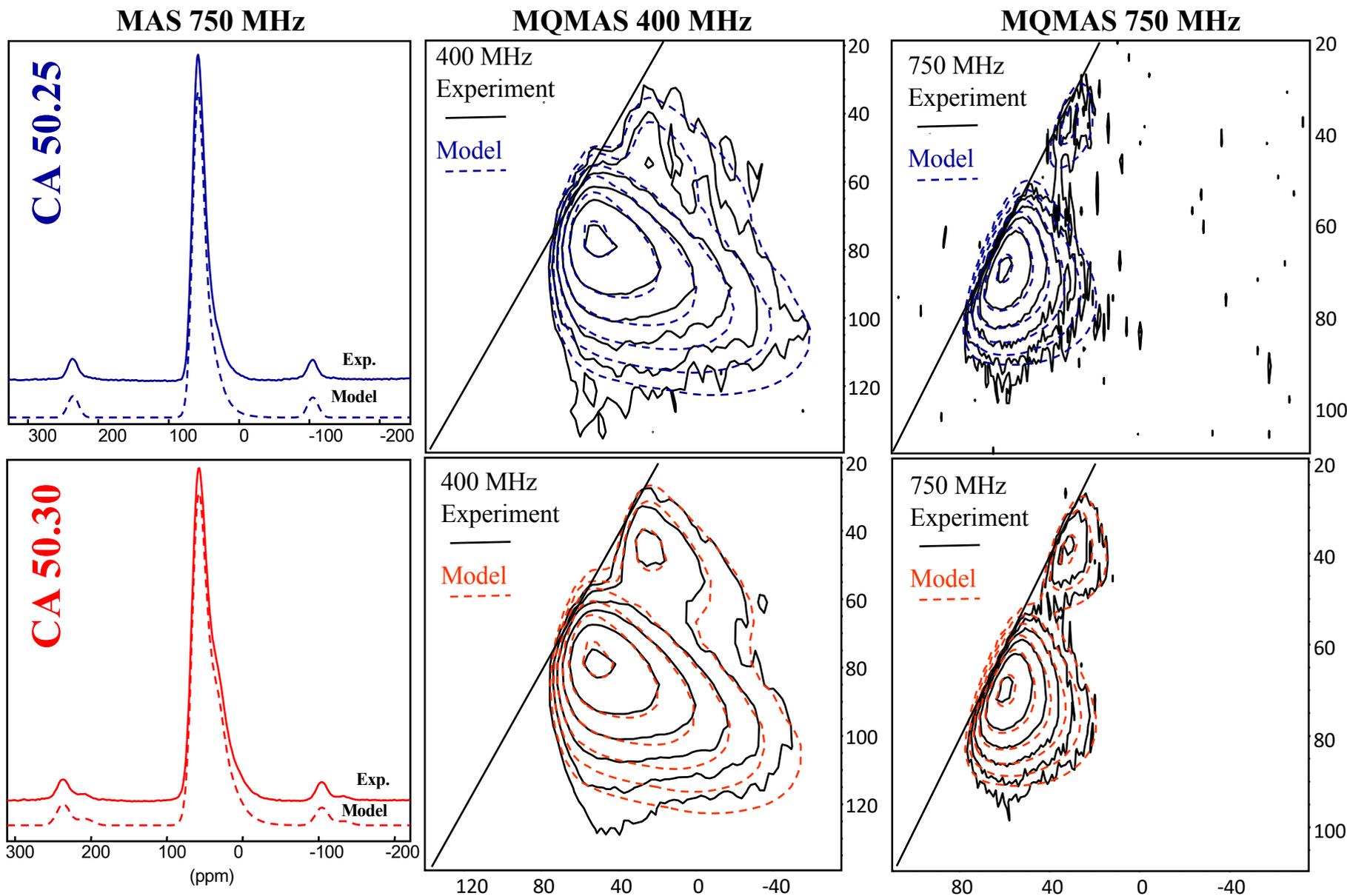
MQ-MAS 400

STMAS 830

NHMFL – Z. Gan



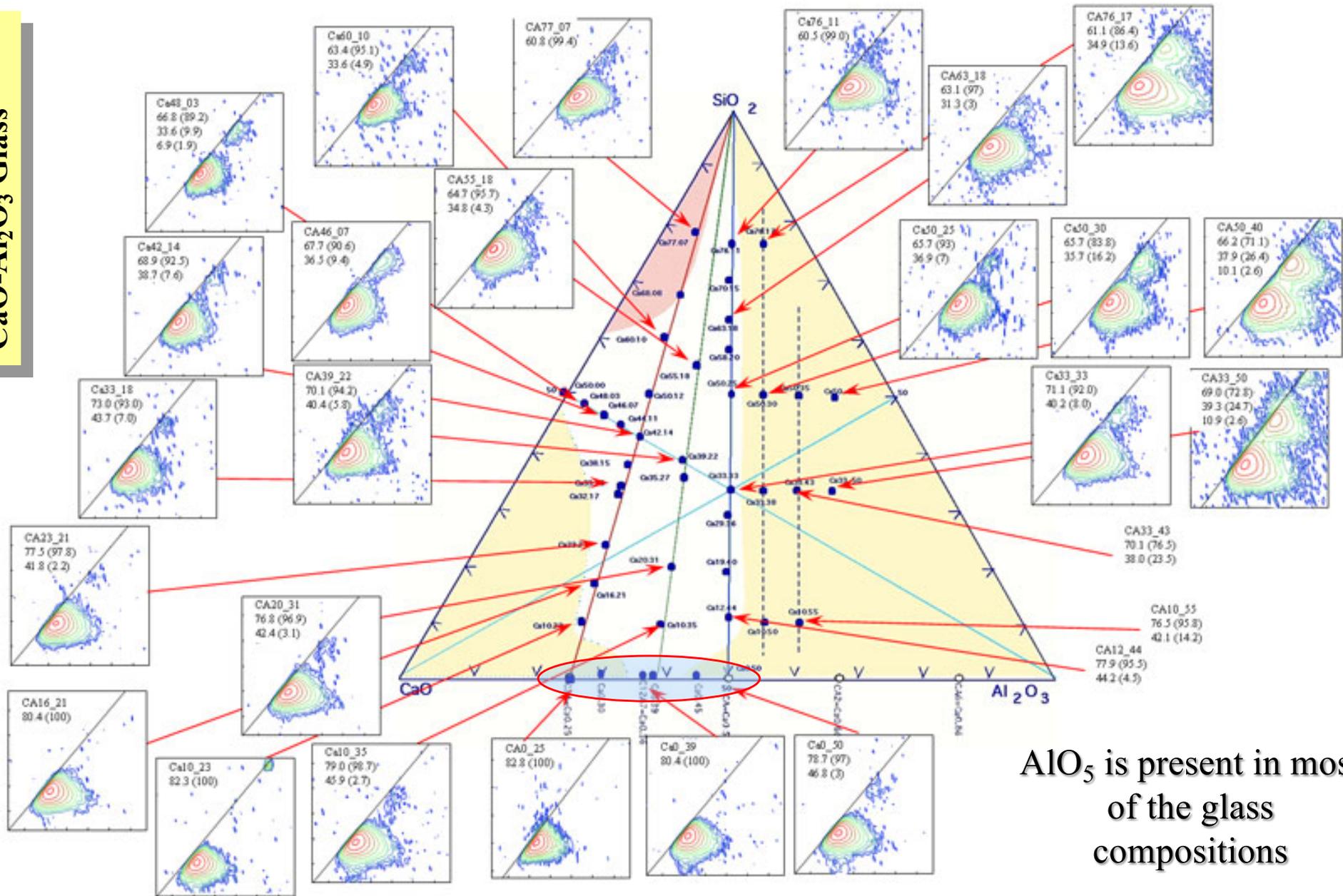
Glass (SiO₂, Al₂O₃, CaO)



D.R.Neuville, L.Cormier, D.Massiot 'Al coordination and speciation in calcium aluminosilicate glasses : effects of composition determined by ²⁷Al MQ-MAS NMR and Raman spectroscopy' [Chem. Geol. 229 173-185 \(2006\)](#)

D.R. Neuville, L. Cormier, D. Massiot 'Al environment in tectosilicate and peraluminous glasses: a NMR, Raman and XANES investigation' [Geochim. Cosmochim. Acta 68 5071-5079 \(2004\)](#)

CaO-Al₂O₃ Glass



AlO_5 is present in most of the glass compositions

1 particle size: crystallization

2 particle sizes: no long-range order

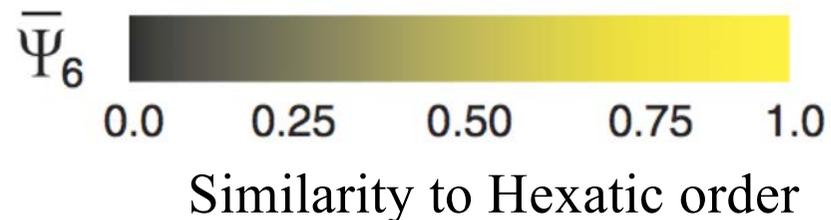
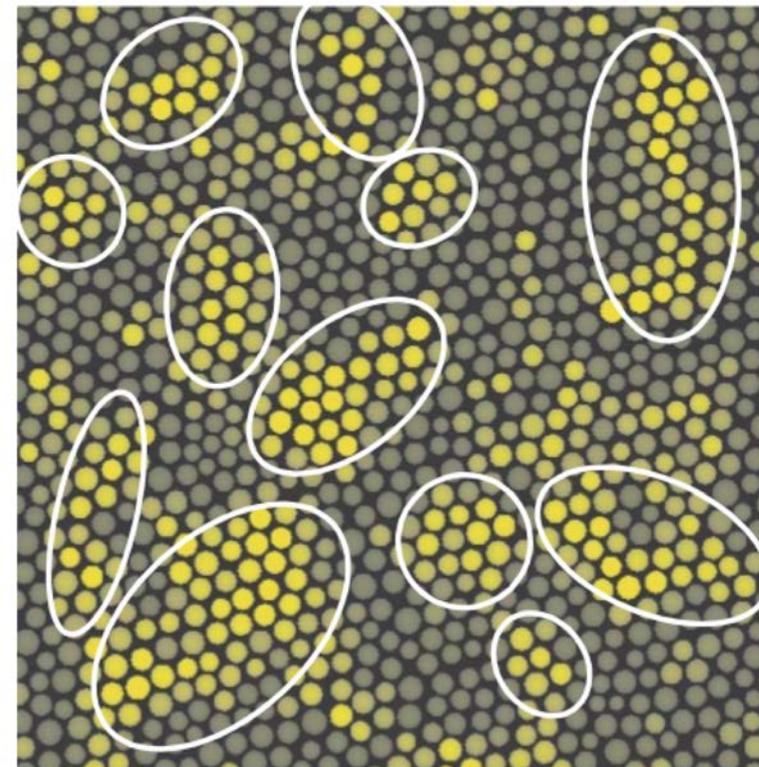
→ *Glass*

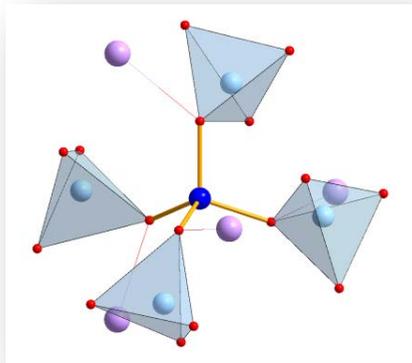
Hexatic order in a 2D colloidal glass with two particle sizes :

**Medium Range Crystalline Order
And
Locally Favored Structures LFS**

Q_n^m , SiO_5 , AlO_5 , $\mu_3\text{O}$, Al-O-Al ...

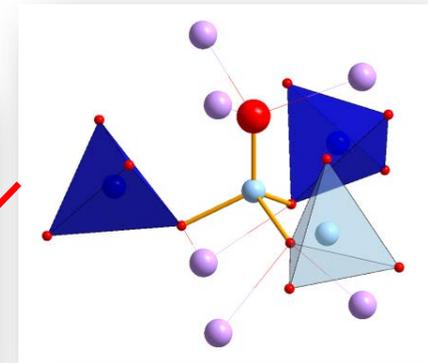
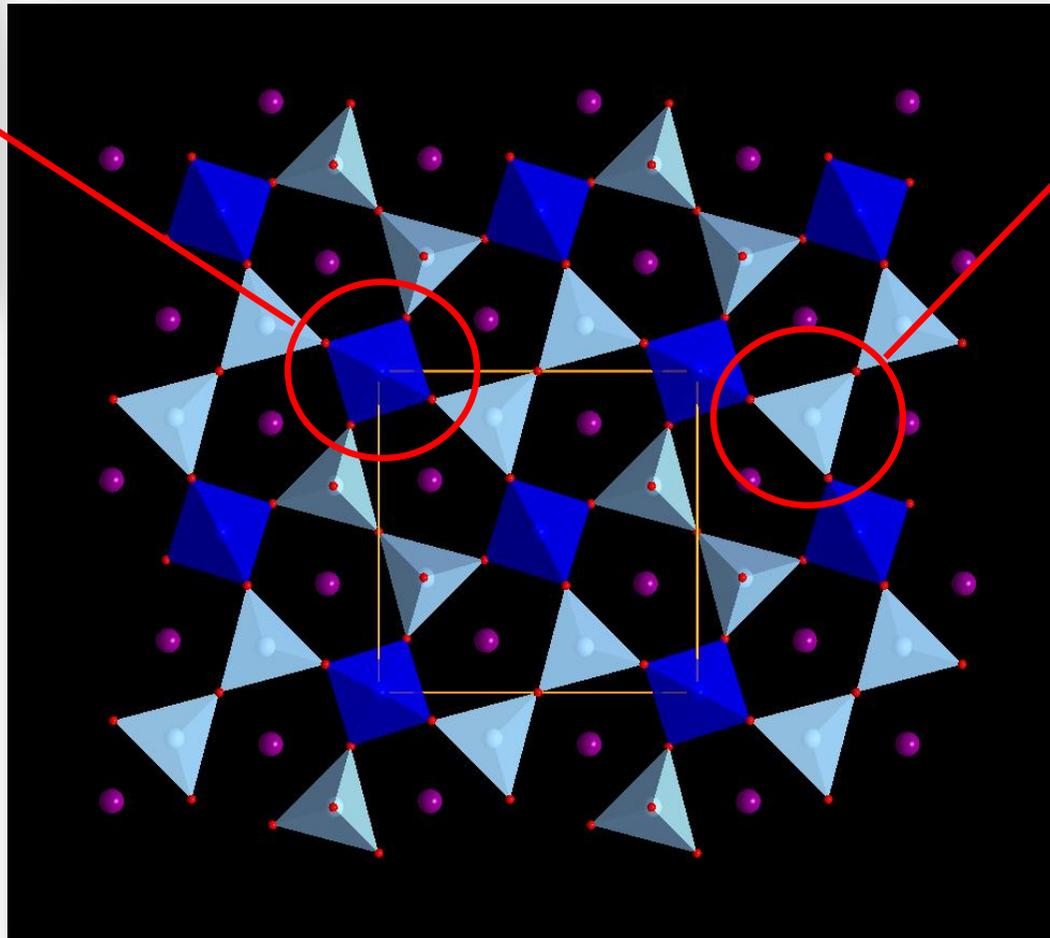
From the model to the real glass?





T_1 : Al only
5 configurations

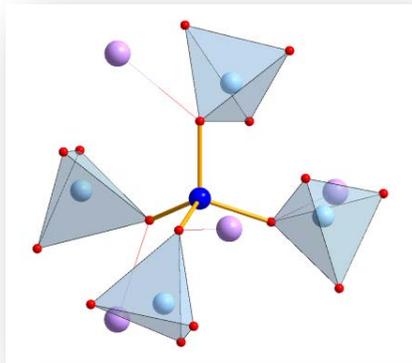
Al-Al₄
Al-SiAl₃
Al-Si₂Al₂
Al-Si₃Al
Al-Si₄



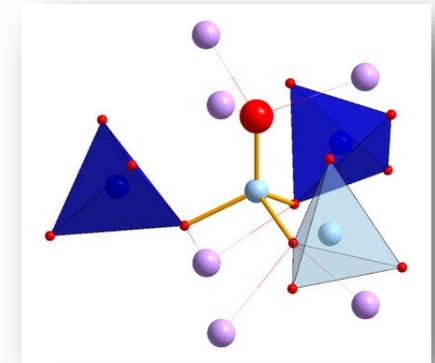
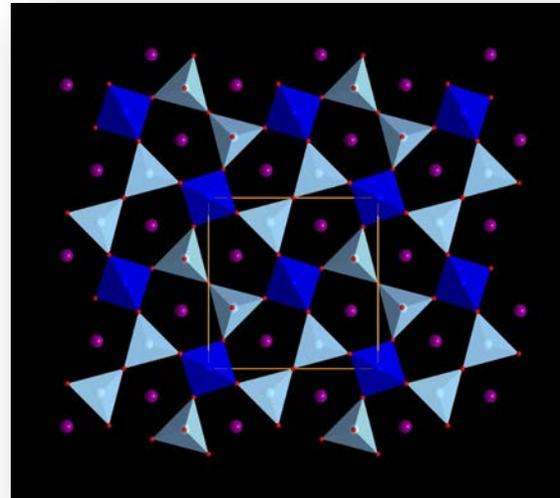
T_2 : (Al_{0.5},Si_{0.5})
2 configurations

Al-SiAl₂ Si-SiAl₂
Al-Al₃ Si-Al₃

Gehlenite Ca₂Al₂SiO₇



Gehlenite $\text{Ca}_2\text{Al}_2\text{SiO}_7$



^{27}Al MAS @ 17.6T

^{29}Si MAS

T_1 & T_2 Al

$\text{Al-Si}_2\text{Al}_2$

T_2 Si

Al-SiAl_2

Al- Al_4

Si-Al_3

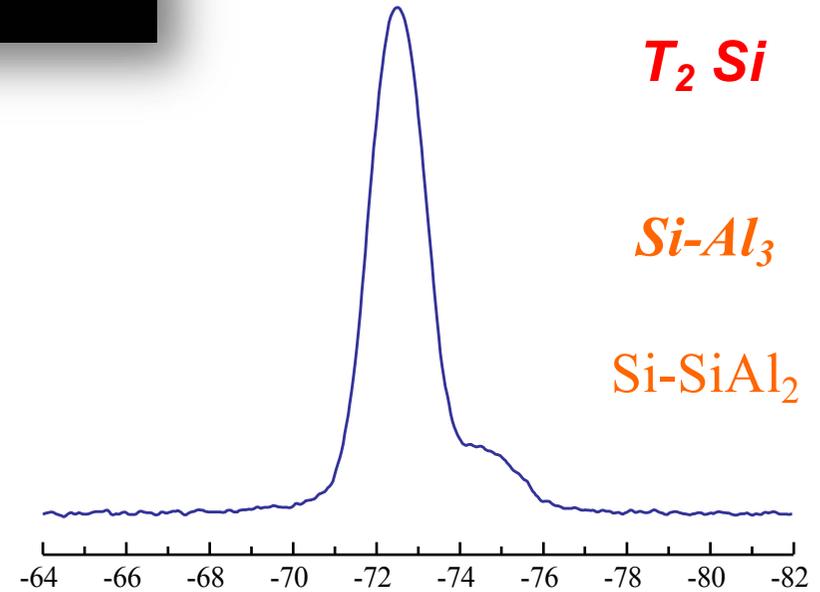
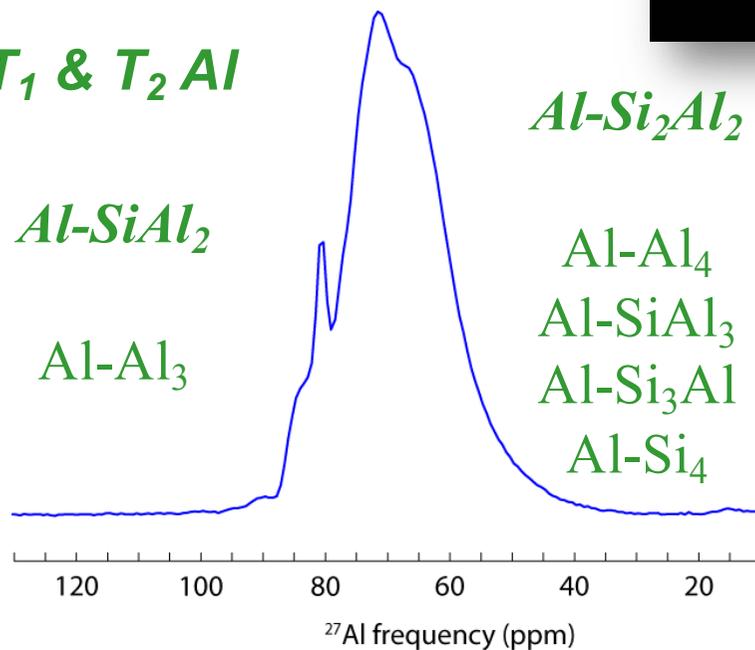
Al- Al_3

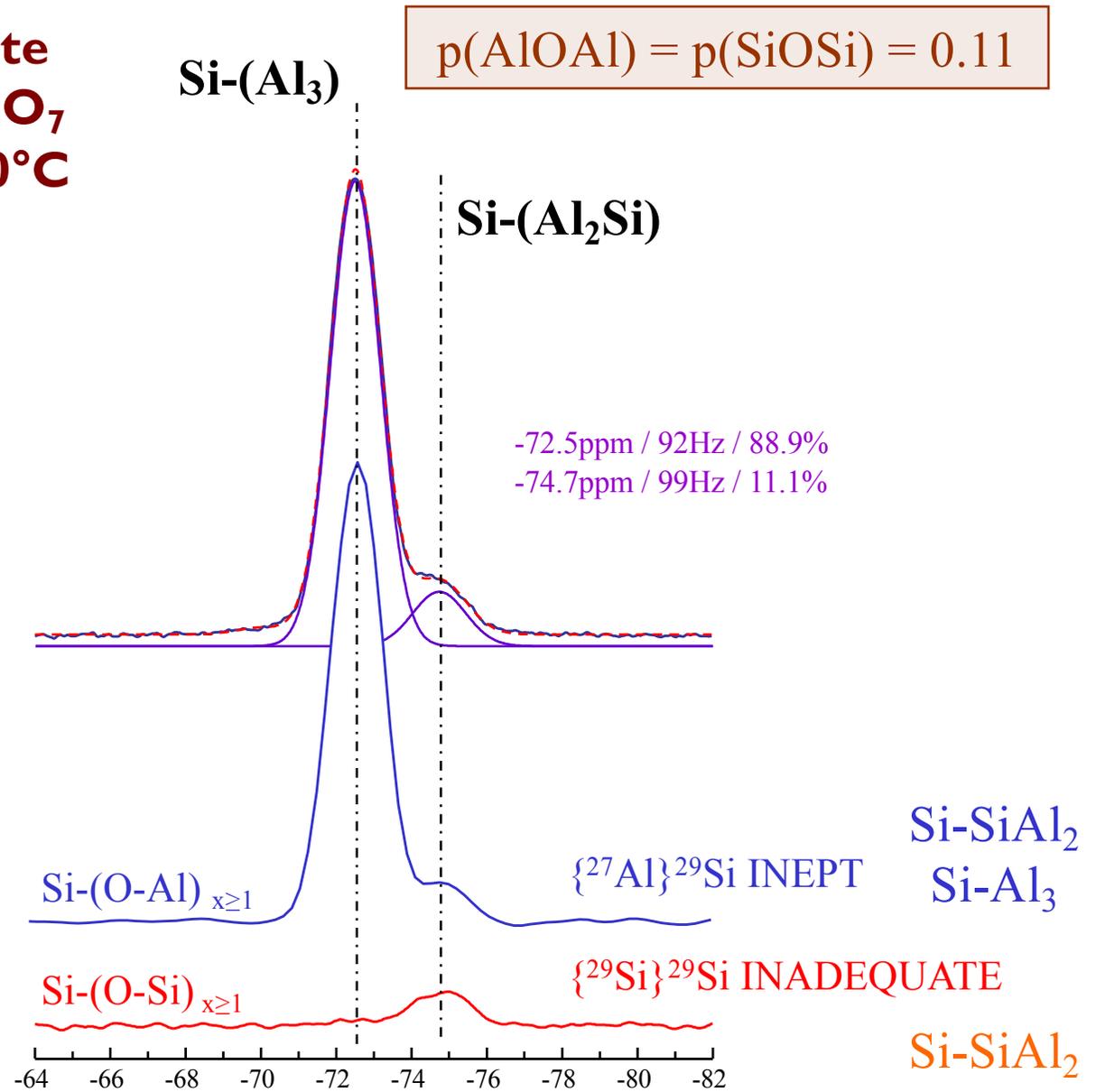
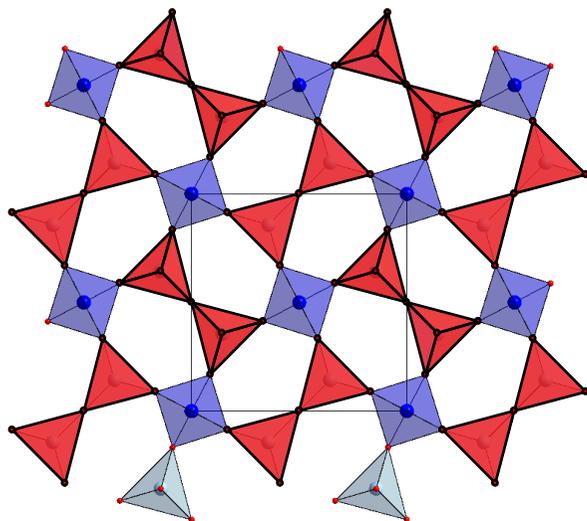
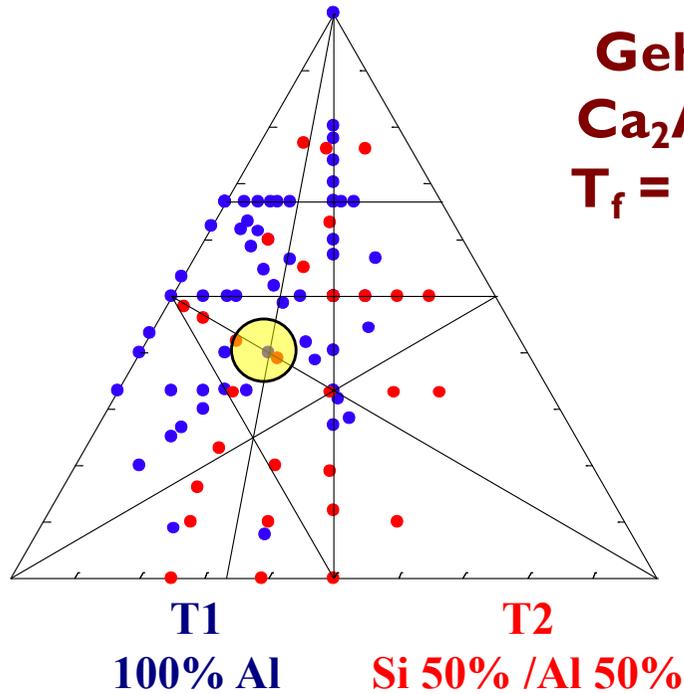
Al-Si Al_3

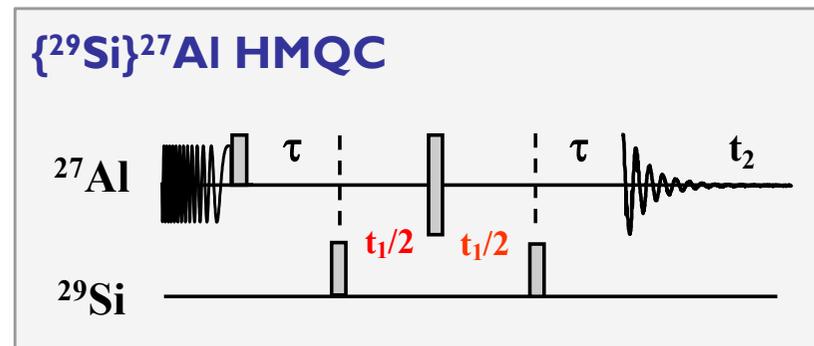
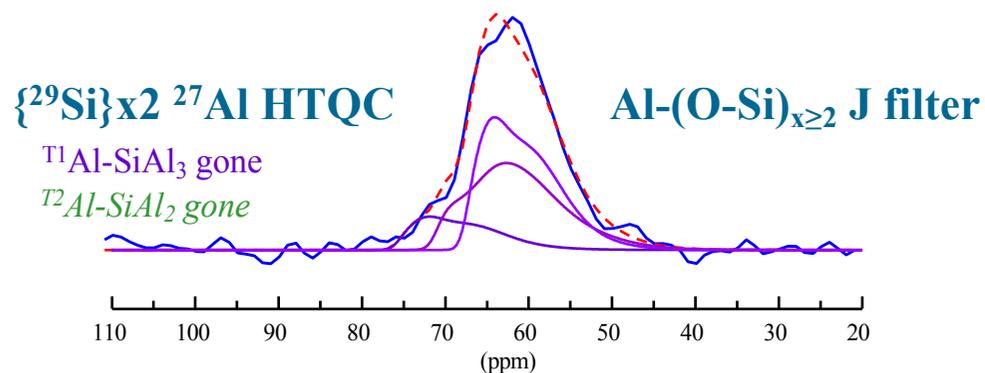
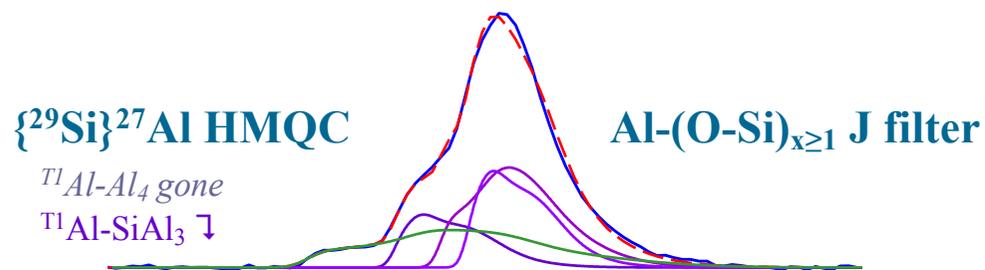
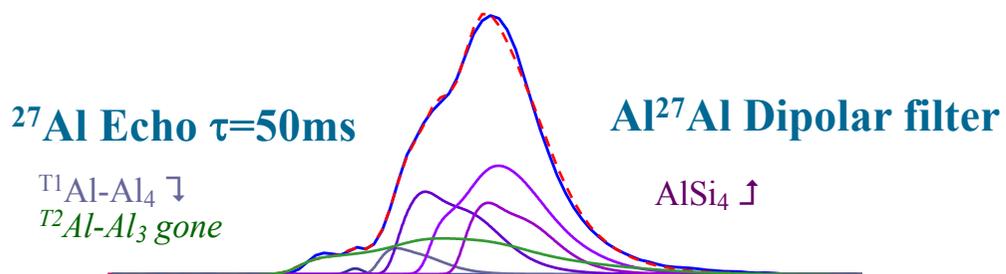
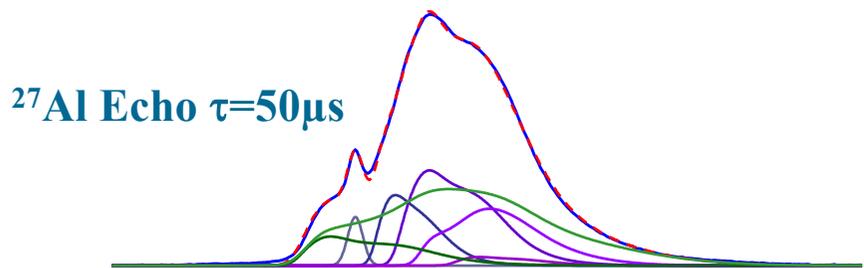
Al-Si $_3$ Al

Al-Si $_4$

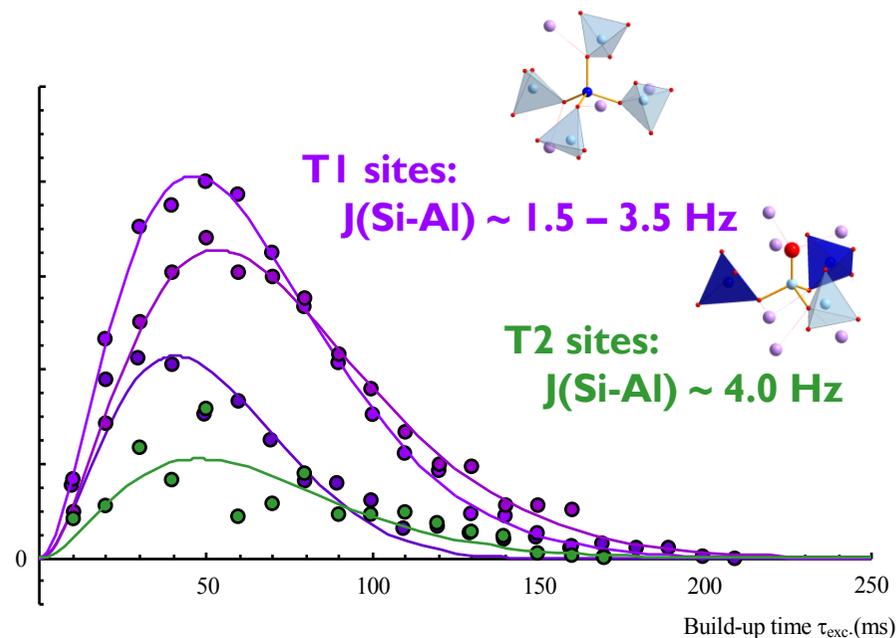
Si-Si Al_2







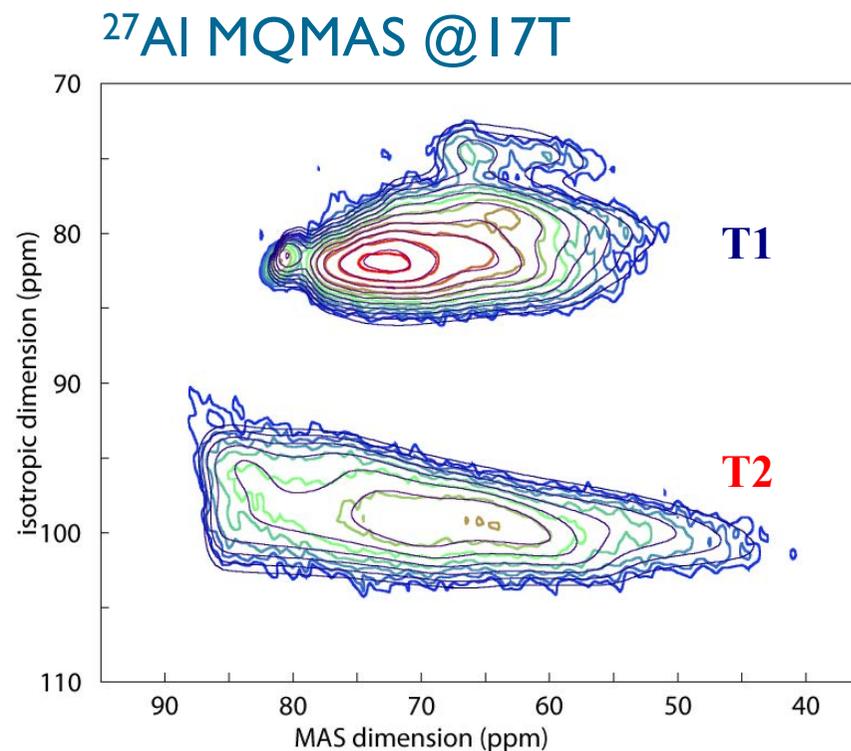
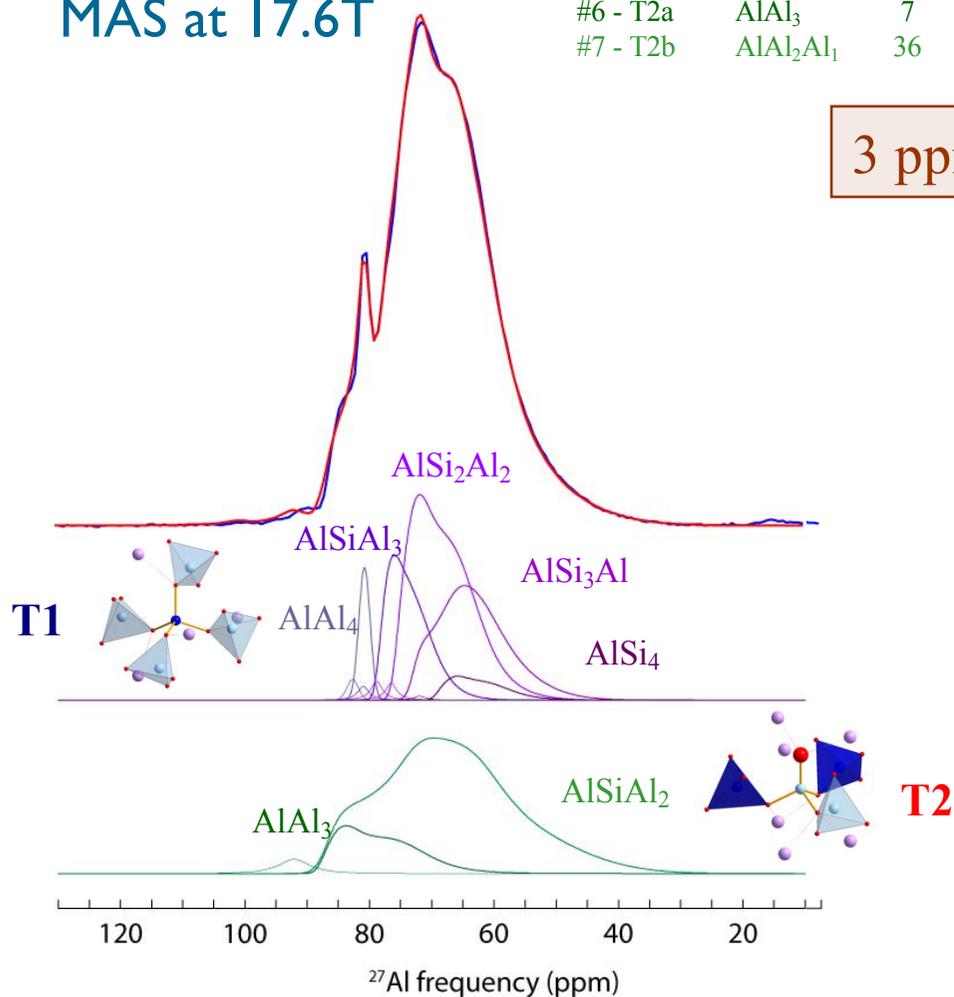
$\{^{29}\text{Si}\}^{27}\text{Al}$ J-coupling
→ Bond angles from DFT



^{27}Al quantitative
MAS at 17.6T

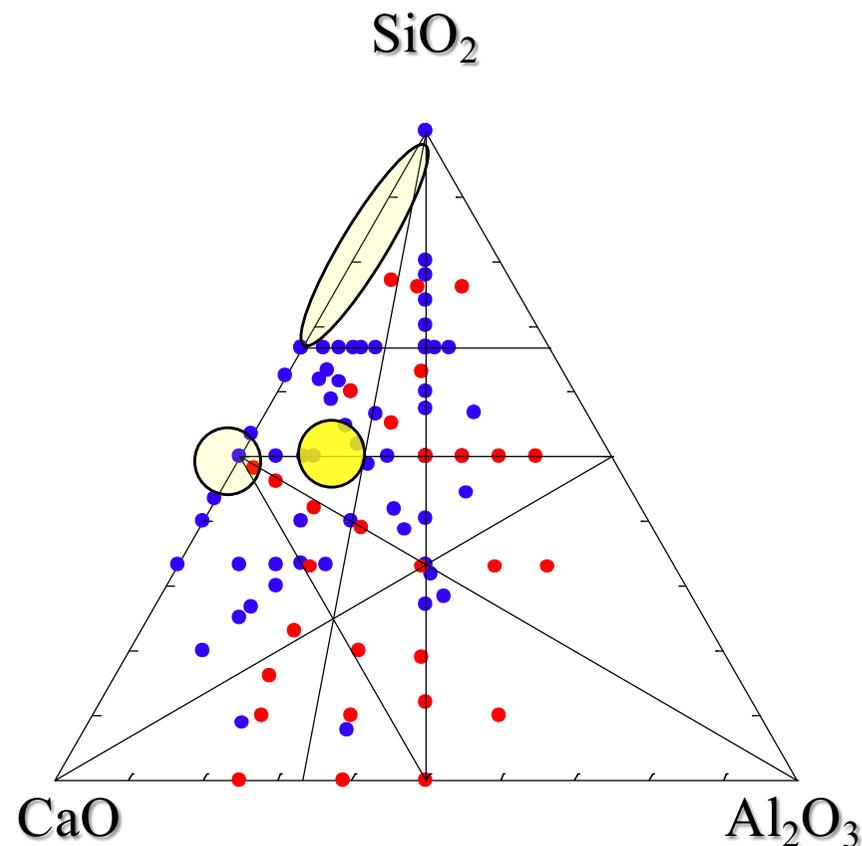
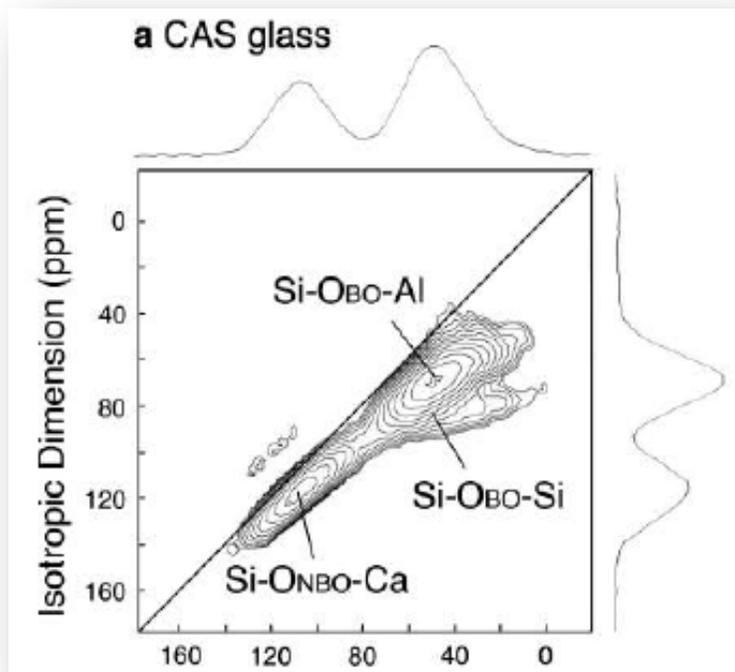
		%	δ_{iso} (ppm)	$\Delta\delta_{\text{iso}}$	C_Q (MHz)	ΔC_Q	η_Q	T_2 (ms)
#1 - T1a	AlAl_4	3	82.5	n/a	1.75	n/a	n/a	9 (± 1)
#2 - T1b	AlAl_3Si_1	11	79.2	1.50	5.82	2.00	0.3	21 (± 1)
#3 - T1c	AlAl_2Si_2	24	76.7	1.50	7.27	2.00	0.3	28 (± 1)
#4 - T1d	AlAl_2Si_3	16	73.4	1.50	7.59	2.00	0.6	47 (± 2)
#5 - T1e	AlSi_4	3	70.2	1.50	6.89	2.00	0.3	86 (± 12)
#6 - T2a	AlAl_3	7	89.4	1.48	8.31	2.20	0.24	12 (± 2)
#7 - T2b	AlAl_2Al_1	36	87.7	1.49	10.7	1.84	0.61	38 (± 2)

3 ppm shift per Si/Al substitution in T1 sites



Oxygen Speciation in Multicomponent Silicate Glasses Using Through Bond Double Resonance NMR Spectroscopy

Sohei Sukenaga,^{*,†} Pierre Florian,^{*,‡} Koji Kanehashi,[§] Hiroyuki Shibata,[†] Noritaka Saito,^{||} Kunihiko Nakashima,^{||} and Dominique Massiot[‡]



Lee et al., *Geochim. Cosmochim. Acta* **70** 4275-4286 (2006)

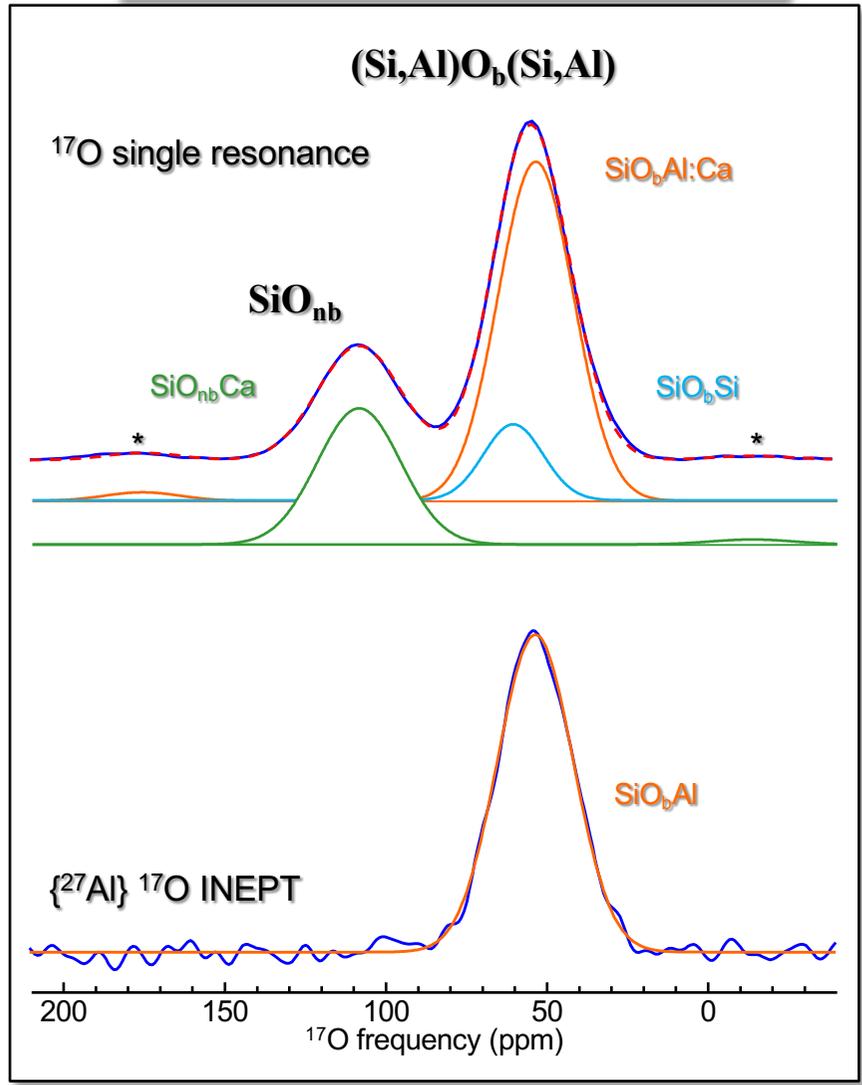
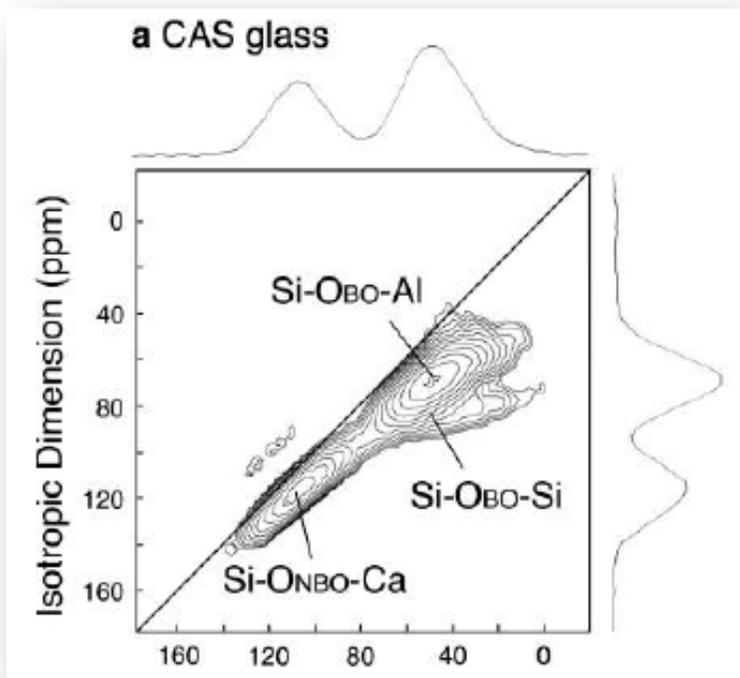
36.5 CaO – 51 SiO₂ – 12.5 Al₂O₃

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Oxygen Speciation in Multicomponent Silicate Glasses Using Through Bond Double Resonance NMR Spectroscopy

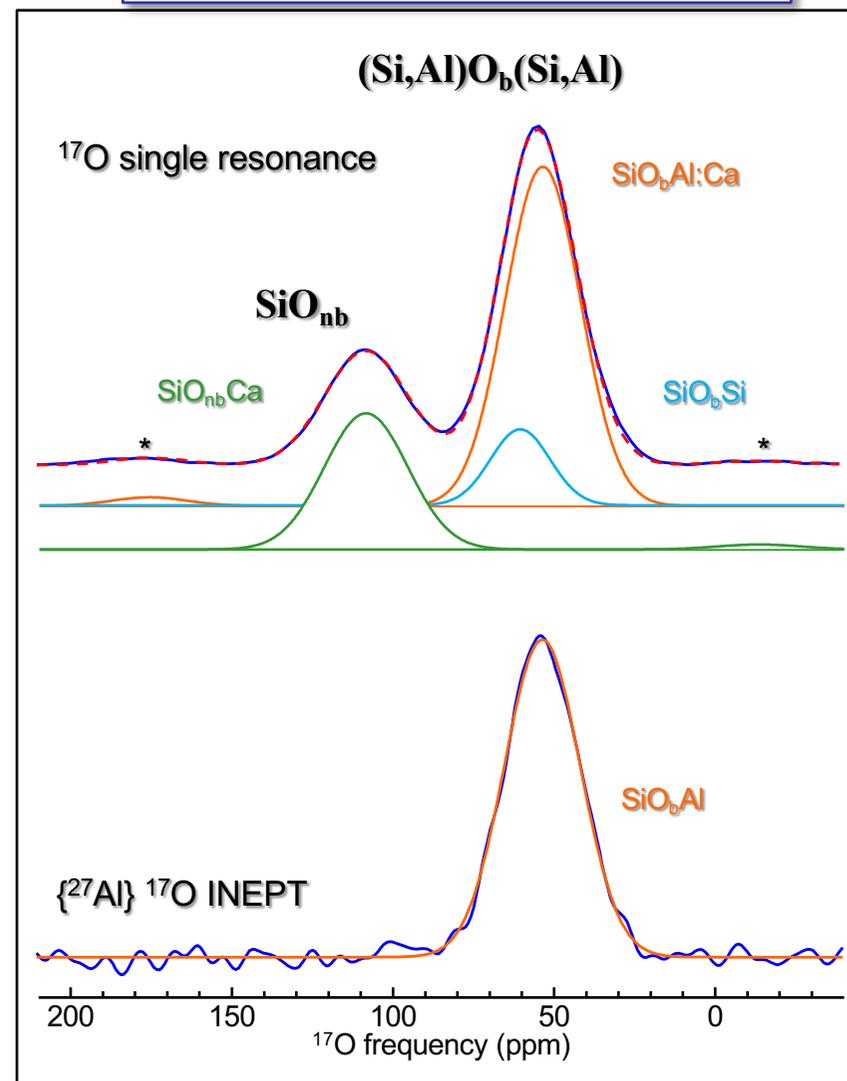
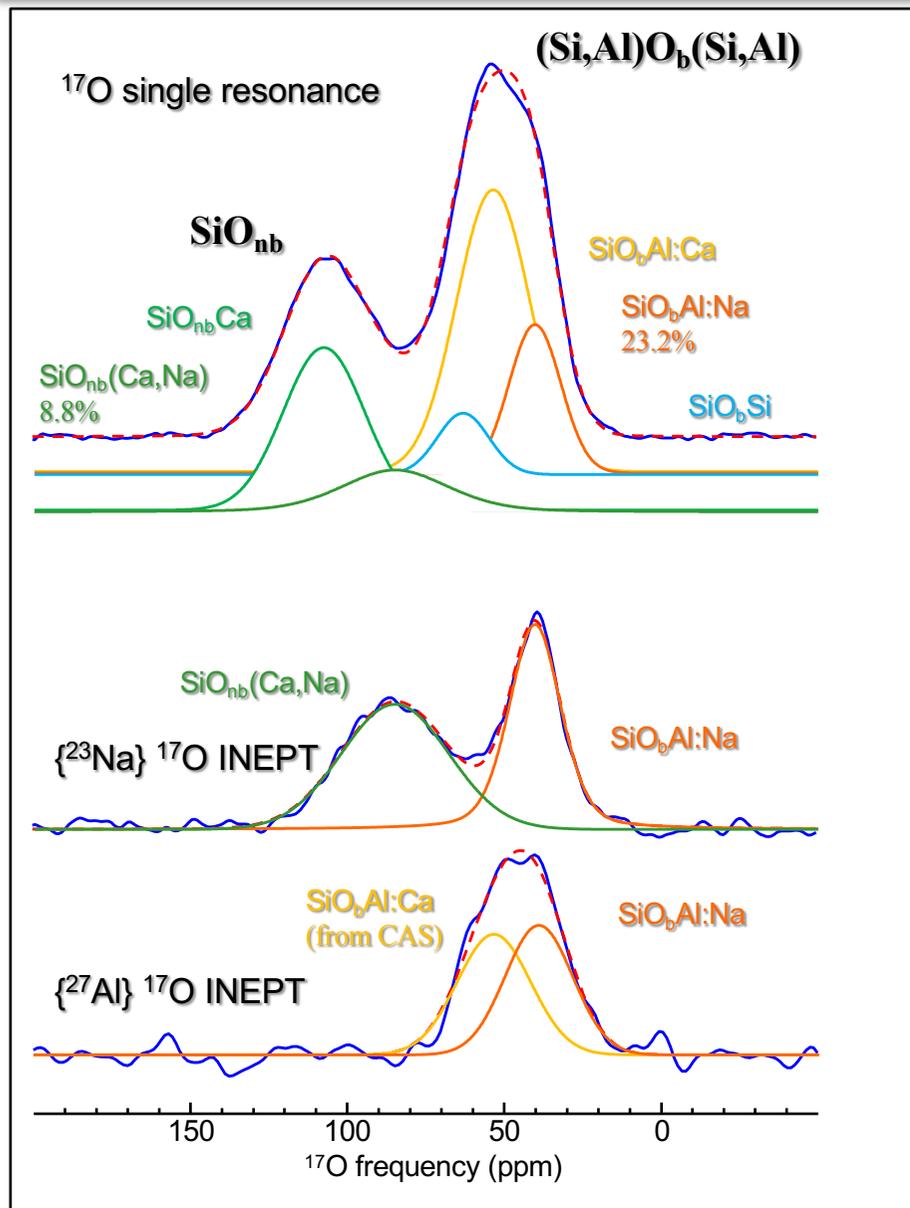
Sohei Sukenaga,^{*†} Pierre Florian,^{*‡} Koji Kanehashi,[§] Hiroyuki Shibata,[†] Noritaka Saito,^{||} Kunihiro Nakashima,^{||} and Dominique Massiot[‡]

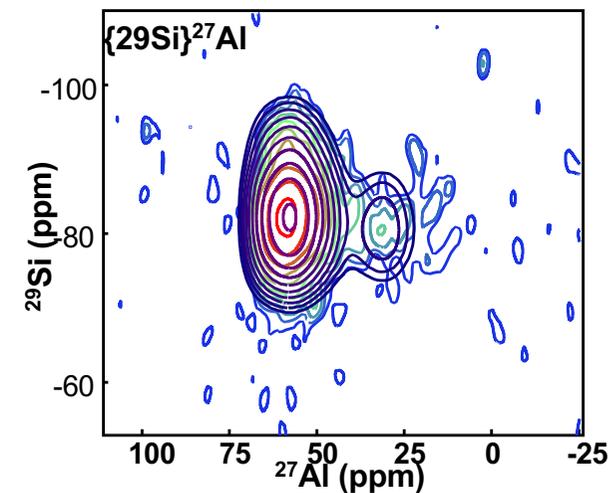
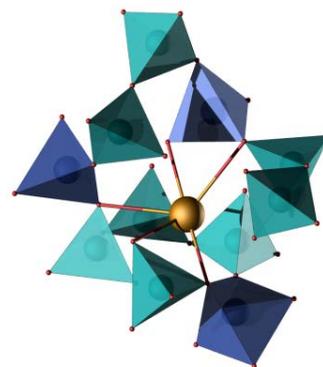
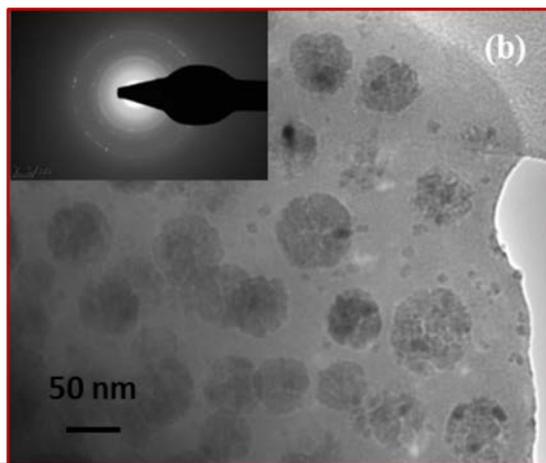
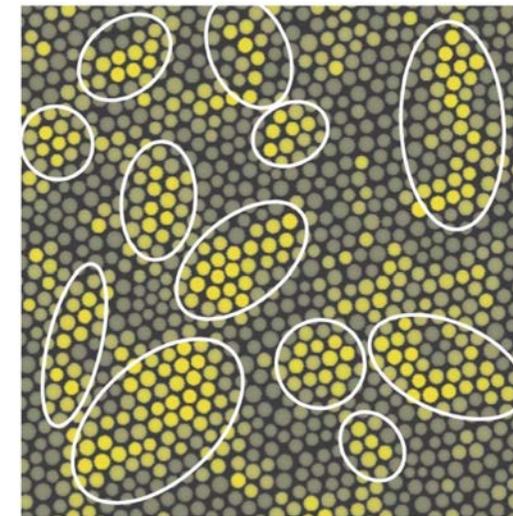
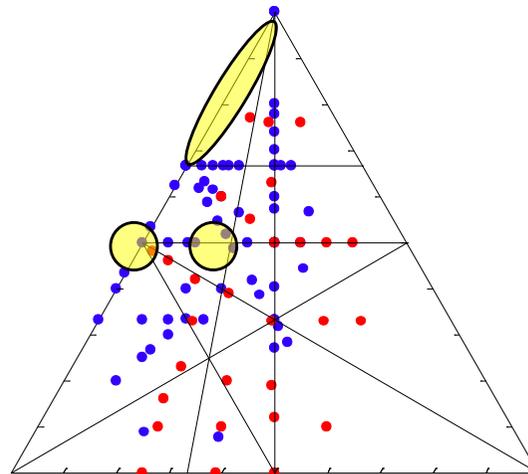
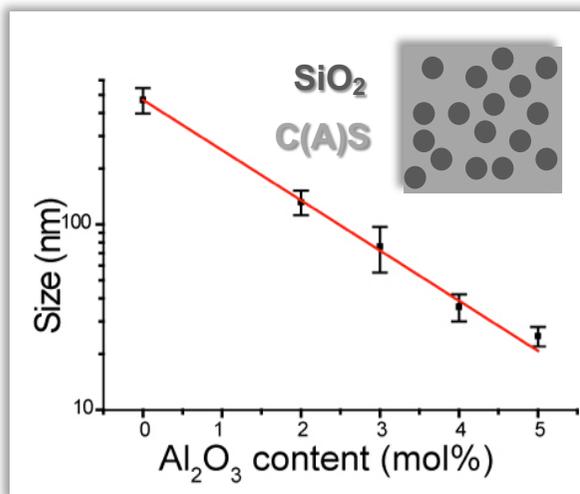


Lee et al., *Geochim. Cosmochim. Acta* **70** 4275-4286 (2006)

10.8 Na₂O – 31.9 CaO – 44.7 SiO₂ – 12.6 Al₂O₃

36.5 CaO – 51 SiO₂ – 12.5 Al₂O₃

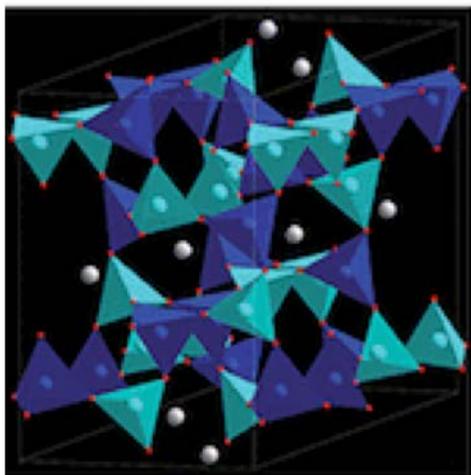
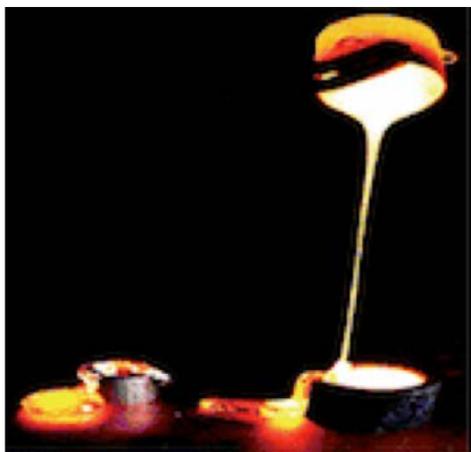




Phase separation

Network Topology

Local structures



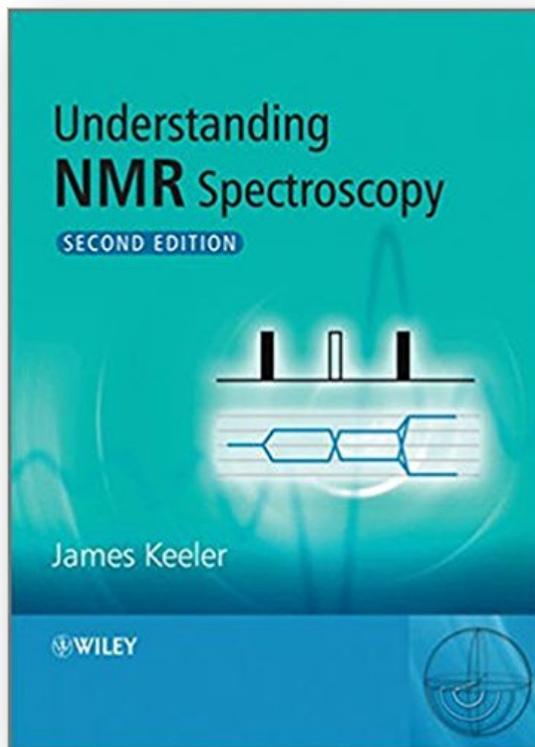
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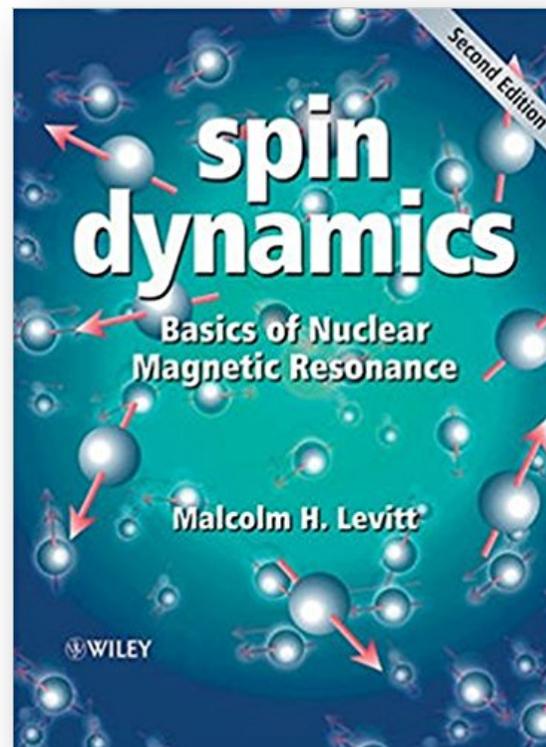
Topological, Geometric, and Chemical Order in Materials: Insights from Solid-State NMR

DOMINIQUE MASSIOT,* ROBERT J. MESSINGER,
SYLVIAN CADARS, MICHAËL DESCHAMPS,
VALERIE MONTOUILLOUT, NADIA PELLERIN,
EMMANUEL VERON, MATHIEU ALLIX, PIERRE FLORIAN, AND
FRANCK FAYON

Accounts Chem. Res. **46** 1975–1984 (2013)



James Keeler – University of Cambridge
<http://www-keeler.ch.cam.ac.uk/lectures/>



Malcom Levitt – University of Southampton
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